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## **Composition Dependence of the Flory-Huggins Interaction**

# Parameters of Block Copolymer Electrolytes and the Isotaksis Point

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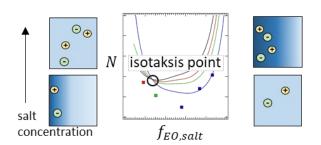
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#### ABSTRACT

The thermodynamics of block copolymer/salt mixtures were quantified through the application of Leibler's Random Phase Approximation to disordered small angle X-ray scattering profiles. The experimental system comprises of polystyrene-block-poly(ethylene oxide) (SEO) mixed with lithium bis(trifluoromethanesulfonyl) imide salt (LiTFSI), SEO/LiTFSI. The determined Flory-Huggins interaction parameter from scattering experiments,  $\chi_{SC}$ , was found to be a function of block copolymer composition, chain length, and temperature for both salt-free and salty systems. In the absence of salt,  $\chi_{0,SC}$  is a linear function of  $(Nf_{EO})^{-1}$ ; in the presence of salt, a linear approximation is used to describe the effect of salt on  $\chi_{eff,SC}$  for a given copolymer composition and chain length. The theory of Sanchez was used to determine  $\chi_{eff}$  from  $\chi_{eff,SC}$  in order to predict the boundary between order and disorder as a function of chain length, block copolymer composition, salt concentration, and temperature. At fixed temperature (100 °C),  $N_{crit}$ , the chain length of SEO at the order-disorder transition in SEO/LiTFSI mixtures, was predicted as a function of the volume fraction of the salt-containing poly(ethylene oxide)-rich microphase,  $f_{EO,salt}$ , and salt concentration. At  $f_{\rm EO, salt}$  >0.27, the addition of salt stabilizes the ordered phase; at  $f_{EO,salt}$  < 0.27, the addition of salt stabilizes the disordered phase. We propose a simple theoretical model to predict the block copolymer

composition at which phase behavior is independent of salt concentration (  $f_{EO,salt}$ =0.27). We refer to this composition as the "isotaksis point".

## INTRODUCTION

There is continued interest in understanding the thermodynamics of polymer/salt mixtures due to their applications as solid electrolytes in rechargeable batteries.<sup>1-7</sup> It is well known that the addition of salt to diblock copolymers greatly affects their thermodynamics and there have been many theoretical and experimental studies on quantifying these effects. A model system is polystyrene-*block*-poly(ethylene oxide) (SEO) mixed with lithium bis(trifluoromethanesulfonyl) imide salt (LiTFSI), SEO/LiTFSI. The transport of Li<sup>+</sup> ions in poly(ethylene oxide) (PEO) has been fully characterized<sup>8</sup> and the addition of polystyrene (PS) introduces mechanical support for the electrolyte.<sup>9</sup>

The phase behavior of pure diblock copolymers has been thoroughly studied, both experimentally and theoretically.<sup>10,11</sup> The equilibrium phase behavior is dictated by two parameters: the volume fraction of one polymer block,  $f_A$ , and the segregation strength,  $\chi N$ , where N is the overall degree of polymerization and  $\chi$  is the Flory-Huggins interaction parameter, a measurement of the thermodynamic compatibility between the two polymer constituents. The temperature (*T*) dependence of  $\chi$  is often given by

$$\chi = \frac{A}{T} + B (1)$$

where *A* and *B* are empirically determined constants.<sup>10,12</sup> At high temperatures, entropy dominates leading to the formation of a disordered phase. As temperature decreases, the importance of the energy of

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interactions between the two polymer blocks increases, and leads to microphase separation and the formation of ordered morphologies. For a given block copolymer, characterized by  $f_A$  and N, the transition from disorder-to-order occurs at a critical value of  $\chi$ . Self-consistent field theory (SCFT) is a powerful tool that can predict the phase behavior of neat block copolymers.<sup>13</sup> At the order-disorder transition (ODT), the product  $\chi N$  may be given by

$$(\chi N)_{odt} = g(f_A) = C_0 + C_1 (f_A - 0.5)^2 + C_2 (f_A - 0.5)^4 + C_3 (f_A - 0.5)^6 + C_4 (f_A - 0.5)^8$$
(2)

where the coefficients,  $C_i$ , are obtained by fitting Equation 2 to the SCFT results of Cochran et al.<sup>11</sup>:  $C_0=10.5$ ,  $C_1=47.9$ ,  $C_2=782$ ,  $C_3=-3567$  and  $C_4=24700$ . In this theory,  $\chi$  is an implicit property of the chemical structure of the constituent monomers in the block copolymer and does not depend on N or  $f_A$ . For a given block copolymer system at a particular temperature, T, Equation 2 can be used to calculate a critical chain length,  $N_{crit}$ , as a function of  $f_A$ :

$$N_{crit} = \frac{g(f_A)}{\chi(T)}$$
 (3)

Block copolymers of a given composition,  $f_A$ , will be ordered if  $N \ge N_{crit}$ . A common way to measure  $\chi$  is through the application of Leibler's Random Phase Approximation (RPA)<sup>14</sup>, where small angle X-ray (SAXS) experiments can be used to measure concentration fluctuations by fitting the structure

factor, S(q), to disordered scattering profiles. The interaction parameters derived from scattering are called  $\chi_{sc}$ . There has been considerable debate about how the measured  $\chi_{sc}$  relates to the value of  $\chi$  that should be used in the SCFT calculations.<sup>15-20</sup> An attractive feature of  $N_{crit}$  is that it can be measured directly and interpreted without any debate.

The addition of salt is known to alter the thermodynamics of block copolymers due to the introduction of new interactions between the polymer chains and ions, e.g., electrostatic interactions, charge dissociation, ion solvation, ion translational entropy and physical cross-linking between the ions and polymer chains.<sup>21-29</sup> Ions tends to segregate in the phase with higher permittivity, which increases segregation strength between the salt-free and the salt-containing blocks. This was captured in models developed by Wang and coworkers using the concept of Born solvation energy.<sup>30-33</sup> In these models,  $\chi$  is replaced with an effective interaction parameter,  $\chi_{eff}$ , to account for the interactions introduced by salt. In the simplest case

## $\chi_{eff} = \chi_0 + mr$ (4)

where  $\chi_0$  is the Flory-Huggins parameter for the salt-free system, r is a suitable measure of salt concentration and m is a system-dependent proportionality constant. This form for  $\chi_{eff}$  in salty systems was first proposed in the pioneering work of Mayes and coworkers.<sup>34</sup>

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In this paper, which builds on our previous study of the phase behavior of block copolymer electrolytes<sup>35</sup>, we use the standard RPA-based analysis to determine the effective interaction parameter,  $\chi_{eff,SC}$ , from SAXS profiles of a series of disordered SEO/LiTFSI mixtures. Our experiments cover a wide range of block copolymer compositions,  $0.18 \le f_{EO} \le 0.84$ , and chain lengths,  $49 \le N \le 414$ . We also determine  $N_{crit}$  as a function of block copolymer composition. It is generally observed that adding salt stabilizes the ordered phase, i.e., *m* in Equation 4 is positive. This would imply that  $N_{crit}$  must decrease with added salt. We show that this is only true over a finite range of copolymer compositions,  $0.27 \le f_{EO} \le 0.90$ . In the range  $0.15 \le f_{EO} < 0.27$ ,  $N_{crit}$ increases with added salt and *m* is negative. We find that *m* is a smooth function of  $f_{EO}$ . For reasons that we clarify below, we propose using the term isotaksis composition to refer to the point where m = 0. For SEO/LiTFSI systems, the isotaksis composition is  $f_{EO} = 0.27$ .

#### **EXPERIMENTAL SYSTEMS**

#### Polymer Synthesis and Characterization

The polystyrene-*block*-poly(ethylene oxide) (SEO) copolymers in this study were synthesized via anionic living polymerization<sup>36</sup> and purified according to ref 37. The copolymers used in this study are called SEO(*xx*-*yy*), where *xx* and *yy* are the number-averaged molecular weights of the PS,

 $M_{PS}$ , and PEO,  $M_{PEO}$ , in kg mol<sup>-1</sup>, respectively. Chain length, N, was calculated by  $N = N_{PS} + N_{PEO}$  where

$$N_i = \frac{M_i}{\rho_i N_A v_{ref}}$$
 (*i* = PS or PEO) (5)

and  $N_A$  is Avogadro's number and  $v_{ref}$  was fixed at 0.1 nm<sup>3</sup>. The volume fractions of each block of the copolymers were calculated by

$$f_{EO} = \frac{v_{EO}}{v_{EO} + \frac{M_{PS}M_{EO}}{M_{S}M_{PEO}}} v_{S}}$$
(6)

where  $v_{EO}$  and  $v_S$  are the molar volumes of ethylene oxide and styrene monomer units, respectively, and  $M_{EO}$  and  $M_S$  are the molar masses of ethylene oxide (44.05 g mol<sup>-1</sup>) and styrene (104.15 g mol<sup>-1</sup>), respectively. Molar volumes were calculated by  $v=M/\rho$ . In this study, the densities (g cm<sup>-3</sup>) of the PEO and PS blocks were given by  $\rho_{PEO}=1.13$  and  $\rho_{PS}=1.05$ , measured values at 90 °C. The neat copolymers are completely transparent and colorless. Table 1 gives the properties of the 8 SEO copolymers used in this study.

polymer	MP	MP	N	f <sub>EO</sub>
	S	EO		
SEO(17.4- 3.9)	17. 4	3.9	342	0.18
SEO(9.4-	9.4	2.4	189	0.20

Table 1: Properties of copolymers in the study.

2.4)				
SEO(9.4- 4.0)	9.4	4.0	214	0.29
SEO(1.7- 1.4)	1.7	1.4	49	0.44
SEO(2.9- 3.3)	2.9	3.3	99	0.52
SEO(3.8- 8.2)	3.8	8.2	189	0.68
SEO(5.1- 12.8)	5.1	12.8	281	0.71
SEO(4.0- 22.4)	4.0	22.4	414	0.84

### Electrolyte Preparation

The block copolymer electrolytes used in this study were prepared according to ref 38. The copolymers were dried at 90 °C under vacuum in a glovebox antechamber for at least 12 hours and then immediately brought into an argon environment. Lithium bis(trifluoromethanesulfonyl)imide (LiTFSI) salt (Novolyte) was transferred from its air-free packaging into a vial inside of a glovebox, and then dried at 120 °C under vacuum in a glovebox antechamber for three days. Due to the hydroscopic nature of the salt, Argon environment gloveboxes (Vacuum Atmosphere Company) with low oxygen and water levels were used for all sample preparation. The salt containing samples were prepared by blending SEO/benzene solutions with the required amount of a 75 wt% solution of LiTFSI/tetrahydrofuran (THF) solution to achieve the calculated salt concentrations. For the salty samples, we assume that all of the salt resides in the PEO domain.<sup>39-41</sup> Block copolymers containing salt are considered to be pseudobinary systems where the volume fraction of the salt + PEO component is given by

$$f_{EO,salt} = \frac{v_{EO,LiTFSI}(r)}{v_{EO,LiTFSI}(r) + \left(\frac{M_{PS}M_{EO}}{M_{S}M_{PEO}}v_{S}\right)}$$
(7)

where *r* is the molar ratio of Li to ethylene oxide (EO) moieties ( $r = \frac{[Li]}{[EO]}$ ) and

 $\nu_{{\scriptscriptstyle EO, {\scriptscriptstyle LiTFSI}}}$  is the molar volume of the salt-containing PEO phase calculated by

$$v_{EO, LiTFSI}(r) = \frac{M_{EO} + r M_{LiTFSI}}{\rho_{EO, LiTFSI}}$$
(8)

where

$$\rho_{EQ,LiTESI}(r) = 2.008r + 1.13$$
 (9)

is derived from measured density values at 90 °C taken from ref 8.

## Small Angle X-Ray Scattering (SAXS) Measurements

SAXS samples were prepared by pressing/melting the polymer into a 1/16 or 1/32 in. thick annular Viton rubber spacer (McMaster Carr) with an inner diameter of 1/8 in. at 120 °C in an argon glovebox and annealing at 120 °C for at least 12 hours under vacuum. The samples were cooled to room temperature under vacuum over the course of 24 hours. The samples were sealed with Kapton windows in custom-designed airtight Aluminum sample

holders. The samples were mounted in a custom-built 8-sample stage and annealed at each temperature for 20 minutes before performing measurements. SAXS measurements were conducted at the Advanced Light Source beamline 7.3.3 at Lawrence Berkeley National Lab<sup>42</sup> and Stanford Synchrotron Radiation Light Source beamline 1-5 at SLAC National Accelerator Laboratory. Scattering was performed using 10-12 keV X-rays.

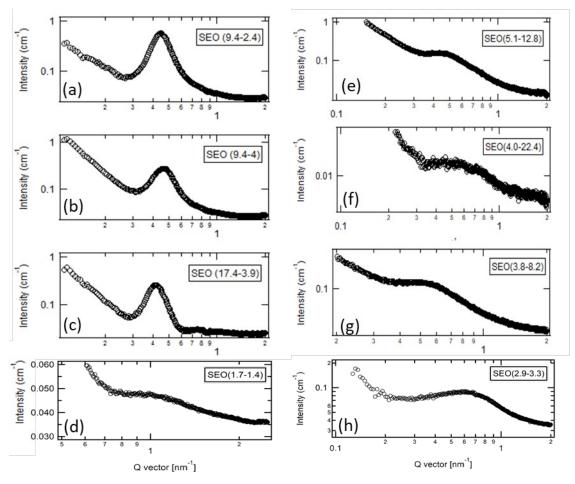
Silver behenate was used to determine the beam center and sampleto-detector distance. The scattered intensity was corrected for beam transmission, empty cell scattering, as well as for unavoidable air gaps in the system. Glassy carbon (NIST) was used to determine the scaling calibration to obtain absolute intensity scattering. Two-dimensional scattering patterns were integrated azimuthally using the Nika program for IGOR Pro to produce one-dimensional (1D) scattering profiles.<sup>43</sup> In order to compare data collected at each beamline, temperature calibrations were conducted to measure the absolute temperature of the samples by making separate electrolyte samples with a thermocouple running through the sample holder. The data presented in the main text reflects the absolute temperatures of the samples.

### **EXPERIMENTAL RESULTS**

Figure 1 shows the measured absolute scattering intensity (I) as a function of scattering vector, q, of the eight SEO copolymers in the salt-free

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state at 90 °C. All eight copolymers are disordered, indicated by a single broad scattering peak, at all accessible temperatures. As the molecular weight of the copolymer increases, the location of the primary scattering peak,  $q^{i}$ , moves to a lower value of q, indicative of an increase in chain dimensions. However, the overall intensity of the scattering is not proportional to *N*. For example, SEO(4.0-22.4) has the lowest scattering intensity, but is the longest copolymer in the study.



**Figure 1**: Absolute scattering for asymmetric copolymers in the neat state (r = 0) at 90 °C for (a) SEO(9.4-2.4), (b) SEO(9.4-4.0), (c) SEO(17.4-3.9), (d) SEO(1.7-1.4), (e) SEO(5.1-12.8), (f) SEO(4.0-22.4), (g) SEO(3.8-8.2), and (h) SEO(2.9-3.3).

The scattering theory of monodisperse disordered diblock copolymers was developed by Leibler.<sup>14</sup> The scattering function I(q) proposed by this theory for a perfectly monodisperse AB diblock copolymer with degree of polymerization *N* can be written as

$$I_{dis}(q) = C \left[ \frac{S(q)}{W(q)} - 2\chi \right]^{-1}$$
 (10)

where C is the electron contrast calculated by

$$C = v_{ref} (B_A - B_B)^2 (11)$$

where  $v_{ref}$  is the reference volume,  $B_i$  is the X-ray scattering length density of

block *i* given by  $B_i = \frac{b_i}{v_i}$ , and  $v_i$  and  $b_i$  are the monomer volumes and X-ray scattering lengths of block *i*, respectively; W(q) and S(q) are the determinant and sum of the elements of the structure factor matrix  $||S_{ij}||$ . The expressions for W(q) and S(q) are given by

$$W(q) = S_{AA} \circ S_{BB} \circ - (S_{AB} \circ)^2$$
 (12)

 $S(q) = S_{AA}^{\circ} + S_{BB}^{\circ} + 2S_{AB}^{\circ}$  (13)

where

 $S_{ii}^{\circ} = f_i N_i v_i P_i(q)$  (i=A,B) (14)

$$S_{AB}^{\circ} = S_{BA}^{\circ} = (N_A f_A N_B f_B)^{\frac{1}{2}} F_A(q) F_B(q)$$
 (15)

and

$$P_{i}(\boldsymbol{q}) = 2 \left[ \frac{\exp(-\boldsymbol{x}_{i}) - 1 + \boldsymbol{x}_{i}}{\boldsymbol{x}_{i}^{2}} \right] (16)$$

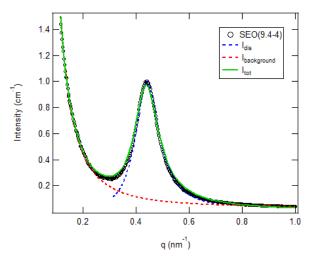
$$F_i = \frac{1 - \exp(-x_i)}{x_i}$$
(17)

with  $x_i = q^2 R_{g,i}^2$ . Both blocks are modeled as flexible Gaussian chains and

$$R_{g,i}^2 = \frac{N_i (\alpha a_i)^2}{6} (i = PS, PEO) (18)$$

where  $a_i$  is the statistical segment length of block *i*. In order to account for the conformational asymmetry between PS and PEO, we set  $a_{PS} = 0.50$  nm

and  $a_{PEO} = 0.72$  nm (conformational asymmetry parameter,  $\epsilon = \frac{a_{EO}}{a_s} = 1.44$ ).<sup>44-46</sup> In Equation 18, a chain stretching parameter,  $\alpha$ , is introduced to match experimental and theoretical values of  $R_{g,i}$ .  $N_i$  in Equations 14-18 is the number-average degree of polymerization for block *i* based on  $\nu_{ref}$ , and the calculated values are provided in Table 1. Equations 10-18 are used to analyze the scattering profiles from disordered block copolymer/salt mixtures. We ignore the fact that these equations were only developed for pure disordered block copolymers.<sup>37,47-49</sup>



**Figure 2**: Example of RPA fit on neat SEO(9.4-4.0) at 75 °C. The open circles show the raw data and the green, blue, and red dashed curves show the fits for the total scattering, RPA fit, and background correction.

Figure 2 shows a typical SAXS profile obtained from salt-free SEO(9.4-4.0) at 75 °C. The open symbols show the data and the solid green line represents a fit to the equation

$$I_{tot}(q) = I_{dis}(q) + I_{bkgrd}(q)$$
 (19)

where  $I_{dis}(q)$  (shown in blue) is Equation 10 with  $\chi$ ,  $\alpha$  and C as adjustable parameters and  $I_{bkgrd}(q)$  (shown in red) is an exponential function to compensate for imperfect background subtraction. We find excellent agreement between the fitted function and the data.

We begin with a discussion of the fitted parameter, *C*. If we assume that the salt molecules remain strongly correlated with the EO segments in the disordered state, then we can consider our mixtures to comprise of two "components": PS and PEO + salt. The density of homopolymer PEO/LiTFSI mixtures,  $\rho_{EO,salt}$ , as a function of r has been measured experimentally and is provided in Equation 9.<sup>8</sup> This can be used to calculate the theoretical scattering length density of these mixtures,  $B_{EO,salt}$ . If we assume that  $B_s$  is given by the known value obtained from homopolymer PS, then the only unknown in Equation 11 is  $B_{EO,salt}$ . The measurements of C thus provide a measurement of  $B_{EO,salt}$ , which we refer to as  $B_{EO,salt}^{fit}$ . This parameter can be used to calculate the effective density of the PEO/LiTFSI component,  $\rho_{EO,salt}^{fit}$  using the following equations

$$\rho_{EO,salt}^{fit} = B_{EO,salt}^{fit} \left( \frac{\rho_{EO,salt}}{B_{EO,salt}} \right)$$
(20)

and

$$B_{EO,salt} = Y_{LiTFSI} B_{LiTFSI} + (1 - Y_{LiTFSI}) B_{EO}$$
(21)

where  $Y_{LiTFSI}$  is the volume fraction of LiTFSI in the PEO+salt phase calculated by

$$Y_{LiTFSI} = \frac{r v_{LiTFSI}}{(1+r) v_{EO, LiTFSI}}$$
(22)

where  $v_{LiTFSI}$  was calculated from  $\rho_{LiTFSI}$ =2.023 g cm<sup>-3</sup>. These calculations were performed at 90 °C, the applicable temperature for Equation 9.

Figure 3a shows the averaged values of  $\rho_{EO,salt}^{fit}$  for all eight copolymers as a function of salt concentration at 90 °C. The error bars represent the standard deviation of the averaged data set. First, it is important to note that the individual values for  $\rho_{EO,salt}^{fit}$  for all SEO copolymers collapse on to a single point for each salt concentration. This provides justification for our assumption that SEO/LiTFSI mixtures can be approximated as twocomponent systems. In other words, the change in density of PEO from homopolymer values is dependent only on salt concentration and not on copolymer properties, *N* and  $f_{EO}$ . The fitted density of PEO/LiTFSI is a weak function of salt concentration, decreasing by < 1% over the experimental salt concentration window (Figure 3a). This decrease is opposite to what has been observed in homopolymer PEO/LiTFSI mixtures.<sup>8</sup> Figure 3b is a plot of

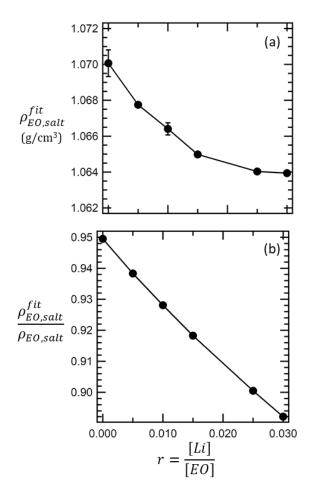
$$rac{
ho_{EO,salt}^{ht}}{
ho_{EO,salt}}$$
 versus salt concentration. The fitted density values are consistently

lower than the homopolymer values. It is important to note that even in the

absence of salt, r=0,  $\frac{\rho_{EO,salt}^{fit}}{\rho_{EO,salt}}=0.95$  indicating that the density of PEO is affected by the presence of the PS block. The deviations between the fitted density and that of homopolymer/salt mixtures is less than 10% over the entire salt concentration window. These seemingly insignificant deviations in  $\rho_{EO,salt}$  are important during the calculation of the scattering contrast because the electron density of the two phases (PS and PEO/LiTFSI) are

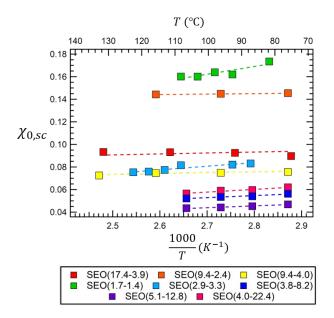
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similar: a 7% decrease in  $\rho_{EO,salt}$  results in a 100% increase in  $(B_{EO,salt}-B_s)^2$ when r=0.01.



**Figure 3**: Results from fitting contrast during RPA fits: averaged values for (a) calculated salty PEO density,  $\rho_{EO,salt}^{fit}$ , and (b) ratio of fitted density to actual density,  $\frac{\rho_{EO,salt}^{fit}}{r}$ , as a function of salt concentration, *r*, taken at 90 °C. Error bars show the

standard deviations for the data sets. Lines are used to guide the eye.

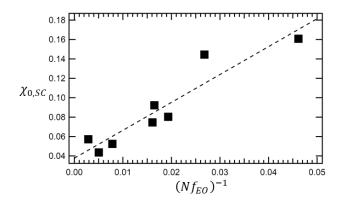


**Figure 4**: Temperature dependence for the Flory-Huggins interaction parameter for the neat SEO copolymers,  $\chi_{0,SC}$ . Dashed lines are fits to Equation 1.

The main parameter of interest is  $\chi_{SC}$ , obtained by the fitting procedure depicted in Figure 2. Figure 4 shows the temperature dependence of  $\chi_{0,SC}$ , the Flory-Huggins interaction parameter determined from scattering for the salt-free block copolymers. The dashed lines are fits of the extracted  $\chi_{0,SC}$  values presented in Figure 4 to Equation 1. The temperature dependence of  $\chi_{0,SC}$  is consistent throughout the copolymers studied:  $\chi_{0,SC}$ decreases with increasing temperature. However, the values of *A* and *B* obtained vary significantly between the copolymers. In other words,  $\chi_{0,SC}$ dependence of  $\chi_{0,SC}$  on composition and chain length. A simple function that is consistent with our data is shown in Figure 5 where we plot  $\chi_{0,SC}$  obtained at 100 °C versus  $(f_{EO}N)^{-1}$ . When data was not taken at exactly 100 °C, the fits to Equation 1 were used to interpolate  $\chi_{0,SC}$  to 100 °C. The squares represent the data and the dashed line is a linear regression fit through the data,

$$\chi_{0,SC} = K_1 + \frac{K_2}{Nf_{EO}}$$
 (23)

The regression analysis gives  $K_1$ =0.038 and  $K_2$ =2.85.



**Figure 5**: Chain length, *N*, and composition,  $f_{EO}$ , dependence on the Flory-Huggins interaction parameter of the neat copolymers,  $\chi_{0,SC}$ , at 100 °C. The dashed line is a linear regression through the data:  $\chi_{0,SC} = \frac{2.85}{Nf_{EO}} + 0.038$ .

The functional form given in Equation 23 is motivated by the theory of Fredrickson and Helfand who theoretically examined the effect of concentration fluctuations on the phase behavior of block copolymers.<sup>50</sup> In this theory, the product  $\chi N$  at the ODT,  $(\chi N)_{ODT}$ , decreases with increasing N.

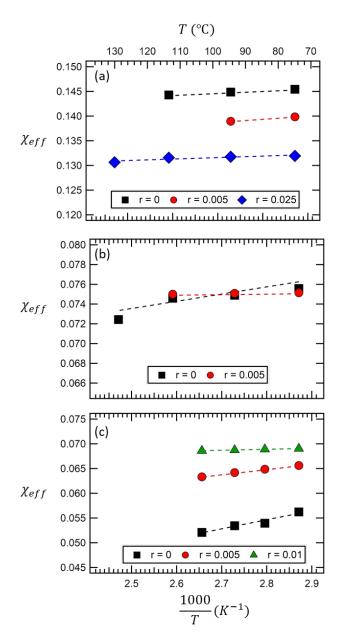
For example, in symmetric systems with  $f_A = 0.50$ ,  $(\chi N)_{ODT} = 10.495 + 41.002 N^{\frac{-1}{3}}$ 

. The mean-field value of  $(\chi N)_{ODT}$  is only obtained in the limit of infinite chain length.<sup>14</sup> Equation 23 is similar in spirit and for finite  $f_{EO}$ , a composition- and chain-length-independent  $\chi$  is obtained in the limit of infinite chain length. The product  $Nf_{EO}$  is equal to  $N_{EO}$ , the degree of polymerization of the PEO block. Equation 23 implies an implicit asymmetry in the thermodynamics between PS and PEO: a longer PEO block reduces the thermodynamic incompatibility between the two polymer blocks.

We now move to a discussion of the thermodynamics of the saltcontaining SEO copolymers. Figure 6 shows the temperature dependence of  $\chi_{_{eff,SC}}$  , the interaction parameter of the salt-containing species derived from scattering, for three SEO copolymers of similar chain lengths (N 200) with varying PEO compositions. Solid squares represent experimental measurements and the dashed lines are a fit through the data according to Equation 1. The temperature dependence of  $\chi_{\rm eff,SC}$  matches what was seen in the salt-free copolymers; A in Equation 1 remains positive. Figure 6a shows the temperature dependence of  $\chi_{eff,SC}$  of SEO(9.4-2.4) with  $f_{EO}$ =0.20. As salt is added to the system,  $\chi_{eff,SC}$  steadily decreases from 0.145 at r=0to 0.140 at r = 0.005 and finally to 0.13 at r = 0.025 at 93 °C. In most cases reported in literature, increases with increasing the  $\chi_{eff.SC}$ salt concentration.<sup>51-55</sup> Although the behavior of SEO(9.4-2.4) is unexpected, this trend has been previously reported in ref 37 for a SEO copolymer of a similar composition (SEO(1.9-0.8) with N=42 and  $f_{EO}=0.29$ ). Figure 6b shows

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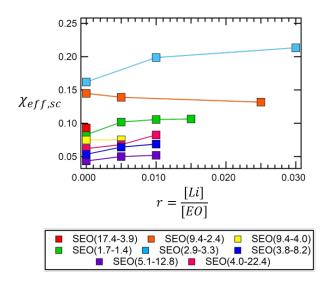
the temperature dependence of  $\chi_{eff,SC}$  of SEO(9.4-4.0) with  $f_{EO}$ =0.29. Here, we do not see a significant change in  $\chi_{eff,SC}$  upon salt addition to r=0.005 indicating that at this copolymer composition,  $\chi_{eff,SC}$  is not a strong function of r. Figure 6c shows the temperature dependence of  $\chi_{eff,SC}$  of SEO(3.8-8.2) with  $f_{EO}$ =0.67. This copolymer shows a significant increase in  $\chi_{eff,SC}$  upon salt addition (almost a 30% increase, from 0.50 at r=0 to 0.62 at r=0.005 at 100 °C). Figure 6 shows that the effect of salt addition on  $\chi_{eff,SC}$  is highly dependent on copolymer composition.



**Figure 6**: Temperature dependence of effective Flory-Huggins interaction parameter,  $\chi_{eff,SC}$ , for neat and salty samples of (a) SEO(9.4-2.4), (b) SEO(9.4-4.0), and (c) SEO(3.8-8.2). All copolymers have chain lengths around N=200. Symbols represent the data and dashed lines are fits to Equation 1.

Figure 7 shows the salt concentration dependence of  $\chi_{eff,SC}$  for the eight SEO copolymers taken at 100 °C. Solid lines are used to connect the data points and serve to guide the eye. In general,  $\chi_{eff,SC}$  increases with r,

except for SEO(9.4-2.4) as described above (shown in orange). For the low molecular weight SEO copolymers, SEO(1.7-1.4) (cyan) and SEO(2.9-3.3) (green), the effect of salt on  $\chi_{eff,SC}$  seems to level-off at a given salt concentration as reported in ref 37, while for the copolymer with the highest PEO composition, SEO(4.0-22.4) (pink),  $\chi_{eff,SC}$  increases linearly with salt concentration over the observed salt concentration window. It is obvious from Figure 7 that the effect of salt on  $\chi_{eff,SC}$  is dependent on both *N* and  $f_{EO}$ .



**Figure 7**: Salt concentration dependence for the effective Flory-Huggins interaction parameter,  $\chi_{eff,SC}$ , at 100 °C for the SEO/LiTFSI mixtures in the study. Lines are drawn to connect data points as a guide for the eye.

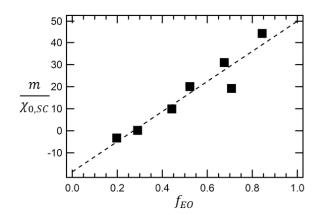
As a first approximation, each of the datasets presented in Figure 7 were fit to Equation 4, where m is a copolymer dependent proportionality constant. Figure 8 shows the composition dependence of m taken at 100 °C

weighted by the interaction parameters of the salt-free systems,  $\chi_{0,SC}$ . The dashed line is a linear regression through the data according to

$$\frac{m}{\chi_{0,SC}} = K_3 f_{EO} + K_4$$
 (24)

where  $K_3$ =68.3 and  $K_4$ =-18.6. Note, only seven data points are presented in Figure 8 because there is no salt containing data for SEO(17.4-3.9), which orders immediately upon salt addition ( $r \ge 0.005$ ). Equation 24 quantifies the dependence of rate of change in  $\chi_{eff,SC}$  upon salt addition, given by m, on copolymer composition. We see excellent agreement between the data in

Figure 8 and Equation 24. As the volume fraction of PEO increases,  $\frac{m}{\chi_{0,sc}}$  increases. It is important to note that at the lowest value of  $f_{EO}$ , the ordinate of Figure 8 becomes negative, characterized by  $K_4$ .



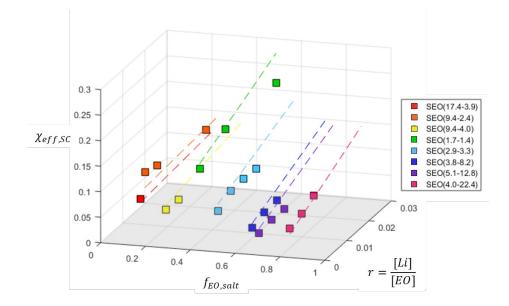
**Figure 8**: Composition,  $f_{EO}$ , dependence on the weighted slope of Equation 4,  $\frac{m}{\chi_{0,SC}}$ 

, taken at 100 °C. Dashed line is a fit to  $\frac{m}{\chi_{0,SC}}$ =68.2 $f_{EO}$ -18.6.

Combining Equations 4, 23 and 24, we arrive at

$$\chi_{eff,SC} = \left[ K_1 + \frac{K_2}{Nf_{EO,salt}} \right] \left[ 1 + K_3 f_{EO,salt} r + K_4 r \right]$$
(25)

The comparison between Equation 25 and the data are presented in Figure 9 on a three-dimensional plot where  $\chi_{eff,SC}$  is shown as a function of  $f_{EO,salt}$  and r. The solid squares represent the extracted  $\chi_{eff,SC}$  replicated from Figure 7, and the dashed lines represent fits to Equation 25 in the range of  $0 \le r \le 0.03$ . Overall, we see good agreement between the measured values of  $\chi_{eff,SC}$  and Equation 25.



**Figure 9**: Model predictions for effective Flory-Huggins interaction,  $\chi_{eff,SC}$ , as a function of salt concentration taken at 100 °C. Dashed lines are the fits to Equation 25,  $\chi_{eff,SC} = \left[ K_1 + \frac{K_2}{Nf_{EO,salt}} \right] [1 + K_3 f_{EO,salt} r + K_4 r]$ . Data points are replicated from Figure 7. The slopes of the lines are negative for the block copolymers with  $f_{EO,salt} < 0.25$ .

The final parameter extracted from the RPA fits is the chain-stretching parameter,  $\alpha$ . In most cases,  $\alpha > 1$  implying that chains are stretched. In the neat copolymers, as temperature increases,  $\alpha$  decreases (Figure SI1). At fixed tempeature,  $\alpha$  for neat copolymers decreases with increasing  $f_{EO}$ . Figure SI2 shows data at 100 °C. At the highest  $f_{EO}$ ,  $\alpha$  is less than one. The addition of salt generally leads to an increase in  $\alpha$  as shown in Figure SI3. The dependence of  $\alpha$  on r, shown in Figure SI3, is very similar to the dependence of  $\chi_{eff,SC}$  on r shown in Figure 7. In general, the increase in

effective repulsion between the blocks, quantified by  $\chi_{eff,SC}$ , leads to a larger value of  $\alpha$ .

#### **ORDER-DISORDER TRANSITION**

The discussion thus far has focused on scattering from disordered SEO/ LiTFSI mixtures. Whether or not a particular mixture is ordered depends on four variables: *N*,  $f_{EO,salt}$ , *r*, and *T*. Our discussion below is restricted to a fixed temperature of 100 °C. At fixed values of  $f_{EO,salt}$  and *r*, one can, in principle, access a transition from disorder to order by increasing *N*. We define  $N_{crit}$  as the chain length at that transition. If  $\chi$  is known, then Equation 3 can be used to determine  $N_{crit}$ .

As a first approximation, we assume that the Flory-Huggins interaction parameter extracted from scattering is equivalent to the theoretical parameter,  $\chi_{eff,SC} = \chi$ , and then use Equation 25 in conjunction with Equation 2 to calculate  $N_{crit}$ :

$$(\chi N)_{odt} = \left[ \kappa_1 N_{crit} + \frac{\kappa_2}{f_{EO, salt}} \right] \left[ 1 + \kappa_3 f_{EO, salt} r + \kappa_4 r \right] = g(f_{EO, salt}) (26)$$

Equation 26 can be solved to obtain  $N_{crit}$  as a function of  $f_{EO, salt}$  and r.

The solid curve in Figure 10a shows the calculated values of  $N_{crit}$  as a function of  $f_{EO}$  in the neat state, r=0. The model prediction is asymmetric with respect to copolymer composition, unlike predictions from mean-field

theory for conventional block copolymer systems<sup>14</sup>, with a minimum at  $f_{EO}$ =0.37. The  $N_{crit}$  versus  $f_{EO}$  curve may be considered a phase boundary: systems below the curve are predicted to be disordered while those above the curve are predicted to be ordered. The squares in Figure 10a represent the SEO copolymers considered in this study, characterized by N and  $f_{EO}$ . All of the SEO copolymers are disordered in the neat state; however, several of the squares lie above the  $N_{crit}(f_{EO})$  curve in Figure 10a. Therefore, it is evident that our assumption that  $\chi_{SC}=\chi$  leads to an inconsistency in the model.

There are no published results for relating  $\chi_{sc}$  to  $\chi$  in block copolymers. However, for the case of polymer blends, Sanchez<sup>15</sup> presented a simple expression relating these two parameters by recognizing that :

$$\chi_{sc} = \frac{-1}{2} \left( \frac{\partial^2 (f_A (1 - f_A) \chi)}{\partial f_A^2} \right)$$
(27)

The solution for Equation 27, which first proposed by de Gennes<sup>16</sup>, with appropriate boundary conditions is

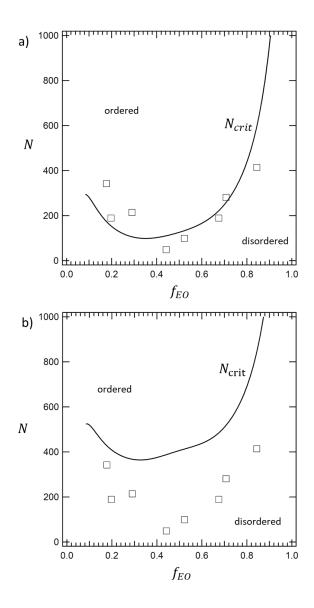
$$\chi = \frac{2}{1 - f_{A}} \int_{0}^{1 - f_{A}} \left( 1 - f_{A}^{'} \right) \chi_{SC} \left( f_{A}^{'} \right) d \left( 1 - f_{A}^{'} \right) + \frac{2}{f_{A}} \int_{0}^{f_{A}} \left( f_{A}^{'} \right) \chi_{SC} \left( f_{A}^{'} \right) d f_{A}^{'}$$
(28)

When Equation 25 for  $\chi_{eff,SC}(f_{EO,salt})$  is substituted into Equation 28 for  $\chi_{SC}(f_A)$ , we arrive at the following result

$$\chi_{eff} = \frac{(K_2 + K_1 N f_{EO,salt})(1 - 2f_{EO} + 2f_{EO,salt}^2)(1 + K_3 f_{EO,salt} r + K_4 r)}{N f_{EO,salt}}$$
(29)

Equation 29 in conjunction with Equation 2 is used to calculate  $N_{crit}$ :

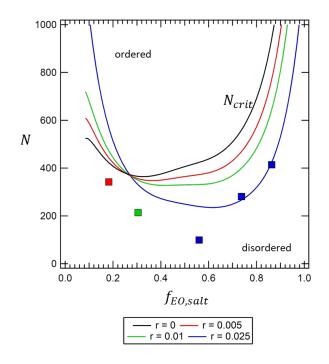
$$(\chi N)_{odt} = \frac{(K_2 + K_1 N_{crit} f_{EO,salt})(1 - 2f_{EO} + 2f_{EO,salt}^2)(1 + K_3 f_{EO,salt} r + K_4 r)}{f_{EO,salt}} = g(f_{EO,salt})(30)$$



**Figure 10**: A plot of *N* versus  $f_{EO}$  for neat SEO copolymers (r = 0i. The curve represents  $N_{crit}$  versus  $f_{EO}$  using a) Equation 25 or b) Equation 29 as  $\chi_{eff}$ . Squares represent the values of *N* and  $f_{EO}$  of the SEO copolymers covered in this study. All of the copolymers are disordered in the neat state and thus we expect the data to lie below the  $N_{crit}$  versus  $f_{EO}$  curve. This is only the case when Equation 29 is used (b).

Equation 30 can be solved to obtain  $N_{crit}$  as a function of  $f_{EO, salt}$  and r. Figure 10b shows the results of these calculations for the salt-free case. The predicted phase boundary remains asymmetric with respect to composition, but it moves to higher values of  $N_{crit}$  for all values of  $f_{EO}$  relative to Figure 10a. In Figure 10b, all of the data points lie underneath the  $N_{crit}(f_{EO})$  curve indicating that all of the copolymers are predicted to be disordered in the absence of salt, consistent with our experiments. It is evident that accounting for the difference between  $\chi_{eff,SC}$  and  $\chi_{eff}$  is essential for quantifying the thermodynamic interactions in SEO/LiTFSI mixtures.

Based on the findings of Figure 10, we take Equation 29 to represent the Flory-Huggins interaction parameter in our system and use it to determine the effect of salt on critical chain length. The curves in Figure 11 show the composition dependence of  $N_{crit}$  for the salt concentrations of interest: r = 0 (black), r = 0.005 (red), r = 0.01 (green) and r = 0.025 (blue). The solid squares in Figure 11 represent the values of N and  $f_{EO, salt}$  for the SEO/LiTFSI mixtures that order at  $r \le 0.025$ . In these mixtures, the addition of salt leads to a transition from disorder to order. The color of each square in Figure 11 indicates the salt concentration at which the SEO/LiTFSI mixture first forms an ordered morphology and matches the color of the  $N_{crit}$  curves. There is excellent quantitative agreement between theory and experiment on the high  $f_{EO,salt}$  side. For  $f_{EO,salt} \leq 0.58$ , the experimental data lie well below the theoretical curves. There is, however, some correspondence between the theoretical curves and experimental data. The curve for r = 0.025 (blue) dips to the lowest value of  $N_{crit}$ , consistent with the data at $f_{EO,salt}=0.58$ . At  $f_{EO,salt}=0.29$ , the curve for r = 0.01 (green) is above the r = 0.025 curve (blue), consistent with the data.



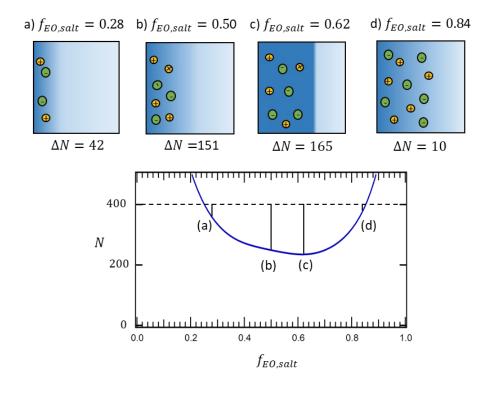
**Figure 11**: Critical chain length for ordering,  $N_{crit}$ , as a function of composition,  $f_{EO,salt}$ , for SEO/LiTFSI mixtures at 100 °C for different salt concentrations: r=0 (black), r=0.005 (red), r=0.01 (green), and r=0.025 (blue). The lines represent the theoretical curves and the squares represent the SEO copolymers discussed in this study. The color of the squares denotes the salt concentration needed to achieve an ordered phase.

We see the emergence of two regimes of salt-dependent phase behavior on either side of  $f_{EO,salt}$ =0.27, which we refer to as the isotaksis point. At this composition,  $N_{crit}$  is independent of salt concentration (all curves for  $N_{crit}$  intersect at  $f_{EO,salt}$ =0.27), i.e., the addition of salt does not affect order in the SEO/LiTFSI mixtures of this composition. We chose the term "isotaksis point" because the word "taksis" means order (spelled  $\tau \alpha \xi_{IG}$ in Greek). At  $f_{EO,salt}$ >0.27, the addition of salt stabilizes the ordered phases and the order-disorder boundary drops to lower values of  $N_{crit}$ . At  $f_{EO,salt}$ <0.27, the addition of salt stabilizes the disordered phase and the order-disorder boundary moves to higher values of  $N_{crit}$ .

The segregation strength in pure block copolymers is characterized by  $\Delta(\chi N)$  defined as  $\Delta \chi N = \chi N - (\chi N)_{odt}$ . At a given temperature and chain length,  $\Delta(\chi N)$  is maximum at  $f_A = 0.50$ . However, this is only valid if  $\chi$  is independent on composition and chain length. This simplification is not valid for SEO/LiTFSI mixtures, and the effect of composition on segregation strength is, perhaps, non-intuitive. We consider copolymers of different compositions at a single chain length and salt concentration. Such systems are represented by a

horizontal line at the chosen value of N in Figure 11. For concreteness, we choose N = 400 and r = 0.025. In Figure 12, we re-plot the  $N_{crit}(f_{EO, salt})$  curve for r = 0.025. Note that the salt concentration under these constraints is proportional to the product  $rf_{EO, salt}$ . The degree of segregation at 100 °C in systems with different  $f_{EO,salt}$  can be quantified by  $\Delta N = N - N_{crit}$ . We show values of  $\Delta N$  at selected  $f_{EO,salt}$  values in Figure 12. Also shown in Figure 12, are schematics that illustrate the degree of segregation at each  $f_{\rm EO, salt}$  value. For simplicity, we focus on the compositions of the salt-rich and salt-poor microphases. The salt is assumed to lie in the PEO-rich microphase, but the degree of mixing between the PS and PEO blocks, governed by  $\Delta N$ , is indicated by the color contrast and interfacial width in each schematic. At f<sub>EO,salt</sub>=0.28(Figure 12a), the SEO/LiTFSI mixture is barely microphase separated as  $\Delta N$ =42. Here we see a microphase separated morphology with a broad interfacial region. The interface sharpens and the compositional contrast between the microphases increases when  $f_{\rm EO, salt}$  is increased to 0.50 (Figure 12b);  $\Delta N$ =151 at this composition. Increasing  $f_{EO, salt}$  further to 0.62 results in the highest degree of segregation possible in systems with N = 400and r = 0.025 (Figure 12c);  $\Delta N = 165$  at this composition. This is in contrast with conventional block copolymer melts, where the highest degree of segregation occurs at  $f_A = 0.50$ .<sup>14</sup> Finally, at  $f_{EO, salt} = 0.84$  (Figure 12d), the degree of segregation decreases and  $\Delta N=10$ . Outside of  $0.28 \le f_{EO, salt} \le 0.84$ , the SEO/LiTFSI mixtures at r = 0.025 and N = 400 form disordered phases.

Qualtitatively different behavior would be obtained at other values of *N* and *r*. It is evident from Figures 11 and 12 that the relationship between segregation strength and copolymer composition in SEO/LiTFSI is complex.



**Figure 12**: Schematic depicting effect of salt on the degree of segregation,  $\Delta N$ , of microphase separated of SEO/LiTFSI mixtures with increasing  $f_{EO, salt}$  at a constant salt concentration, r = 0.025, and total chain length, N = 400. The curve represents  $N_{crit}$  versus  $f_{EO, salt}$  taken from Figure 11.  $\Delta N$  is defined as the difference between the chosen N and the curve, and the vertical line segments indicate its value at selected values of  $f_{EO, salt}$ . The salt is assumed to lie in the PEO-rich microphase, but the degree of mixing between the PS and PEO blocks, governed by the magnitude of  $\Delta N$ , is indicated by the color contrast and interfacial width in each schematic.

### MODEL FOR THE ISOTAKSIS POINT

As mentioned in the Introduction, the thermodynamic properties of block copolymer/salt mixtures are complex and governed by many effects including electrostatics, ion solvation, ion-ion and ion-polymer correlations. Constructing a quantitative theory that is consistent with our data is outside the scope of this paper. In fact, the origin of the composition dependence of  $\chi_{0,SC}$ , the interaction parameter in SEO without salt, is unclear. Nevertheless, it seems important to present a plausible explanation for the existance of the isotaksis composition.

We assume that the ions are distributed uniformly in the disordered state, but they are confined to the PEO microphase in the ordered state, similar in spirit to the original work of Marko and Rabin.<sup>56</sup> If we assume that the salt ions are ideal, then the ion entropy change due to order formation,  $\Delta S_{order}$ , is given by

$$\Delta S_{order} = nk_B \ln f_{EO}$$
 (31)

where *n* is the number of independent ions and  $k_B$  is the Boltzmann constant. For simplicity we ignore the difference between  $f_{EO}$  and  $f_{EO,salt}$  in this analysis which is a reasonable approximation for  $r \le 0.025$ , the range of salt concentrations covered in Figure 11. If we assume a lattice model wherein  $n_t$ is the total number of lattice sites,  $n_{EO}$  is the number of lattice sites occupied by EO and each lattice site is either occupied by a polymer segment or a salt ion, then

37

$$\frac{\Delta S_{order}}{n_t k_B} = r \frac{n_{EO}}{n_t} \ln(f_{EO}) = r f_{EO} \ln(f_{EO})$$
(32)

Note that the entropic contribution is a linear function of *r*.

We assume that this tendency to disorder is balanced by the solvation energy<sup>33,34</sup> that induces ordering, and is quantified by Equation 4. We thus define a theoretical  $\chi$ ,  $\chi_{th}$ , that is the sum of the two contributions:

$$\chi_{th} = \chi_0 + m_n r + \frac{r f_{EO} \ln(f_{EO})}{f_{EO}(1 - f_{EO})}$$
, (33)

where  $m_n$  is a nominal value of m that is assumed to be independent of composition. In many previous studies on mixtures of salt and SEO copolymers<sup>35,51,55</sup>,  $m_n$  has been shown to be about 1.7. At the isotaksis point ( $f_{EO,it}=0.31$ ),

$$\frac{d\chi_{th}}{dr} = 0 = m_n + \frac{f_{EO,it} \ln(f_{EO,it})}{f_{EO,it} (1 - f_{EO,it})}, (34)$$

and

$$m_n = \frac{-f_{EO,it} \ln(f_{EO,it})}{f_{EO,it}(1 - f_{EO,it})}$$
 (35)

For a system with  $m_n=1.7$ , the isotakis point according to Equation 35 is predicted to occur at  $f_{EO}=0.31$ . The experimentally measured isotaksis composition is 0.27. The quantitative agreement between experiments and the model is probably fortuitious, as the model is highly simplified. For example, one could envision placing a prefactor of 2 on the right hand side of Equation 32 to account for salt dissociation. However, the extent to which positively and negatively charged ions are independent in low dielectric media like PEO remains unclear. In addition to this, a complete theory for the isotaksis composition would require inclusion of electrostatic interactions, ion correlations, and physical crosslinking due to interactions between ions and polymer backbones.

### CONCLUSIONS

We have characterized thermodynamics of SEO/LiTFSI mixtures by analyzing SAXS scattering patterns from disordered systems. RPA was used to determine  $\chi_{sc}$ , as a function of block copolymer composition, chain length, temperature and salt concentration. In the neat copolymers,  $\chi_{0,sc}$  is a linear function of  $(Nf_{EO})^{-1}$ . At a given temperature, block copolymer composition and chain length,  $\chi_{eff,SC}$  increases linearly with added salt. The data are only weak functions of temperature as shown in Figures 4 and 6, and therefore these conclusions apply across the entire temperature window studied (  $75 < T(^{\circ}C) < 130$ ). The framework of Sanchez<sup>15</sup> was used to determine  $\chi_{eff}$  from  $\chi_{eff,SC}$  in both salty and salt-free systems. We use the term  $\chi_{eff}$  as it represents interactions between PEO/LiTFSI and PS. SCFT results on pure block copolymers<sup>11</sup> are used to determine the relationship between  $\chi_{eff} N$  and composition,  $f_{EO, salt}$ , at the order-disorder transition. We refer to the value of N at the order-disorder transition as  $N_{\rm crit}$ . All of the SEO copolymers used in this study were disordered in the neat state and order upon salt addition. This enables a direct comparison between experimentally determined values of  $N_{crit}$  with theoretical predictions. At  $f_{EO, salt}$  values greater than 0.27, the addition of salt decreases  $N_{crit}$ , i.e., the ordered phase is stabilized. At  $f_{EO, salt}$ values less than 0.27, the addition of salt increases  $N_{crit}$ , i.e., the disordered phase is stabilized. We propose calling  $f_{EO, salt} = 0.27$  the isotaksis point. A simple theoretical model is proposed to predict the existence of this point. The use of  $\chi$  to describe the phase behavior of neat block copolymers is strictly valid in the limit of infinite chain length. Fluctuation effects become important at finite chain lengths and this leads to non-trivial changes in phase behavior.<sup>50,57,58</sup> Strictly speaking, the implication of the isotaksis point is that the phase behavior of SEO/LiTFSI mixtures at this composition should be independent of salt concentration. Instead, experiments show anamalous phase behavior in the vicinity of this composition. In particular the addition of salt to a disordered phase in the vicinity of the isotaksis point gives rise to two coexisting body centered cubic lattices which then disorder before finally ordering into the expected hexagonnally packed cylinder morphology.<sup>59</sup> While further work is required to determine the underpinnings of such observations, the present framework provides a platform to do so.

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Author Contributions:

The manuscript was written through contributions of all authors. All authors have given approval to the final version of the manuscript.

Notes:

The authors declare no competing financial interest.

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## SUPPORTING INFORMATION

Additional information on the salt concentration and temperature dependence of the chain-stretching parameter,  $\alpha$ , for the SEO/LiTFSI mixtures presented in the main text.

# LIST OF SYMBOLS

- *a<sub>i</sub>* statistical segment length of species *i* (nm)
- $b_i$  X-ray scattering length of species *i* (nm mer<sup>-1</sup>)
- $B_i$  scattering length density of species *i* (nm<sup>-2</sup> mer<sup>-1</sup>)
- *C*<sub>i</sub> Self consistent field theory fitting parameters
- *C* electron density contrast (cm<sup>-1</sup>)
- d domain size (nm)
- *f*<sub>A</sub> volume fraction of species A
- *f<sub>EO</sub>* volume fraction of PEO block
- *I(q)* scattering intensity (cm<sup>-1</sup>)
- *I*<sub>dis</sub>(q) disordered copolymer scattering intensity (cm<sup>-1</sup>)

k <sub>B</sub>	Boltzmann constant
т	proportionality constant
$m_{ m it}$	proportionality constant at isotaksis composition
$M_i$	number-averaged molecular weight of species $i$ (kg mol <sup>-1</sup> )
<b>N</b> <sub>i</sub>	number-averaged degree of polymerization of species $i$ (sites chain <sup>-1</sup> )
NA	Avogadro's number
N <sub>crit</sub>	critical chain length for ordering
n	number of independent ions
n <sub>EO</sub>	number of EO sites
n <sub>t</sub>	total number of lattice sites
q	scattering vector (nm <sup>-1</sup> )
$q^*$	primary peak of scattering vector (nm <sup>-1</sup> )
r	salt concentration ([Li <sup>+</sup> ] [EO] <sup>-1</sup> )
$R_{g,i}$	radius of gyration of species <i>i</i> (nm)
$\boldsymbol{S}(\boldsymbol{q})$	scattering structure factor
Т	temperature (K)
$Y_{LiTFSI}$	volume fraction of salt in PEO/LiTFSI microphase

# Greek

α	chain-stretching parameter
$\Delta S_{order}$	change in entropy due to ordering
$\Delta N$	degree of segregation
E	conformational asymmetry parameter
χ	Flory-Huggins interaction parameter
X <sub>sc</sub>	Flory-Huggins interaction parameter from scattering
X₀	Flory-Huggins interaction parameter of salt free system
$\chi_{\scriptscriptstyle th}$	theoretical Flory-Huggins interaction parameter
$\chi_{\scriptscriptstyle eff}$	effective Flory-Huggins interaction parameter
χN	segregation strength
$(\pmb{\chi} \pmb{N})_{odt}$	segregation strength at the ODT
$\boldsymbol{v}_i$	molar volume of species <i>i</i> (cm <sup>3</sup> mol <sup>-1</sup> )
$v_{ref}$	reference volume (nm³ site-1)
$\rho_i$	density of species <i>i</i> (g cm <sup>-3</sup> )
• 1	

# LIST OF ABBREVIATIONS

EO ethylene oxide

LiTFSI	lithium bis(trifluoromethanesulfonyl) imide salt
N <sub>A</sub>	Avogadro's number
ODT	order-disorder transition
PEO	poly(ethylene oxide)
PS	polystyrene
RPA	Random Phase Approximation
S	styrene
SAXS	small angle X-ray scattering
SCFT	Self-Consistent Field Theory
SEO	polystyrene- <i>block</i> -poly(ethylene oxide)

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