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Defect-driven redox interplay on anatase TiO₂: surface- structure dependent activation for CO₂ hydrogenation catalysis

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Abstract

Titanium dioxide (TiO₂) is one of the most extensively studied oxides as an active catalyst or catalyst support, particularly in energy and environmental applications, but the atomistic mechanisms governing its dynamic response to reactive environments and their correlation to reactivity remain largely elusive. Using in-situ environmental transmission electron microscopy (ETEM), synchrotron X-ray diffraction (XRD), ambient-pressure X-ray photoelectron spectroscopy (AP-XPS), reactivity measurements, and theoretical modeling, we reveal the dynamic interplay between oxygen loss and replenishment of anatase TiO₂ under varying reactive conditions. Under

29 H₂ exposure, anatase TiO₂ undergoes surface reduction via lattice oxygen loss, forming Ti₃O₅. In contrast, CO₂
30 exposure induces oxygen replenishment, reversing stoichiometry. In mixed H₂/CO₂ environments, the reverse
31 water-gas shift (RWGS) reaction proceeds selectively on stepped and high-indexed TiO₂ surfaces, whereas the
32 thermodynamically stable TiO₂(101) surface remains inactive and intact. Critically, H₂ pre-treatment generates
33 oxygen vacancies on TiO₂(101), transforming it into an active Ti₃O₅ or defect-rich surface that catalyzes RWGS.
34 By correlating surface structure, defect dynamics, and gas-phase interactions, this work deciphers the competition
35 between H₂-driven reduction and CO₂-driven oxidation pathways at the atomic scale. These insights establish defect
36 engineering as a strategic lever to activate inert TiO₂ facets, advancing the design of adaptive catalysts for
37 sustainable fuel synthesis technologies.

38

39 **Keywords:** Anatase TiO₂, O vacancy, Oxygen loss and replenishment, Reverse water-gas shift reaction, Surface
40 reduction, In-situ TEM

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43 Introduction

44 Heterogeneous catalysts, particularly those consisting of oxide-supported metals or metal oxides, are
45 essential for producing commodity chemicals and fuels. Key reactions such as CO₂ hydrogenation, the reverse
46 water-gas shift (RWGS) reaction to methanol, and CO oxidation have been extensively investigated using catalytic
47 systems like Cu/ZnO/Al₂O₃,¹ Ni/CeO₂,^{2,3} Pt/CeO₂,^{4,5} Ni/TiO₂,⁶⁻⁸ Rh/TiO₂,^{9,10} Pt/TiO₂,¹¹ and Au/TiO₂.^{12,13} While
48 significant advancements have been made in identifying active sites and unraveling structure-reactivity
49 relationships,^{14,15} a fundamental atomic-scale understanding of dynamic surface and interfacial processes under
50 reactive conditions, such as interfacial chemical reactions and metal-support interactions, remains elusive. A critical
51 gap persists in the oversimplified view of oxide supports, which are often regarded merely as inert structural
52 stabilizers or oxygen reservoirs, overlooking their intrinsic physicochemical properties and active participation in
53 catalytic cycles.^{16,17} This work focuses on elucidating the dynamic behavior of oxide supports themselves, aiming
54 to uncover the atomic-scale mechanisms by which their surfaces and interfaces govern catalytic reactivity.

55 Anatase TiO₂ is widely utilized as an oxide support in heterogeneous catalysis, owing to its inherent
56 reducibility and sub-stoichiometric composition under reducing conditions. Previous ETEM studies have revealed
57 complex surface reconstructions and chemical dynamics on TiO₂ and other oxides under reactive environments,¹⁸⁻
58 ²⁰ highlighting the highly dynamic nature of metal oxides. These properties critically govern its physicochemical
59 behavior and catalytic performance. Notable examples are metal atoms or nanoparticles loaded onto different
60 anatase TiO₂ facets, which yield site-specific reactivity and stability due to their unique coordination
61 environments.^{21,22} Using anatase TiO₂ as a prototypical system, this study employs in-situ ETEM to capture real-
62 time atomic-scale dynamics of the RWGS reaction on TiO₂ surfaces. By controlling CO₂/H₂ gas mixtures at
63 operationally relevant temperatures and pressures, we track spatiotemporal evolution of atomic configurations, from
64 surface restructuring to bulk transformations, under reactive conditions. Along with complementary synchrotron
65 XRD, AP-XPS and reactivity measurements, we reveal the dynamic interplay of lattice oxygen loss and
66 replenishment, as well as facet-dependent redox competition during catalysis.

67 Density Functional Theory (DFT) and machine-learning interatomic potentials (MLIPs) further map
68 structural phase evolution across anatase surfaces under varying oxygen chemical potentials. These calculations

reveal that the (112) facet undergoes reduction more readily than the (101) facet. Electronic structure analysis demonstrates how surface reduction alters H and CO₂ adsorption/dissociation energetics, governing reactivity and driving oscillatory transitions between oxidized and reduced states during catalysis. Our findings highlight the crucial role of oxygen vacancies in modulating the catalytic reactivity of anatase TiO₂—a behavior further governed by surface orientations and defect structure. By bridging atomic-scale dynamics to macroscopic reactivity, this work establishes a framework for engineering oxide and oxide-supported catalysts with adaptive surfaces for chemical synthesis and energy conversion.

Results and Discussion

2.1 Structural and Chemical Evolutions of Anatase TiO₂ under Reactive Conditions

Anatase TiO₂ nanoparticles (NPs) (99.9% purity, ~ 30 nm average diameter) were purchased from Research Nanomaterials Inc. Brunauer-Emmett-Teller (BET) analysis yielded a surface area of 49.4 m²/g using physical adsorption of N₂ at 77 K, corresponding to an estimated particle diameter of 28.8 nm (assuming anatase density: 4.23 g/cm³). This aligns with TEM observations, which show an average particle diameter of ~ 31.5±2 nm (Supplementary Fig. S1). Prior to in-situ TEM experiments, NPs were pretreated in an O₂ environment (pO₂ = 5×10⁻³ Torr, 400 °C, 30 min)) to combust surface carbon contamination.

Figs. 1(a-f) depict a time-sequenced series of high-resolution TEM (HRTEM) images capturing the dynamic surface evolution of TiO₂ under alternating H₂ or CO₂ environments, prior to combining them for RWGS conditions. Initial O₂ pretreatment at 400 °C produces a stoichiometric TiO₂ surface with ordered atomic columns, confirmed by the fast Fourier transform (FFT) in Fig. 1a. Subsequent exposure to 3×10⁻³ Torr H₂ at 500 °C induces progressive O loss, initiating at the surface and propagating inward (Fig. 1b, c). The advancing reduction front (marked by cyan dashed lines in Fig. 1) correlates with the formation of Ti₃O₅(312) lattice fringes (more examples in Supplementary Fig. S2, S3), consistent with lattice O removal under reducing conditions. The transition from TiO₂ to Ti₃O₅ is accompanied by a theoretical volume shrinkage ~ 10.2% (see calculation details in Supplementary information Section 2), reflecting the structural compaction associated with oxygen vacancy formation and lattice

reconfiguration. With the formation of Ti_3O_5 layer, the further reduction becomes increasingly limited and eventually halts at a critical thickness, because the following oxygen removal may become diffusion-limited, governed by the transport of oxygen vacancies through the existing Ti_3O_5 layer (see more examples and discussion in Supplementary information Fig. S2). This self-limiting behavior is further supported by temperature-programmed reduction (TPR) measurements (Fig. S4), which reveal two distinct H_2 consumption peaks at 410 °C and 590 °C. These peaks are attributed to surface and subsequent bulk reduction processes, respectively, underscoring the kinetically hindered, depth-dependent reduction of anatase TiO_2 under the H_2 reduction process at 500 °C.

Upon switching to 3×10^{-3} Torr CO_2 , dissociative adsorption of CO_2 occurs on the reduced Ti_3O_5 surface, regenerating lattice O and releasing CO gas.²³ This O replenishment is evidenced by the resulting transformation of the Ti_3O_5 lattice into the TiO_2 structure, as seen in the recovery of well-defined atomic columns and relatively uniform image contrast in Figures 1(d–f). The interface between Ti_3O_5 and bulk TiO_2 (marked by cyan dashed lines) shifts from the bulk toward the surface (red arrows, Figs. 1(d–f)), driven by O reintegration. The corresponding FFTs confirm this structural healing, as reflections corresponding to $\text{Ti}_3\text{O}_5(312)$ vanish (Fig. 1h). However, the reoxidation ($\text{Ti}_3\text{O}_5 \rightarrow \text{TiO}_2$) often remains incomplete, particularly in larger particles (see more examples in Fig. S6), due to kinetic limitations in oxygen transport. The smeared contrast at the center shown in Fig. 1(f) is consistent with residual reduced domains or defect-rich regions, indicating non-uniform reoxidation and potential trapping of vacancy clusters or defects within the particle core. Notably, the CO_2 -driven healing process completes in ~ 60 s—three times faster than the time period (~ 200 s) of the Ti_3O_5 formation from H_2 -induced reduction—underscoring the enhanced reactivity of the reduced Ti_3O_5 surface toward CO_2 dissociation. More experiments at 400 °C corroborate this trend, demonstrating consistent CO_2 adsorption and dissociation on Ti_3O_5 (Supplementary Fig. S5). Figure 1(g) presents a representative HRTEM image (from in situ TEM observation in Supplementary Fig. S3) of the reduced anatase TiO_2 surface, revealing Ti_3O_5 formation viewed along the $[352]$ zone axis. The corresponding FFT (Fig. 1h) aligns with Ti_3O_5 formation observed during H_2 exposure (Figs. 1a–c). Fig. 1i shows the atomic configurations of the parent TiO_2 and reduced Ti_3O_5 , matching the crystallographic orientations in the HRTEM images in Fig. 1(g).

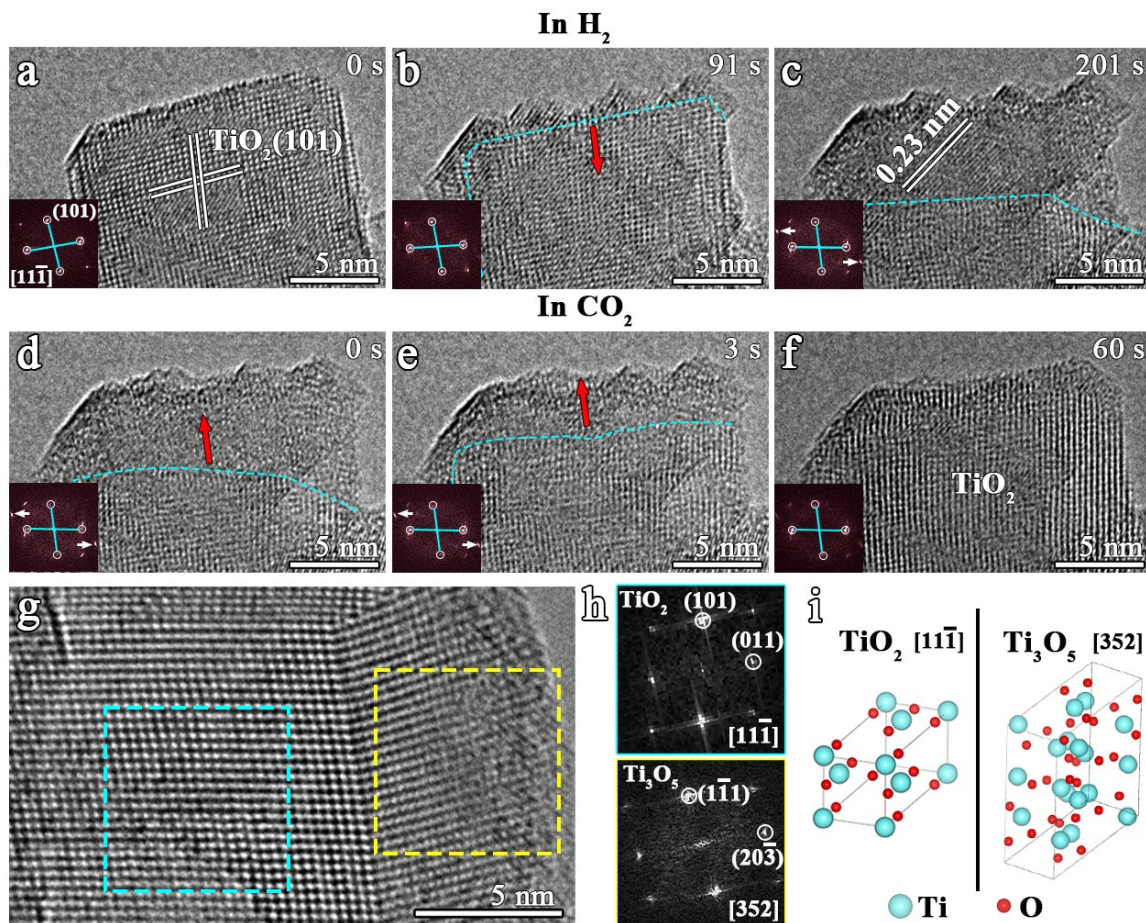


Figure 1. In-situ HRTEM imaging of TiO_2 reduction under H_2 and dynamic recovery under CO_2 . **a-c**, Time-sequence HRTEM images capturing the progressive reduction of anatase TiO_2 to Ti_3O_5 (500°C , $\approx 3 \times 10^{-3}$ Torr H_2). The reduction front propagates from the surface into the bulk. **d-f**, Structural recovery of reduced TiO_2 under CO_2 (500°C , $\approx 3 \times 10^{-3}$ Torr CO_2). O replenishment drives the $\text{Ti}_3\text{O}_5/\text{TiO}_2$ interface (cyan dashed lines) from bulk to surface (red arrows), restoring stoichiometry and eliminating Ti_3O_5 diffraction features in FFTs. **g, h**, Representative HRTEM image and the corresponding FFT of reduced TiO_2 (500°C , $\approx 2 \times 10^{-4}$ Torr H_2), showing the TiO_2 bulk along the $[11\bar{1}]$ zone axis and the Ti_3O_5 surface layer along the $[352]$ zone axis. **i**, Schematic of the TiO_2 and Ti_3O_5 structures mirroring the lattice matching in (h).

To further investigate surface structural evolutions under varying reactive gas environments, in-situ high-angle annular dark-field scanning transmission electron microscopy (HAADF-STEM) and electron energy loss spectroscopy (EELS) were performed on anatase TiO_2 NPs exposed to O_2 , H_2 , and CO_2 , respectively. Figures 2(a–c) present HAADF-STEM images under these conditions. O_2 -treated TiO_2 exhibits uniform contrast (Fig. 2(a)),

132 while H₂ exposure induces localized contrast variations (bright/dark regions, white arrow in Fig. 2(b)), signaling O
133 loss and structural evolution during reduction. Subsequent CO₂ exposure partially restores uniformity (Fig. 2(c)),
134 reflecting O replenishment and structural reconstruction with kinetic limitations (see examples in Supplementary
135 Fig. S6). EELS spectra from regions marked the dashed red boxes in Figs. 2(a-c) reveal chemical-state dynamics.
136 Under H₂, Ti-L_{2,3} and O K edges (Figs. 2d, e) show spectra shape changes (gray shaded regions), consistent with
137 Ti⁴⁺ to Ti³⁺ reduction and O depletion.²⁴ Notably, a more reduced Ti oxidation state is observed at the outermost
138 surface compared to the particle interior (see Supplementary Fig. S7). This is also supported by in-situ HRTEM
139 observations in Figs. 1(a-c) and Fig. S8, with pronounced disordering on the outermost surface representing more
140 heavily reduced, oxygen-deficient TiO_x regions. CO₂ exposure reverses these trends, restoring Ti⁴⁺ and O signatures
141 (Figs. 2d, e), aligning with structural recovery shown in Figs. 1(a-f). Surface-sensitive AP-XPS measurements (Fig.
142 2(f)) corroborate these findings: reduction in H₂ introduces electronic states within the band gap (<3.2 eV),
143 attributed to oxygen vacancy defect states,²⁵ consistent with the formation of Ti₃O₅ on the surface, while O₂ re-
144 oxidation eliminates these features. Together, these results resolve the dynamic interplay of O loss and
145 replenishment on TiO₂ surfaces, linking atomic-scale structural changes observed from in-situ TEM imaging to
146 redox-driven chemistry.

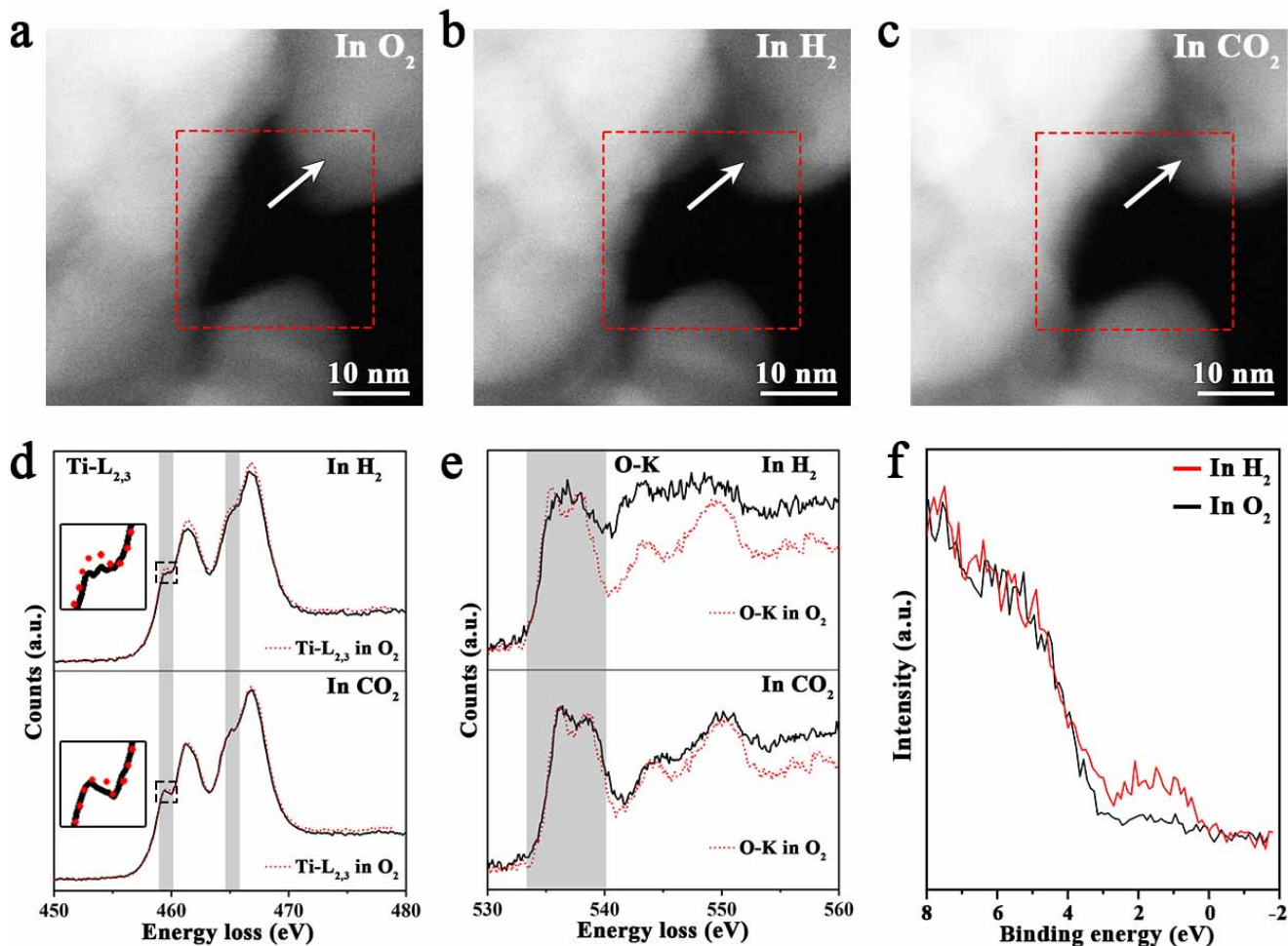


Figure 2. Redox-driven structural and chemical evolution of anatase TiO₂ surfaces under varying gas environments. **a–c**, In-situ HAADF-STEM images of TiO₂ NPs at 500 °C exposed to 5×10⁻³ Torr O₂, H₂, and CO₂, respectively. White arrows highlight localized contrast variations indicative of O loss during H₂ exposure and O replenishment under CO₂. **d, e**, EELS spectra of Ti L_{2,3} and O K edges under alternating H₂ and CO₂ exposure. Under H₂, the Ti L₃ shoulder peak attenuation and O-K edge integrating signal indicate Ti⁴⁺ → Ti³⁺ reduction and O loss. CO₂ reverses these trends, restoring Ti⁴⁺ and O-K edge splitting. **f**, AP-XPS valence band spectra of TiO₂ under O₂ (400 °C, 2×10⁻³ Torr O₂, black) and H₂ (500 °C, 5×10⁻³ Torr H₂, red).

2.2 In-situ Environmental TEM Observations of RWGS Reaction Dynamics on Anatase TiO₂ Surface

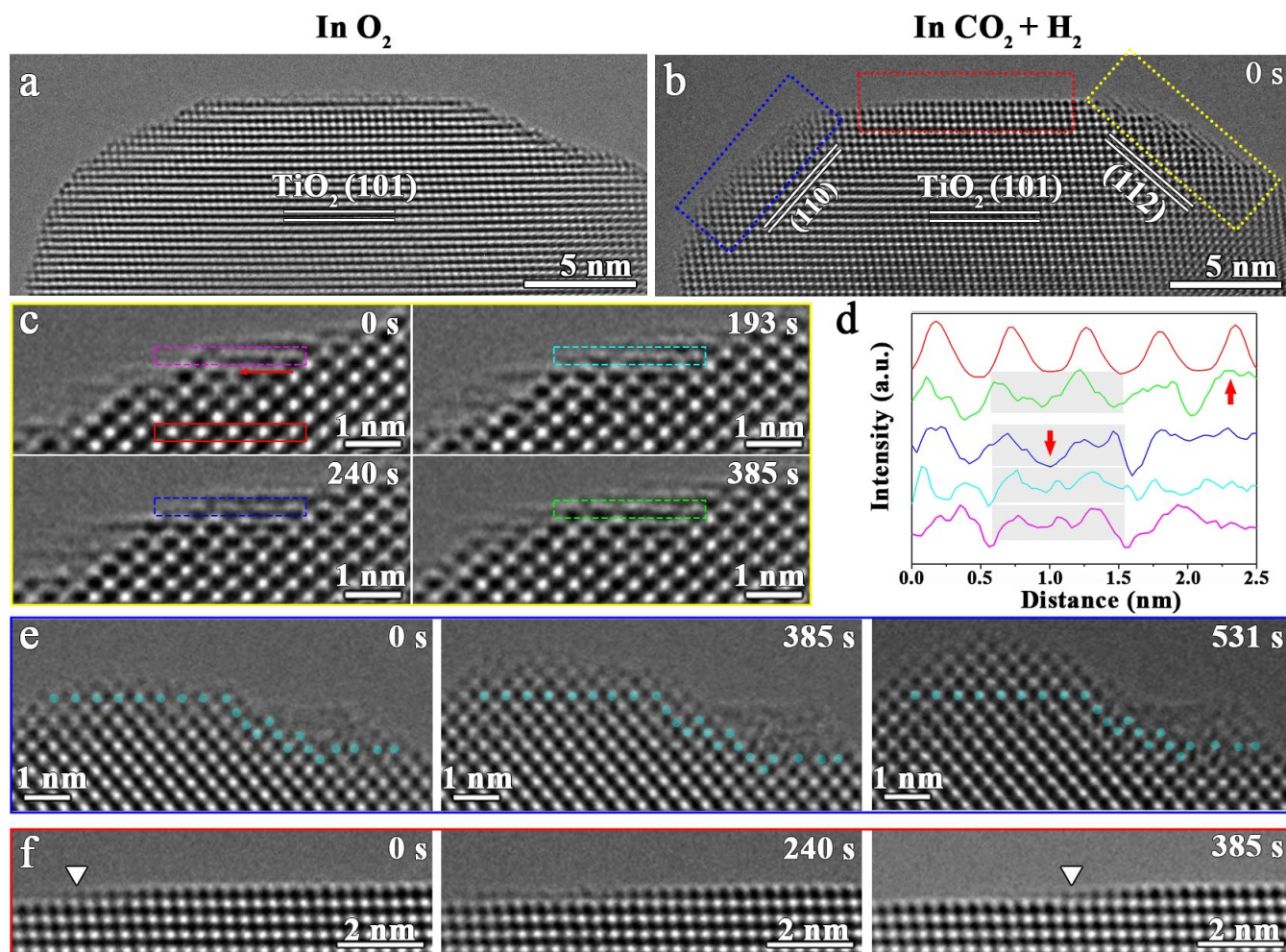
To investigate the dynamic structure behaviors, we performed in-situ TEM experiments on O₂-pretreated TiO₂ NPs (500 °C for 30 min), following a CO₂/H₂ gas mixture (2×10⁻³ Torr CO₂, 6×10⁻³ Torr H₂). Fig. 3(a)

corresponds to the initial state, revealing well-defined (101), (112), and stepped (110) facets. Notably, slight atomic column displacements observed under O₂ (2×10^{-3} Torr, 400 °C) align with prior studies,^{18–20,26–28} confirming O-adsorption-induced surface reconstruction. Time-resolved HRTEM images (Figs. 3(b-e)) capture dynamic surface evolution: (i) O loss and facet-specific activity: initial H₂/CO₂ exposure triggers O loss, inducing atomic displacements on the (110) facet and atom column merging on (112), while the (101) remains relatively stable (Fig. 3b). These changes are pronounced at undercoordinated (110) and (112) sites, prevalent at NP edges and corners. (ii) Oscillatory dynamics: progressive O exchange drives cyclic merging and splitting of atomic columns on the (112) facet (Fig. 3c), reflecting redox-driven lattice restructuring. Fig. 3d illustrates the intensity profiles of marked regions in Figs. 3c reveal atom column merging (uniform contrast) and splitting (peaks/valleys, red arrows), directly correlating with dynamic surface rearrangements.

Figure 3(e) highlights the dynamic evolution of the stepped TiO₂(110) facet marked by the dashed blue box in Fig. 3b. Cyan dots track morphological changes, revealing catalytic mass transport: new TiO₂ layers form on the left (indicative of surface growth), while the right exhibits semi-crystalline regions with disordered atomic columns display pronounced atom displacements. The appearance, disappearance, and movement of these atomic columns and pronounced displacements. These transient features arise from oscillatory O loss/gain dynamics at the gas-solid interface during CO₂ hydrogenation. Fig. 3(f) contrasts this activity with the stability of the TiO₂(101) facet. Real-time HRTEM imaging shows no structural reorganization on the (101) (more examples in Supplementary Fig. S9), indicating its limited catalytic role. Atomic step retraction initiates at the step edge (0 s image) on the TiO₂(101), followed by progressive receding (240 s image), evidenced by diminishing column contrast.

It is crucial to evaluate potential electron beam effects on in-situ TEM observations through control experiments, including "blank-beam" (Supplementary Fig. S10) and "vacuum annealing" (Supplementary Fig. S11) tests. These confirm that the electron beam dose rates used here do not measurably influence TiO₂ surface reaction dynamics (see Experimental methods for details). To complement the localized structural changes observed by in-situ TEM, synchrotron X-ray diffraction (XRD) is utilized to monitor bulk structural evolution during the RWGS reaction (Supplementary Fig. S12). XRD patterns of TiO₂—as-synthesized, O₂-pretreated, and post-reaction—confirm the stability of the anatase TiO₂. While in-situ HRTEM observations reveal dynamic surface restructuring,

185 XRD detects no bulk structural changes, underscoring that redox activity is confined to the surface. This surface-
 186 limited behavior aligns with reactivity measurements of pure anatase TiO_2 in the RWGS reaction, discussed below.



187 **Figure 3.** In-situ HRTEM visualization of RWGS reaction dynamics on TiO_2 surfaces under CO_2/H_2 gas flow. (a) HRTEM
 188 image of TiO_2 after O_2 pretreatment (4×10^{-3} Torr O_2 , 400 °C, 30 min), showing pristine (101), (112), and stepped (110) facets.
 189 (b), TiO_2 exposed to CO_2/H_2 gas mixture ($\approx 2 \times 10^{-3}$ Torr CO_2 , $\approx 6 \times 10^{-3}$ Torr H_2 , 500 °C), initiating surface restructuring. (c, e, f),
 190 Zoom-in views (Supplementary Movie 1) of the dashed regions, capturing RWGS reaction dynamics on the (112) facet (c, d),
 191 the stepped (110) facet (e), and the (101) facet (f). In (e), cyan dots mark the atom columns in the initial 0-s image, overlaid
 192 onto subsequent snapshots to track atomic rearrangements.
 193

194
 195 The competitive interplay between reduction (H_2 -driven) and oxidation (CO_2 -driven) processes on TiO_2
 196 surfaces is critically mediated by O vacancies, with CO_2 -driven O replenishment enhanced by O vacancy presence.²³

197 To further probe how O vacancies modulate RWGS dynamics, we intentionally introduce O vacancies via H₂
198 pretreatment (500°C, $\sim 6 \times 10^{-3}$ Torr H₂) and track surface evolution under RWGS conditions by adding 2×10^{-3} Torr
199 CO₂ at 500°C. Figure 4a reveals a pristine TiO₂ structure after O₂ retreatment ($\sim 4 \times 10^{-3}$ Torr O₂, 400°C). H₂
200 exposure induces surface reduction (Fig. 4b), forming a Ti₃O₅ layer on the (101) facet and amorphous regions near
201 the (110) corner due to oxygen loss. Figure 4(c) (zoomed-in view of the dashed blue box in Figure 4(b)) captures
202 oscillatory atomic dynamics on the Ti₃O₅ surface during RWGS, where CO₂ introduction partially restores the TiO₂
203 surface, as evidenced by the initial absence of significant image contrast on the topmost surface (white arrow in
204 Figure 4(b)) and the recovered lattice fringes in the 230-s HRTEM image (Fig. 4c). Unlike the stoichiometric
205 TiO₂(101) facet (Figure 3(f)), Ti₃O₅ exhibits dynamic evolution, transitioning from the outermost surface to the
206 subsurface. This is evident from regions where atomic resolution is lost and subsequently regained, as marked by
207 purple shading in Figure 4(c). Figure 4(d) provides a detailed view of the amorphous regions identified in Figure
208 4(b), revealing structural evolution and competing redox processes during the RWGS reaction, similar to untreated
209 TiO₂ in Figure 3(c). Critically, H₂-pretreated TiO₂ exhibits deeper affected zones (Fig. 4(d) vs. Fig. 3(e)),
210 highlighting how preexisting O vacancies generated from the H₂ pretreatment amplify RWGS reaction dynamics.

211 RWGS reactivity measurements in a plug flow reactor (see more details in experimental session)
212 corroborate the critical role of oxygen vacancies in the reaction mechanisms, as inferred from in-situ TEM
213 observations described above. Both pristine and H₂-pretreated anatase TiO₂ (10% H₂/Ar at 500°C for one hour)
214 exhibit 100% selectivity for CO, achieving steady-state CO₂ conversion rates of ~ 1.2 % and CO production rates
215 of $\sim 3 \times 10^{-6}$ mol/g(TiO₂)/s (Figures 4(e, f)). During thermal ramp-up and stabilization, the catalyst surface undergoes
216 structural reorganization—including increased defect density and surface restructuring—that activates additional
217 sites, while concurrent equilibration of oxygen vacancies and partial Ti reduction under reaction conditions drives
218 the system toward the quasi-steady-state structure with enhanced catalytic activity. The observed RWGS reactivity
219 and the reactivity enhancement from H₂ pretreatment, validated by the repeated measurements (Supplementary Fig.
220 S13), align well with the in situ HRTEM observations presented in Figures 3 and 4, respectively. To bridge the gap
221 between these two approaches, due to the differences in gas pressures and timescale, we have performed ex-situ
222 TEM characterizations of the TiO₂ following both H₂ reduction and CO₂/H₂ reaction treatments in the flow reactor.

223 These results (see Supplementary Fig. S14) reveal surface restructuring and morphological features indicative of a
224 more reduced state, yet chemically and structurally comparable to those observed in the ETEM experiments,
225 including partial reduction, defect formation, and phase evolution. The relatively limited forward CO production
226 rate is consistent with the in-situ XRD data (Supplementary Fig. S12) and HRTEM observations (Fig. 3),²⁹ which
227 confirm no bulk structural changes, restricting redox reactivity to surface layers.

228 The initial reactivity boost in CO₂ conversion (0.65% to 0.85%) and CO production (1.8×10^{-6} to 2.3×10^{-6}
229 mol/g(TiO₂)/s) (Figures 4(e, f), red arrows) during the first 3 h of RWGS aligns with in-situ HRTEM observations
230 (Figs. 4(a-d)). This enhancement stems from H₂ pretreatment-induced O vacancies and Ti₃O₅ formation, which
231 promote CO₂ activation and lead to a transient increase in RWGS activity. However, CO₂ gradually heals these
232 defects via selective O replenishment, driving the system toward a steady state akin to untreated TiO₂ with relatively
233 minor fluctuations (Figures 4(e, f)). The convergence to similar steady-state performance underscores that transient
234 O vacancies—not permanent structural changes—are pivotal for imparting reactivity. This dynamic equilibrium
235 reflects the competition between H₂-driven vacancy generation and CO₂-driven surface passivation (by structural
236 repair), as captured in real-time HRTEM imaging. These findings are corroborated by AP-XPS (Fig. 2f) and TPR
237 measurements (Supplementary Fig. S4), which also indicate surface oxygen depletion upon H₂ exposure.
238 Quantitative analysis of the first TPR peak suggests oxygen removal equivalent to approximately 18 % of a
239 TiO₂(101) monolayer, based on an H₂ uptake of 86 μmol/g and a surface area of 50 m²/g for anatase TiO₂. This
240 value is comparable to the ~ 8% monolayer intrinsic vacancy concentration reported for rutile TiO₂(110) surfaces.³⁰
241 Notably, the observed enhancement in CO production (~ 6 %) on H₂-pretreated TiO₂ correlates well with the extent
242 of oxygen depletion, underscoring the role of surface vacancies in modulating reactivity.

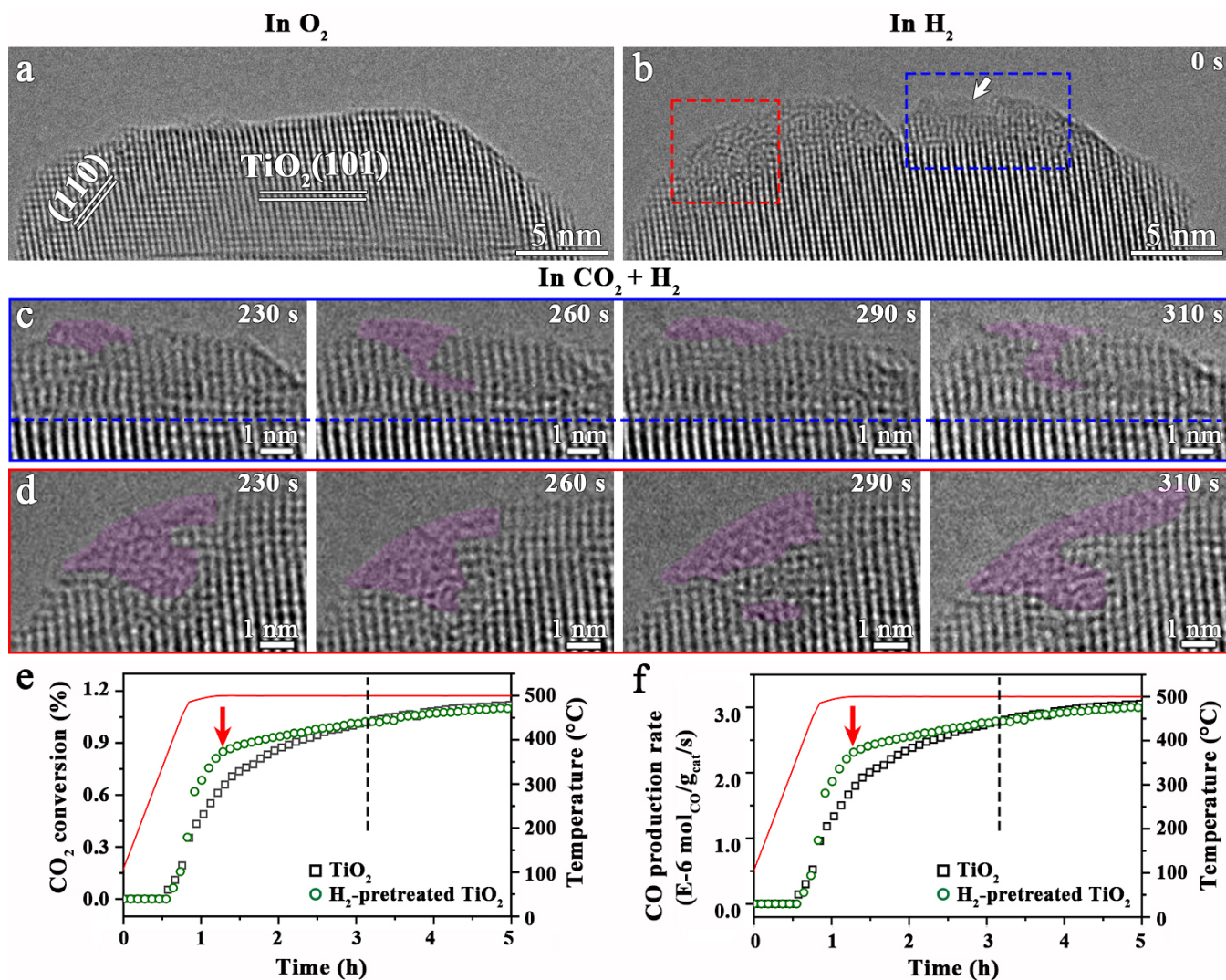


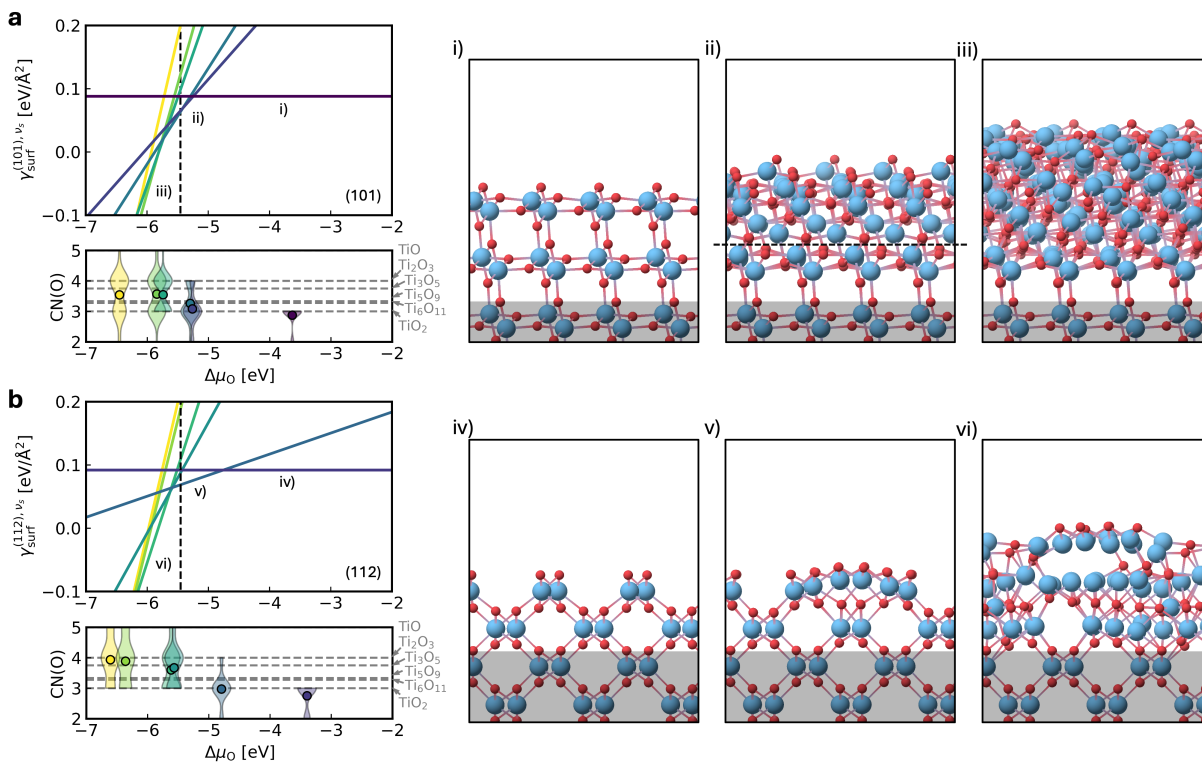
Figure 4. In-situ atomic-scale dynamics of anatase TiO₂ surfaces during H₂ pretreatment and RWGS reactions. **a, b**, HRTEM images of anatase TiO₂ under O₂ and then H₂ treatment (500 °C, $\sim 6 \times 10^{-3}$ Torr H₂), respectively. **c, d**, In-situ zoomed-in views (Supplementary Movie 2) showing the dynamic evolution of the dashed blue and red box regions (marked in b) of H₂-pretreated TiO₂ during RWGS reaction (500 °C, $\sim 2 \times 10^{-3}$ Torr CO₂, $\sim 6 \times 10^{-3}$ Torr H₂). Purple shading highlights semi-crystalline regions with diminished atomic contrast, reflecting redox-driven structural disorder. **e**, CO₂ conversion profiles for pristine TiO₂ and H₂-pretreated TiO₂ under 1:3 CO₂:H₂ flow (32 sccm), measured during a temperature ramp of 10 °C/min to 500 °C. **f**, Steady-state forward CO production rates (mass-normalized). The red arrows in (**e, f**) highlight the point exhibiting the most significant reactivity differences between the pristine TiO₂ and H₂-pretreated TiO₂ samples. The dashed black lines in (**e, f**) delineate the consistent long-term trends.

2.3 Mechanistic Insights into Catalytic Reaction Dynamics

To further elucidate the atomic origins of the reaction dynamics on anatase TiO_2 surfaces, we conducted DFT calculations and extensive structural sampling using MLIPs (see more details in Supplementary Figs. S15, S16). Building on training data from our previous work³¹, we identify three reductive surface reconstructions phases for the (101) surface, governed by oxygen chemical potential: (i) anatase retention phase ($\Delta\mu_{\text{O}} > -5.26$ eV), preserving bulk-like anatase structure, (ii) shear plane phase ($-5.74 < \Delta\mu_{\text{O}} < -5.26$ eV), featuring off-lattice shear planes, and (iii) highly reduced phase ($\Delta\mu_{\text{O}} < -5.74$ eV), dominated by disordered TiO_3 , TiO_4 , and TiO_5 . Representative atomic structures for each phase are illustrated in Figure 5. In the off-lattice reconstruction (ii), the shear plane is indicated by a dashed line, with atoms outside this boundary deviating from crystalline symmetry, while interior regions retain the crystal order. The highly reduced structure (iii) exhibits a more disordered surface, predominantly composed of TiO_3 , TiO_4 , and TiO_5 polyhedra, with a smaller fraction of TiO_2 and TiO_6 polyhedra. The O coordination number ($\text{CN}(\text{O})$) is a key descriptor of reduction and reconstruction behaviors. During TiO_2 reduction, O tends to coordinate with an increasing number of Ti atoms, leading to a higher $\text{CN}(\text{O})$. The violin plot in Figure 5 presents $\text{CN}(\text{O})$ distributions to illustrate surface reorganization. For reference, bulk anatase TiO_2 has an ideal $\text{CN}(\text{O})=3$, while the stoichiometric (101) surface exhibits a slightly lower value of 2.875 due to surface O binding to only two Ti atoms. As reduction progresses, $\text{CN}(\text{O})$ increases: the off-lattice shear plane structure has an average $\text{CN}(\text{O})=3.082$, while the highly reduced reconstruction reaches $\text{CN}(\text{O})=3.566$, approaching the $\text{CN}(\text{O})=3.750$ characteristic of bulk Ti_3O_5 , consistent with the structural evolution in Figure 1.

For the (112) surface, the stoichiometric surface (iv) features a prominent cavity between line-shaped TiO_5 units, which fills with Ti-O atoms as reduction begins ($\Delta\mu_{\text{O}} \cong -4.79$). The first reduced structure (v) stabilizes within the range $-5.56 < \Delta\mu_{\text{O}} < -4.79$ eV, forming a highly symmetric 1T structure with $P3m1$ symmetry in its outermost layer. At $\Delta\mu_{\text{O}}$ below -5.56 eV, the structure undergoes a transition into a highly disordered, amorphous-like structure, as depicted in Figure 5 (vi). Similar to the (101) surfaces, the $\text{CN}(\text{O})$ increases during the reduction process. The stoichiometric (112) surface initially has a $\text{CN}(\text{O})$ of 2.750, which rises to 2.972 in the 1T-like reduced structure (v) and further to 3.600 in the highly reduced reconstruction (vi).

280 Compared to the (101) surface, reduction on the (112) surface initiates at a significantly higher $\Delta\mu_{\text{O}}$ —
 281 approximately 0.5 eV higher. This behavior indicates that the (112) surface is more prone to reduction than the most
 282 stable (101) surface, consistent with the in-situ HRTEM observations in Fig. 3, suggesting the reactivity differences
 283 among various surface facets. It is important to note that in theoretical calculations, the oxygen chemical potential
 284 for a given surface free energy or reconstructed structure is typically treated as a constant value. However, real
 285 catalytic reactions involve both forward and reverse processes due to the presence of mixed gas environments,
 286 leading to complex surface and interfacial dynamics. As shown in Figs. 3 and 4, the evolution of surface structures
 287 includes oscillatory atomic-scale dynamics, emphasizing the interplay between competing oxidation and reduction
 288 reactions. These phenomena are further modulated by varying surface coordination conditions, which exhibit
 289 distinct reactivities toward gas-solid interfacial reactions involving adsorption and dissociation processes. The
 290 observed differences in potential transitions between pristine and reduced TiO_2 surfaces under identical reaction
 291 conditions highlight the variability in surface reactivity, driven by structural and electronic property variations.



292
 293 **Figure 5.** DFT-calculated surface free energies ($\gamma_{\text{surf}}^{(hkl),v_s}$) as a function of O chemical potential ($\Delta\mu_{\text{O}}$) for anatase **a**, (101) and
 294 **b**, (112) surfaces. The lowest-energy structures, identified through MLIP-assisted grand canonical samplings, were further

295 optimized via DFT and are represented by colored solid lines. The vertical dashed line indicates the O-poor limit ($\Delta\mu_{\text{O}} = -5.47$
 296 eV), corresponding to the formation energy of bulk anatase TiO_2 (per O atom). Violin plots below the phase diagrams illustrate
 297 the coordination numbers of O (CN(O)) in the surface region, with average CN(O) values shown as circles. Horizontal gray
 298 dashed lines denote average CN(O) values for bulk Ti-O phases. Representative atomic structures are displayed in the right
 299 panel, where Ti and O atoms are depicted as large cyan and small red spheres, respectively. The fixed bulk region is shaded in
 300 gray, distinguishing it from the surface region.

301

302 To gain deeper insight into these dynamics, we performed electronic structure analysis of the pristine and
 303 reduced $\text{TiO}_2(101)$ surfaces, focusing on the binding characteristics of H and CO_2 . Figure 6(a) presents the Bader
 304 charge distribution of Ti atoms on both surfaces. On the pristine surface, the average charge of surface Ti atoms is
 305 +2.34, closely matching the bulk Ti charge of +2.49, with a small standard deviation of 0.025. Despite the presence
 306 of five-fold-coordinated surface Ti sites, their oxidation state remains largely unaltered. However, on the reduced
 307 surface, Ti atoms undergo a significant reduction in charge, with an average Ti charge of +1.55 and a much larger
 308 standard deviation of 0.355, indicating substantial local charge redistribution upon surface reduction.

309 This phenomenon is clearly reflected in the spin-resolved partial density-of-states (PDOS) shown in Figure
 310 6(b). The pristine structure exhibits a well-defined band gap with no defect states, leading to a symmetric occupation
 311 of spin components. In contrast, the reduced surface shows the emergence of significant mid-gap states that are
 312 spin-polarized and predominantly localized in the up-spin channel. This spin anisotropy in the filled mid-gap defect
 313 states, primarily derived from Ti $3d$ orbitals, indicates the formation of Ti^{3+} polarons at the surface. On the other
 314 hand, mid-gap contributions from oxygen $2p$ states are nearly negligible, and the occupied O $2p$ states below the
 315 Fermi level remain relatively symmetric between spin channels. This suggests that hole localization on oxygen sites
 316 ($\text{O}^\cdot$ polarons) is not significant, as no unoccupied spin-polarized states characteristic of hole polarons are observed
 317 near the valence band edge. The high concentration of Ti^{3+} polarons at the reduced surface implies distinct
 318 adsorption characteristics for adsorbates compared to the pristine surface.

319 We then computed the binding energies of H and CO_2 on both pristine and reduced $\text{TiO}_2(101)$ surfaces. To
 320 systematically explore all possible adsorption sites, we employed a graph-based approach to generate symmetry-

inequivalent, unique initial adsorption configurations, ensuring a comprehensive search for the most favorable binding sites and energies. Interestingly, both H and CO₂ exhibit stronger binding on the reduced TiO₂(101), with binding energies of -5.581 and -6.671 eV, respectively. In contrast, their adsorption on the pristine surface is much weaker, with binding energies of -0.243 and -0.736 eV, respectively. More specifically, H preferentially binds at the two-fold bridge oxygen site on the pristine surface but relocates to the Ti-Ti bridge site on the reduced surface. Meanwhile, the CO₂ molecule is weakly physisorbed at the five-fold Ti site on the pristine surface, with a bond length of 2.47 Å. In contrast, on the reduced surface, CO₂ adopts a bent configuration, with both oxygen atoms binding to adjacent Ti atoms at atop sites. This bent geometry suggests activation of CO₂, facilitating its dissociation and enhancing its surface reactivity.³²

To elucidate the nature of adsorption interactions, we performed Crystal Orbital Hamilton Population (COHP) analysis for H and CO₂ on both pristine and reduced surfaces. This analysis provides insight into the bonding characteristics by identifying bonding and antibonding states relative to the Fermi energy. As discussed above, H binds at the bridge O site on the pristine surface, with its bonding state located at approximately -8 eV below the Fermi level. Notably, no bonding states appear near the Fermi energy (-5 to 0 eV), indicating that the interaction is deeply localized within the valence band rather than contributing to surface reactivity. The integrated COHP (iCOHP) values for spin-up and spin-down components are highly negative (-3.969 and -3.968, respectively), suggesting a strongly covalent bonding interaction. However, despite this strong local bond formation, the overall binding energy remains weak (-0.243 eV), indicating that H adsorption is not well stabilized on the pristine surface. This discrepancy suggests that while H forms a local covalent bond with O, other factors, such as lack of charge transfer, prevent strong overall adsorption.

In contrast, CO₂ binds weakly at the Ti atop site, with no distinct bonding or antibonding states near the Fermi level. The very small iCOHP values (-0.446 for both spin-up and spin-down) indicate that the interaction is neither covalent nor ionic but instead dominated by weak van der Waals forces, consistent with its physisorption behavior. On the reduced surface, distinct changes in the electronic structure emerge. For H adsorption, a clear bonding state appears closer to the Fermi level, around -4 eV, suggesting increased interaction with the surface. However, the iCOHP values (-1.546 for spin-up and -1.621 for spin-down) are less negative than those on the

pristine surface, despite the much stronger binding energy (-5.581 eV). This suggests that H binding on the reduced surface has a more ionic character, likely influenced by charge redistribution and stronger electrostatic interactions with reduced Ti sites.

For CO₂ adsorption, a strong bonding state emerges at -2 eV, significantly different from the weak physisorption observed on the pristine surface. The highly negative iCOHP values (-4.033 for spin-up and -4.162 for spin-down) indicate that CO₂ interacts via a strong covalent bond rather than simple electrostatic attraction, consistent with its bent adsorption configuration. These findings reveal a fundamental shift in adsorption behavior upon surface reduction. While the pristine surface exhibits weak overall hydrogen binding despite strong localized H-O covalent interactions, the reduced surface promotes significantly stronger, more ionic hydrogen adsorption and enhanced covalent CO₂ binding. This change is crucial for CO₂ activation and dissociation, providing insight into the reduction and reoxidation processes occurring under different reaction conditions. The notably stronger binding of H on the reduced surface suggests that once an initial reduction occurs, the surface becomes highly favorable for further hydrogen adsorption. This autocatalytic reduction process explains why TiO₂ readily reduces in pure H₂ environments—once some oxygen vacancies form through the oxidation of adsorbed hydrogen, they serve as high-affinity sites for further H adsorption, accelerating continued reduction.³³ In contrast, pristine TiO₂ does not strongly bind H, making reduction a slow process until a sufficient number of vacancies accumulate to trigger autocatalysis.

Under RWGS conditions, exposure to CO₂ leads to the exceptionally strong adsorption of CO₂ on the reduced TiO₂ surface. This occurs due to the high concentration of O vacancies and Ti³⁺ centers, which enhance electron transfer to CO₂, promoting its dissociation and supplying oxygen atoms to refill vacancies. This process effectively re-oxidizes Ti³⁺ back to Ti⁴⁺, restoring the pristine TiO₂ phase. Since reoxidation eliminates high-affinity sites for H adsorption, the system undergoes a self-sustaining transition from reduced to pristine TiO₂, preventing further reduction despite the continued presence of H₂. These binding energy trends align well with experimental observations. Under pure H₂ conditions, the weak initial interaction of H with pristine TiO₂ allows reduction to proceed gradually, eventually leading to a fully reduced TiO₂ surface. However, under RWGS conditions, the strong affinity of CO₂ for reduced TiO₂ drives its dissociation, leading to reoxidation and the restoration of a pristine-like phase, as shown in the reactivity trends in Figure 4(e, f). This interplay between adsorption-driven reduction and

reoxidation creates a dynamic surface state, where TiO_2 continuously cycles between reduced and oxidized phases depending on the prevailing reaction environment.

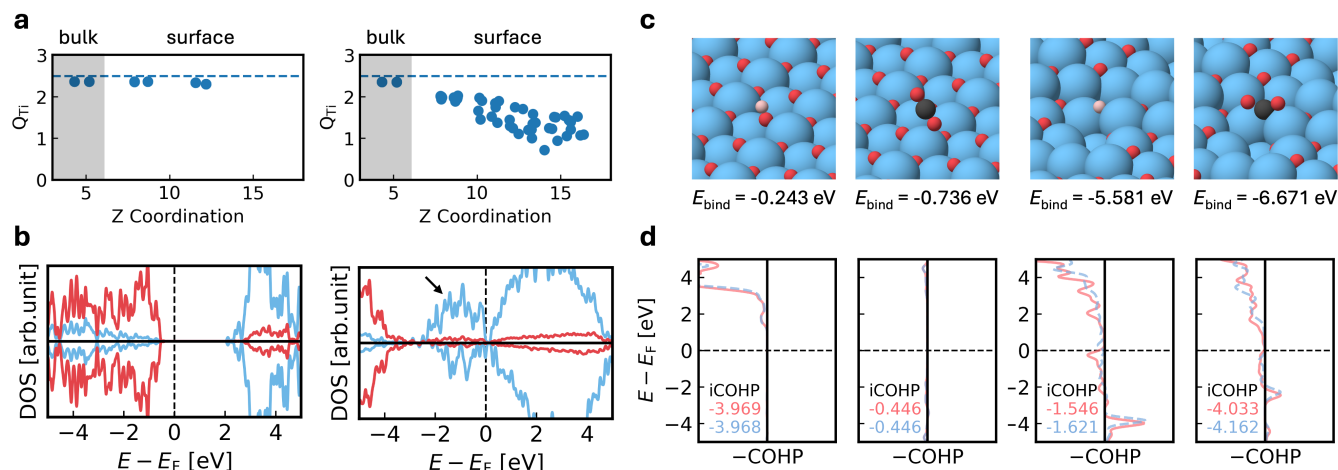


Figure 6. Electronic structure analyses for pristine and reduced $\text{TiO}_2(101)$ surfaces. **a**, Bader charge (Q_{Ti}) of Ti atoms, plotted against their z-coordinates in the supercell for the pristine (left) and reduced (right) surfaces. The horizontal dashed line indicates the reference Bader charge of Ti in bulk anatase TiO_2 . **b**, Spin-resolved partial density-of-states (PDOS) for the pristine (left) and reduced (right) surfaces. Ti 3d and O 2p states are shown in blue and red lines, respectively. The arrow indicates the defect states associated with Ti^{3+} within the band gap of the reduced surface. **c**, Adsorption configurations of H and CO_2 on pristine (left) and reduced (right) $\text{TiO}_2(101)$, with binding energies (E_{bind}) noted below each panel. Ti, O, C, and H atoms are represented by blue, red, black, and white spheres, respectively. **d**, Crystal Orbital Hamilton Population (-COHP) analysis for the bonding interactions between each adsorbate and its neighboring surface atoms. Spin-up and spin-down components are represented as red solid and blue dashed lines, respectively. Positive peaks indicate bonding states, whereas negative values correspond to antibonding interactions. Integrated COHP (iCOHP) values are provided as insets.

The results presented above resolve the atomic-scale redox mechanisms governing the RWGS reaction on anatase TiO_2 , establishing O vacancy dynamics and surface orientation as pivotal factors. The facet-dependent activity arises from the preferential reduction of undercoordinated sites on high-index (112) and stepped surfaces, where H_2 generates O vacancies and Ti^{3+} -rich surface layers while the dominant (101) facet remains inert. The apparent discrepancy in dynamic behaviors of the (101) facet arises from differences in reaction environments: under H_2 alone (Figs. 1 and 4), the (101) surface undergoes reduction via lattice oxygen loss, forming Ti_3O_5 due to

393 favorable thermodynamics. In contrast, during $\text{CO}_2 + \text{H}_2$ catalysis (Fig. 3), competing reduction and reoxidation
394 processes dominate. CO_2 acts as a dynamic oxidant, replenishing lattice oxygen and stabilizing partially reduced
395 regions—particularly on the more stable (101) facet, as supported by DFT (Fig. 5). Consequently, while the (101)
396 facet is reducible under strongly reducing conditions, it remains relatively stable during RWGS, whereas the more
397 reducible (112) surface undergoes greater restructuring driven by redox cycling.

398 CO_2 rapidly heals the O vacancies—threefold faster than H_2 -induced reduction—restoring lattice O and
399 suppressing sustained catalytic activity. This competition creates oscillatory atomic-scale dynamics but limits
400 overall reactivity due to the dominance of inert (101) facets and CO_2 -driven passivation of active sites.
401 Mechanistically, these findings define anatase TiO_2 as a dynamic equilibrium between defect creation and
402 annihilation, governed by local coordination environments. Practically, pure anatase TiO_2 is generally not an effec
403 tive catalyst for the RWGS reaction due to its limited ability to activate CO_2 and dissociate H_2 (Figs. 4(e, f)). How
404 ever, when used as an oxide support for dispersed metal species—such as single atoms or nanoparticles— TiO_2 ca
405 n enable significantly enhanced catalytic performance.^{6,7,21,22,34} These metal–support systems exhibit strong interfa
406 cial interactions that promote key steps in the RWGS reaction, including hydrogen activation and CO_2 reduction. I
407 mportantly, structure–reactivity studies have revealed that the catalytic activity and stability are highly sensitive to
408 the crystallographic facets of the TiO_2 support.^{21,22,34,35} Specific surface orientations can stabilize different metal c
409 onfigurations or defect structures, which in turn modulate the reaction pathway. The work highlights surface
410 engineering as a critical lever for optimization: selectively exposing reactive facets (e.g., (112), (110)) or stabilizing
411 O vacancies (via doping, strain, or heterostructuring) could amplify RWGS efficiency of TiO_2 support. Such
412 strategies could transform TiO_2 into a tunable platform for heterogeneous catalyst development for CO_2 conversion,
413 bridging atomic-scale defect control to industrial-scale sustainability goals like carbon-neutral fuel synthesis. By
414 linking surface geometry, defect lifetimes, and reaction kinetics, this work provides a framework for designing
415 adaptive oxide tailored for dynamic gas-solid interfaces.

Conclusions

This study provides direct evidence of the dynamic interplay of O vacancy generation and replenishment in governing the RWGS reaction on anatase TiO₂ surfaces. Through in-situ atomic-scale imaging and computational modeling, we demonstrate that RWGS proceeds via redox cycle mediated by surface orientation and defect dynamics: (i) reduction-driven O loss: H₂ preferentially reduces undercoordinated sites on high-indexed (112) and stepped surfaces, forming O vacancies and Ti³⁺-rich domains, while the thermodynamically stable (101) facet remains structurally intact and inert. (ii) oxidation-driven O replenishment: CO₂ dissociates at O vacancy sites, restoring lattice O with kinetics approximately three times faster than reduction. This rapid healing passivates active sites, leading to transient activity loss. The competition between these processes drives oscillatory atomic-scale dynamics, yet results in modest overall reactivity due to two key factors: (i) CO₂-induced healing of O vacancies, which suppresses sustained catalytic activity, and (ii) the dominance of inert (101) facets on anatase TiO₂ NPs. These results underscore the critical role of atomic-scale defect control and facet engineering in tailoring oxide catalysts for sustainable CO₂ conversion technologies.

429 **Experimental methods**

430 **In-situ environmental TEM experiments.** The anatase TiO₂ nanoparticles (99.9% anatase, average
431 particle diameter ~ 30 nm) were purchased from Research Nanomaterials Inc. and subsequently loaded onto
432 through-hole SiN_x thin films on DENSsolutions chips. In-situ TEM experiments were performed using a single-tilt
433 DENSsolutions heating holder within an environmental TEM (FEI Titan 80-300) equipped with an objective-lens
434 aberration corrector and a differential pumping system. The annealing processes were conducted at a heating rate
435 of 100 °C/min. EELS measurements were performed in the ETEM using probe scanning mode following in-situ
436 treatments under various gas environments. To thoroughly evaluate the effects of the electron beam, comparative
437 experiments were performed to assess RWGS reaction dynamics and EELS spectrum acquisition process of TiO₂
438 under electron beam blanked conditions (Fig. S10). Additionally, we investigated the effects of vacuum annealing
439 and continuous electron beam irradiation on TiO₂ reduction and found that these effects were less pronounced than
440 those observed in H₂ (Fig. S11).

441 **AP-XPS measurements.** AP-XPS experiments were conducted at the Center for Functional
442 Nanomaterials, Brookhaven National Laboratory. The system utilized a SPECS Phoibos 100 MCD analyzer for
443 XPS measurements, coupled with a monochromatic Al K α X-ray source (1486.6 eV, ~ 0.25 eV line width), focused
444 to a spot size smaller than 300 μ m. To minimize charging effects in the AP-XPS measurements, the conductivity of
445 the sample was improved by pressing a small amount of anatase TiO₂ powders onto a Tantalum (Ta) foil. The Ta
446 foil was subsequently spot welded directly onto the sample holder, which was heated using a laser heater positioned
447 at the backside of the sample holder. The AP-XPS setup, equipped with multiple differential pumping apertures,
448 enables gas pressures in the Torr range near the sample while maintaining ultrahigh vacuum (UHV) conditions in
449 the hemispherical analyzer. All spectra were acquired at a 20° takeoff angle relative to the electron analyzer optics.

450 **In-situ XRD measurements.** X-ray diffraction experiment was performed at beamline BL2-1 at Stanford
451 Synchrotron Radiation Lightsource at SLAC National Accelerator Laboratory using X-ray energy of 17 keV
452 (0.72932 Å). 2.7 mg of pressed and sieved anatase TiO₂ is packed in a 1 mm capillary on a flow cell reactor. The
453 sample was first pretreated in 20% O₂ in He for a half hour. After cooling and purging the cell with helium, the

sample was then exposed to a reaction mixture (24: 8 sccm H₂: CO₂) and heated from 200 °C to 700 °C, with a 30-minute dwell at every 50 °C increment, then from 700 °C to 720 °C with a 20-minute dwell at every 10 °C increment.

Physical adsorption measurements. The surface area of the as-received anatase TiO₂ was measured by physical adsorption at 77 K with nitrogen as the adsorbent gas on a Micromeritics 3Flex. The quantity of N₂ adsorption over relative pressure (P/P₀) range of 0.05–0.35 was used to calculate the BET surface area. A carbon black surface area reference material (Micromeritics part # 004-16833-00) with a reported multi-point surface area of 21.5±0.8 m²/g was run simultaneously to ensure the accuracy of the measured surface area of the as-received anatase TiO₂ support.

RWGS reactivity measurements. Reactivity measurements were conducted in a packed-bed flow reactor using a 3/8-inch outer diameter stainless steel tube coated with an inert Silcolloy® coating. The reactor tube was heated within a copper block furnace using 5-inch resistive cartridge heaters, and the temperature was monitored by a K-type thermocouple positioned in the catalyst bed. Anatase TiO₂ samples (20 mg, US research Nanomaterials Inc, 30 nm) were diluted with SiO₂ (500 mg, Sigma Aldrich, 40–150 mesh, 90–425 µm particle size) and thoroughly mixed using a vortex mixer to prevent the formation of hotspots. The diluted catalyst mixture was loaded into the inert stainless steel reactor tube and positioned between two plugs of quartz wool. The catalyst was pretreated in 100% O₂ at 400°C for 30 minutes, with a ramp rate of 10°C/min, and then cooled to room temperature under an O₂ atmosphere. For the H₂-pretreated TiO₂ catalysts, an additional treatment with 10% H₂/Ar at 500°C for one hour was performed after the O₂ pretreatment. The reactor was then purged with Ar to remove residual pretreated gas. A reactant mixture (32 sccm with a 1: 3 CO₂: H₂ ratio) was introduced at room temperature and ramped to 500 °C over one 1 hour, where it was maintained for 15 hours to conduct the reaction study. The product gases were analyzed using gas chromatograph (SRI 8610C) with a methanizer-equipped flame ionizing detector and thermal conductivity detector.

Density functional theory calculations. Spin-polarized density functional theory (DFT) calculations were carried out using the projector-augmented wave (PAW) method³⁶ within the Vienna Ab initio Simulation Package (VASP).^{37,38} The valence states explicitly included the 3*p*, 3*d*, and 4*s* orbitals of Ti, along with the 2*s* and 2*p* orbitals

of O. The electronic exchange and correlation were modeled using the semi-local generalized gradient approximation (GGA) with the Perdew-Burke-Ernzerhof (PBE) functional,³⁹ enhanced by van der Waals (vdW) corrections following the Tkatchenko-Scheffler (TS) scheme.⁴⁰ To accurately describe the electronic structure of TiO₂, an on-site Coulomb interaction correction (DFT+*U*) was applied to the Ti 3*d* electrons using the Dudarev approach,⁴¹ with the effective Hubbard parameter of 3.5 eV.^{42,43} The calculations employed a plane-wave basis set with a kinetic energy cutoff of 500 eV. Brillouin zone sampling was performed with a reciprocal distance of 0.03 Å⁻¹. The lattice parameters were optimized by iteratively minimizing lattice stress and atomic displacements until the external pressure fell below 0.5 kbar. Convergence criteria were defined as total energy changes below 10⁻⁵ eV and 10⁻⁴ eV for electronic and ionic relaxations, respectively. Under these conditions, the atomic forces in the optimized bulk TiO₂ structure were minimized to less than 0.01 eV/Å. The aforementioned DFT setups are consistent with Ref [26] to ensure the continuous use of the data for further training of the machine-learning potentials.

Surface calculations utilized supercell geometries with periodic boundary conditions, incorporating a vacuum gap of at least 20 Å to eliminate unintended interactions between repeating images. For the TiO₂(101) (TiO₂(112)) surface, asymmetric slab models were designed using (3×1) ((2×1)) surface unit cells, comprising a minimum of 24 (12) atomic layers, with the bottom 12 (6) layers fixed at their bulk positions. The energies used in the DFT surface phase diagrams were computed with a dipole correction applied perpendicular to the surface to counteract artificial electrostatic effects. However, no dipole correction was included in the DFT results used for training the machine-learning potentials to prevent inaccuracies arising from artificial energy modifications.

MACE architecture. The MLIP model was implemented using the MACE architecture, an equivariant message-passing graph tensor network designed to encode many-body atomic geometry information at each layer.⁴⁴ The initial training set for the TiO₂ system was obtained from the database provided in Ref [26]. This training set includes rutile and anatase TiO₂ surfaces mostly for rutile (110) and anatase (101). In the MACE model used herein, two message-passing layers were utilized, with spherical expansions up to $l_{\max} = 3$, and 4-body messages were incorporated in each layer, corresponding to a correlation order of 3. All models were constructed with a 128-channel dimension for tensor decomposition. A radial cutoff of 6 Å was applied, and the interatomic distances were

expanded into 8 Bessel functions, which were further multiplied by a smooth polynomial cutoff function to generate the radial features. The model adopted maximal message equivariance with an order of $L = 1$. For better accuracy and expandability, fine-tuning the medium materials project⁴⁵ foundation model (MACE-MP-0⁴⁶) was used.

Grand canonical genetic algorithm. To investigate the chemical space of anatase TiO₂ surfaces under reductive conditions, we employed a global optimization strategy using the grand canonical genetic algorithm (GCGA), which was implemented in our proprietary GOCIA package.⁴⁷ In this framework, the system is treated as a grand canonical ensemble where Ti and O adsorbates interact with the anatase slab, aiming to minimize the coverage-dependent grand canonical free energy Ω , which is expressed as:

$$\Omega = U - TS - \sum_i n_i \mu_i = G_{\text{surf}}^{(110),v_s} - G_{\text{surf}}^{(110),v_0} - \Delta n_{\text{Ti}} \mu_{\text{Ti}} - \Delta n_{\text{O}} \mu_{\text{O}},$$

Here, $G_{\text{surf}}^{(110),v_s}$ and $G_{\text{surf}}^{(110),v_0}$ correspond to the Gibbs free energies of the newly generated (termination v_s) and reference (termination v_0) TiO₂ slabs, respectively. The variables n_i and μ_i denote the number of atoms of species i in the periodic system and their respective chemical potentials. The chemical potentials of Ti and O (μ_{Ti} and μ_{O}) are constrained by the equilibrium with bulk anatase TiO₂, satisfying the relation: $g_{\text{TiO}_2,\text{bulk}} = \mu_{\text{Ti}} + 2\mu_{\text{O}}$. Additionally, μ_{O} is determined by the equilibrium with an oxygen gas-phase reservoir: $\mu_{\text{O}} = \frac{1}{2} E_{\text{O}_2} + \Delta\mu_{\text{O}}$, where E_{O_2} represents the total energy of an isolated O₂ molecule, incorporating zero-point energy (ZPE) corrections.⁴⁸

For GCGA sampling, a population of 30 structures and a mutation rate of 40% were chosen. The initial set of candidates was generated using the bond length distribution algorithm (BLDA), which randomly constructs structures based on covalent radii.⁴⁹ A pre-optimization step utilizing Hookean potentials was applied to ensure physically reasonable initial geometries before local optimization with MLIP and energy evaluation. The generation of offspring followed a split-and-splice approach,⁵⁰ in which parent slabs were sectioned along random planes and then recombined.

Each candidate was assigned a fitness score based on mating frequency and grand canonical free energy, prioritizing structures with lower free energy and infrequent mating for further evolution. A similarity check was implemented to prevent duplicate structures from being added to the population. Mutation operations included the

addition or removal of adsorbates, random displacement of surface atoms, translation of the buffer slab along the x- or y-axis by $1/n$ (where $n = 2, 3, 6$) of the cell length, and permutation of half of the buffer region. If offspring exhibited excessive similarity to their parents, their mutation rate was increased to 100%. To maintain a constant population size, the candidate with the lowest fitness was achieved after each addition.

Ab initio thermodynamics. To evaluate the stability of different reconstructed terminations of the anatase TiO_2 surface, we employed an ab initio thermodynamics approach.⁵¹ Similar to the grand canonical free energy approach discussed earlier, the surfaces were assumed to be in thermodynamic equilibrium with an oxygen-containing gas phase. This allowed for the computation of the surface free energy, $\gamma_{\text{surf}}^{(hkl),\nu_s}$, as follows:

$$\gamma_{\text{surf}}^{(hkl),\nu_s} = \frac{1}{A^{(hkl)}} \left[G_{\text{surf}}^{(hkl),\nu_s} - \sum_i n_i \mu_i \right]$$

Here, the Gibbs free energies of the condensed phases were approximated using DFT-derived total energies,⁵² enabling the equation to be rewritten as:

$$\gamma_{\text{surf}}^{(hkl),\nu_s}(\Delta\mu_{\text{O}}) = \frac{1}{A^{(hkl)}} \left[E_{\text{surf}}^{(hkl),\nu_s} - n_{\text{Ti}} E_{\text{TiO}_2,\text{bulk}} - (n_{\text{O}} - 2n_{\text{Ti}}) \left(\frac{1}{2} E_{\text{O}_2} + \Delta\mu_{\text{O}} \right) \right].$$

This formulation provides a direct means of assessing the thermodynamic stability of various surface configurations under different environmental conditions.

543 **SUPPORTING INFORMATION**

544 Supporting information is available free of charge on the ACS publication website:

545 Supplementary figures, and Supplementary in-situ TEM movies

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549 Notes

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551

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