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THE DIFFUSION LAYER ON A ROTATING DISK ELECTRODE

John Newman

November 1966

The Diffusion Layer on a Rotating Disk Electrode

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November, 1966

In a recently published work,¹ the author treated the current distribution on a rotating disk under conditions where ohmic potential drop, surface overpotential, and concentration variations near the electrode must all be taken into account. The diffusion layer, which is then characterized by a nonuniform current density and concentration at the surface, was handled by a power-series expansion in r, the radial distance from the axis of rotation:

$$c = c_{\infty} \left[1 + \sum_{m=0}^{\infty} A_m (r/r_0)^{2m} \theta_m(\zeta) \right], \qquad (1)$$

where c is the concentration of the reactant, c_{∞} is the bulk concentration, r_{o} is the radius of the electrode, and θ_{m} represents functions of $\zeta = y (a\nu/3D)^{1/3}\sqrt{\Omega/\nu}$, where y is the normal distance from the disk, a = 0.51023, v is the kinematic viscosity, D is the diffusion coefficient, and Ω is the rotation speed of the disk.

The functions θ_{m} and, in particular, the derivatives $\theta'_{m}(0)$ at the surface were obtained by a finite-difference solution of the differential equations

$$\theta_{\rm m}^{\rm u} + 3\zeta^2 \theta_{\rm m}^{\rm i} - 6m\zeta \theta_{\rm m}^{\rm i} = 0, \qquad (2)$$

subject to the boundary conditions

$$\theta_{\rm m}$$
 = 1 at ζ = 0 and $\theta_{\rm m}$ = 0 at ζ = ∞ .

Subsequent numerical difficulties were encountered which required accurate

values of the derivatives $\theta_{\perp}^{\prime}(0)$.

As suggested by the work of Rosner,² one can use the equations in section 58 of Levich's book³ to obtain the concentration at the surface in terms of an integral over the concentration derivative at the surface:

$$c(\mathbf{r},\mathbf{o}) = c_{\mathbf{o}}(\mathbf{r}) = \frac{1}{\Gamma(2/3)} \int_{\mathbf{o}}^{\mathbf{r}} \frac{J(\mathbf{r}')\mathbf{r}' d\mathbf{r}'}{(\mathbf{r}^{3}-\mathbf{r}'^{3})^{2/3}},$$
 (3)

where $J(r) = -\partial c/\partial \zeta$ at $\zeta = 0$. Substitution of a power-series expansion for c into equation (3) yields an expression for $\theta_m(0)$:

$$\theta_{m}^{*}(0) = -\frac{3\Gamma(\frac{5}{3})\Gamma(\frac{2m+3}{3})}{2\Gamma(\frac{4}{3})\Gamma(\frac{2m+2}{3})} \quad . \tag{4}$$

The gamma functions necessary to evaluate $\theta_m^{\prime}(0)$ from equation (4) are given⁴ to 15 significant figures. The values of $\theta_m^{\prime}(0)$ are in satisfactory agreement with those tabulated in table 1 of reference 1. At the same time direct use of equation (3) or its inverse² provides a possibility of avoiding the numerical difficulties encountered in the earlier work.¹

1. John Newman. "Current Distribution on a Rotating Disk below the Limiting Current." Journal of the Electrochemical Society, 113, 1235-1241 (1966).

2. Daniel E. Rosner. "Reaction Rates on Partially Blocked Rotating Disks--Effect of Chemical Kinetic Limitations." Journal of the Electrochemical Society, 113, 624-625 (1966).

3. Veniamin G. Levich. <u>Physicochemical Hydrodynamics</u>. Englewood Cliffs, N. J.: Prentice-Hall, Inc., 1962.

4. Milton Abramowitz and Irene A: Stegun, eds. <u>Handbook of Mathe-</u> <u>matical Functions</u>. Washington: National Bureau of Standards, 1964. p.3.

Acknowledgment

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