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# COLLISIONAL QUENCHING OF $\kappa^*(4_P^2 P)$ AND $\kappa^*(5_P^2 P)$ BY H<sub>2</sub>O, CF<sub>4</sub>, AND CH<sub>4</sub>

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Although a number of workers have studied the gas-phase collisional quenching of the lowest excited configuration of the alkali atoms, 1 very little is known about the quenching of higher energy configurations. In contrast to Hg and rare gas atoms which have also been the subjects of many studies, the excitation energies of the alkali atoms are typically less than the bond energies of molecules which are efficient quenching agents. For example, the excitation energies (in e.V.) of the first six excited configurations of K are: 1.61  $(4p^2P)$ , 2.60  $(5s^2S)$ , 2.67  $(3d^2D)$ , 3.06  $(5p^2P)$ , 3.40  $(4d^2D)$ , and 3.41 (6s<sup>2</sup>S). Nevertheless, a variety of molecules are quite efficient at quenching the lowest excited configuration of the alkali metals, although certain "saturated" molecules, such as  $CH_{\mu}$ ,  $CF_{\mu}$ , and, to a lesser extent,  $H_2O$  are inefficient as quenching agents.<sup>1,2</sup> The work described here was undertaken in order to determine whether this inefficiency would persist for higher energy configurations. Measured rate constants are presented for quenching of  $K^{*}(4p^{2}P)$  and  $K^{*}(5p^{2}P)$  by  $H_{2}O$ ,  $CF_{\mu}$ , and  $CH_{\mu}$ .

The apparatus and experimental procedure was similar to that described in Ref. 2. All measurements were made with ultraviolet radiation from a hydrogen arc which passed through a monochromator with a triangular bandpass function,  $60\text{\AA}$ FWHM, and impinged upon a quartz cell containing KI vapor at 890°K. The K<sup>\*</sup>(4p<sup>2</sup>P) fluorescence intensities at 7665Å and at 7699Å, which were produced by photodissociation of the KI at 2450Å, were isolated by different interference filters. Irradiation at 1925Å permitted measurement of the  $K^*(5p^2P)$  fluorescence through an interference filter which transmitted both the 4044 and 4047Å resonance lines. With the UV intensity held constant at one of these two wavelengths, the rate constant,  $k_q$ , for quenching of the corresponding 4p or 5p potassium fluorescence was determined by admitting a known number density, n, of quenching gas to the KI cell and fitting the measured fluorescence intensity,  $R_n$ , to the Stern-Volmer relation:

 $R_0/R_n = 1 + \tau k_q n.$  (1)

The radiative fluorescence lifetimes,  $\tau$ , were assigned values<sup>3,4</sup> of 2.6x10<sup>-8</sup> sec. for K<sup>\*</sup>(4p<sup>2</sup>P) and 1.4x10<sup>-7</sup> sec. for K<sup>\*</sup>(5p<sup>2</sup>P). Thus, k<sub>q</sub> is determined by assuming that the resonance fluorescence is produced by the direct photodissociation of KI to give the corresponding K<sup>\*</sup>(4p<sup>2</sup>P) or K<sup>\*</sup>(5p<sup>2</sup>P) configuration. This clearly holds for production of K<sup>\*</sup>(4p<sup>2</sup>P) at 2450Å because threshold for production of K<sup>\*</sup>(5s<sup>2</sup>S) by photodissociation of KI in its ground vibrational level is calculated<sup>5</sup> to be 2060Å. Calculated thresholds for production of K<sup>\*</sup>(5p<sup>2</sup>P) and K<sup>\*</sup>(4d<sup>2</sup>D) (or K<sup>\*</sup>(6s<sup>2</sup>S)) are at 1920 and 1820Å respectively. The measured reduced K<sup>\*</sup>(5p<sup>2</sup>P) fluorescence efficiency (i.e., the ratio of fluorescence intensity to incident UV intensity) peaks at ~ 1900Å and decreases to 20% of its peak value at ~ 1760 and 2020Å, indicating that very little of the K<sup>\*</sup>(5p<sup>2</sup>P)

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fluorescence produced by 1925Å can be due to initial production of higher energy configurations. This was further verified in an auxiliary experiment in which no variation was observed in  $k_q$  for CH<sub>4</sub> quenching of the K<sup>\*</sup>(5p<sup>2</sup>P) fluorescence produced by irradiation of KI at 1925, 1950, and 1975Å.

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Measured values of k are listed in Table I. Measurements at 7665Å and 7699Å determined the same  $k_q$  value for  $K^*(4p^2P)$  +  $H_{2}O.$  For  $K^{*}(4p^{2}P) + CH_{4}$  or  $CF_{4}$ , where  $k_{q}$  was too small to measure reliably, some collisional mixing of the two fine structure levels was observed. In contrast to work in Ref. 2 where the translational energy dependence of k was measured,  $k_q$  values listed in Table I were determined for thermal distributions. Relatively large pressures of quenching gas (  $\sim$  20 torr for  $\rm H_2^{}0,$   $\sim$  50 torr for  $\rm CF_4$  and  $\rm CH_4^{})$  were necessary to determine the slow quenching rates of  $K^*(4p^2P)$ , so that the initial K<sup>\*</sup>(4p<sup>2</sup>P) speed distribution should have collisionally relaxed before being quenched. In experiments on  $K^*(5p^2P)$  +  $CH_{4}$  and  $CF_{4}$ , a few torr of the quenching gas was pre-mixed with  $\sim$  650 torr of Ar, and this mixture (at a reduced pressure) was admitted to the KI cell in order to insure that the K<sup>\*</sup>(5p<sup>2</sup>P) speed distribution was thermalized by collisions with Ar. In separate experiments, it was ascertained that pure Ar (at  $\sim$  600 torr) or pure Xe (at  $\sim$  200 torr) failed to produce a measurable attenuation of the  $K^*(5p^2P)$  fluorescence intensity. This "Ar thermalizing" procedure was not used in measuring  $k_{a}$  for  $K^{*}(5p^{2}P) + H_{2}^{0}$  due to experimental difficulties. Here,

however, the distribution in relative collision speeds is primarily determined by the thermal speed distribution<sup>2</sup> of  $H_20$  by virtue of its light mass relative to K. This was verified with  $K^*(5p^2P) + CH_4$  where a value of 6.7  $\pm$  0.7x10<sup>-10</sup> cm<sup>3</sup>/sec was measured for  $k_q$  in the absence of the "Ar thermalizing" procedure, in reasonable agreement with the thermal value given in Table I.

We are not aware of any previous studies on the quenching of K resonance fluorescence by  $CH_4$  or  $CF_4$ . Flame fluorescence studies of  $K^*(4p^2P) + H_20$  have determined quenching cross sections of 2.8  $\pm$  0.9Å<sup>2</sup> at 1400 - 1800°K<sup>6</sup> and 2.6  $\pm$  0.3Å<sup>2</sup> at  $\sim$  2000°K,<sup>7</sup> in reasonable agreement with our result in view of the temperature differences. On the other hand, the only previous study<sup>7</sup> of  $K^*(5p^2P)$  determined (at 2000°K) a quenching cross section for  $H_20$  of 10  $\pm$  4Å<sup>2</sup>, substantially different from the result obtained here.

In Ref. 2, we chose to compare our measured cross sections,  $Q_q$ , for quenching of Na<sup>\*</sup>(3p<sup>2</sup>P) with cross sections,  $Q_D$ , for "close encounters" wherein the incident trajectories surmount the centrifugal barriers in the effective potentials associated with the long-range attractive dispersion forces. In doing so, we did not propose that  $Q_q$  should equal  $Q_D$  because (1) the description of the quenching collision must include some matrix element coupling different electronic states and (2)  $Q_D$  provides only an estimate of the cross section for close encounters because of neglect of shorter range forces.

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Nevertheless, this comparison proved useful in assessing the efficiency of different molecules for quenching Na<sup>\*</sup>(3p<sup>2</sup>P). Molecules with appreciable electron affinities (e.g., I<sub>2</sub> and SO<sub>2</sub>) are "super-efficient" with Q<sub>q</sub> > Q<sub>D</sub>, unsaturated molecules (e.g., N<sub>2</sub> and C<sub>2</sub>H<sub>4</sub>) as well as some potentially reactive compounds (e.g., CF<sub>3</sub>Cl) are efficient quenchers with Q<sub>q</sub>  $\leq$  Q<sub>D</sub>, and certain saturated molecules (e.g., H<sub>2</sub>O, CF<sub>4</sub> and CH<sub>4</sub>) are very inefficient with Q<sub>q</sub> << Q<sub>D</sub>. A similar comparison, shown in Table I in the form of rate constants (k<sub>q</sub>/k<sub>D</sub>) rather than cross sections, indicates that H<sub>2</sub>O, CF<sub>4</sub>, and CH<sub>4</sub> are also very inefficient at quenching K<sup>\*</sup>(4p<sup>2</sup>P). For K<sup>\*</sup>(5p<sup>2</sup>P), however, H<sub>2</sub>O, CH<sub>4</sub>, and, to a lesser extent, CF<sub>4</sub> prove to be quite efficient as quenching agents.

The quenching collisions studied here refer to some unknown combination of inelastic processes producing lower K configurations and reactive processes. For  $H_20$  and  $CF_4$ , especially, the large efficiency for quenching of  $K^*(5p^2P)$ might simply be a consequence of the relatively large reaction exoergicities listed in Table I; this point of view is advanced by Dowling, et. al.<sup>8</sup> in discussing their results on quenching of Na $^*(3p^2P)$  and Tl $^*(7s^2S)$ . At any rate, the failure of Ar or Xe to quench  $K^*(5p^2P)$  suggests that  $H_20$ ,  $CF_4$ , and  $CH_4$  quench this configuration either by reaction or by an inelastic process wherein electronic energy is at least partially converted into vibrational and rotational excitation. Furthermore, the failure of  $CH_{\rm h}$  to efficiently quench  $Hg^*(6p^3P_{0,1})$ ,<sup>9</sup> despite the larger excitation energies, suggests that the dramatic difference shown in Table I is not simply a consequence of higher density of vibrational rotational levels in the quenching molecules corresponding to the higher  $K^*(5p^2P)$  excitation energy.

One possible mechanism for the efficient quenching of  $K^{*}(5p^{2}P)$  would picture the collision as proceeding via a charge-transfer intermediate which could produce lower K configurations or an "ionically bound" KX product. This venerable model of collisional quenching of excited alkali atoms has recently been employed to interpret<sup>10</sup> the quenching of Na<sup>\*</sup>(3p<sup>2</sup>P) by N<sub>2</sub>; it may also be operative<sup>11</sup> in the related (not precisely inverse) process of electronic excitation in collisions of fast alkali atoms. Appearance potentials of 5.0 - 5.7,  $\sim$  7.5, and 4.7 eV have been reported for dissociative electron attachment to  $H_20$ , <sup>12</sup>  $CH_{\mu}$ , <sup>12</sup> and  $CF_{\mu}$ , <sup>13</sup> respectively, whereas vertical electron affinities of  $\sim$  -3 to  $\sim$  -4 eV would appear to be necessary in order to account for the observed contrasting efficiencies for quenching of  $K^{*}(4p^{2}P)$ and  $K^*(5p^2P)$  in terms of this charge-transfer model. However, these dissociative electron attachment results may not refer to excitation of the lowest negative ion resonance states. Indeed, Claydon, et. al.<sup>14</sup> have reviewed the known data on negative ion states of  $H_2^0$  and have concluded that the lowest resonance state is not seen in dissociative electron attachment because it is bound relative to dissociation into OH + H,

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suggesting a vertical electron affinity greater than -3.28 eV, probably between -2 and -3 eV. In terms of the simple charge transfer model, a value near -3 eV would seem to account for the large value of  $k_{a}$  for  $K^{*}(5p^{2}P)$ , the relatively small values of  $k_a$  for  $K^{*}(4p^2P)$  and  $Na^{*}(3p^2P)$ , as well as the trend<sup>1</sup> of increasing  $k_q$  on proceeding from  $K^*(4p^2P)$  to  $Rb^*(5p^2P)$  and Cs<sup>\*</sup>(6p<sup>2</sup>P). Similarly, Christophorou and Stockdale<sup>15</sup> have suggested that the lowest energy dissociative electron attachment observed in  $\text{CH}_{\!\mu}$  proceeds via excitation of a core-excited resonance level. Thus, it seems plausible to attribute the quenching behavior of  $H_20$ ,  $CH_{\mu}$ , and  $CF_{\mu}$  observed here to formation of charge-transfer intermediates involving the lowest energy resonance states of the negative ions. It should also be noted, however, that the increasing density of electronic terms with increasing excitation energy might be expected to lead to avoided crossings and efficient quenching even in the absence of formation of such a charge-transfer intermediate.

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### REFERENCES

For a review of the literature, see: P. L. Lijnse,
 "Review of Literature on Quenching, Excitation and Mixing
 Collision Cross Sections for the First Resonance Doublets
 of the Alkalis". Report i-398, Fysisch Laboratorium,
 Rijks Universiteit Utrecht, The Netherlands, February, 1972.

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- B. L. Earl, R. R. Herm, S.-M. Lin, and C. A. Mims, J.
  Chem. Phys. <u>56</u>, 867 (1972).
- W. L. Wiese, M. W. Smith, and B. M. Miles, <u>Atomic Transition</u> <u>Probabilities</u>. Vol. II, Nat. Stand. Ref. Data Ser., Nat. Bur. Stand. (U.S.) 22, 1964.
- 4. Corrections for radiation imprisonment of K resonance
  fluorescence are negligible because potassium is produced
  in very low concentrations by photodissociation of KI.
- 5. Bond dissociation energies for KI and KH were taken from A. G. Gaydon, <u>Dissociation Energies and Spectra of</u> <u>Diatomic Molecules</u>, Chapman and Hall, London, 1968; other values are from B. deB. Darwent, <u>Bond Dissociation Energies</u> <u>in Simple Molecules</u>, Nat. Stand. Ref. Data Ser., Nat. Bur. Stand. (U.S.), 31, 1970.
- 6. D. R. Jenkins, Proc. Roy. Soc. A303, 453 (1968).
- H. P. Hooymayers and P. L. Lijnse, J. Quant. Spectrosc.
  Radiat. Transfer 9, 995 (1969).
- D. J. Dowling, G. R. H. Jones, and E. Warhurst, Trans.
  Faraday Soc. 55, 537 (1959).

0 0 0 0 3 9 0 5 3 2 1

-9-

- 9. A. C. Vikic and H. C. Moser, J. Chem. Phys. 53, 1491 (1970).
- E. Bauer, E. R. Fisher, and F. R. Gilmore, J. Chem. Phys.
  51, 4173 (1969).
- 11. K. Lacmann and D. R. Herschbach, Chem. Phys. Letters <u>6</u>, 106 (1970).
- 12. L. G. Christophorou, <u>Atomic and Molecular Radiation Physics</u>, John Wiley, New York, 1971.
- 13. K. A. G. MacNeil and J. C. J. Thynne, Int. J. Mass Spectrom. Ion Phys. <u>3</u>, 455 (1970).
- C. R. Claydon, G. A. Segal, and H. S. Taylor, J. Chem. Phys.
  54, 3799 (1971).
- L. G. Christophorou and J. A. D. Stockdale, J. Chem. Phys.
  48, 1956 (1968).

K configurat:	quenching ion gas	10 <sup>9</sup> k <sub>q</sub> (cm <sup>3</sup> /sec) <sup>a</sup>	Q <sub>q</sub> (Å <sup>2</sup> ) <sup>b</sup>	k <sub>q</sub> ∕k <sub>D</sub> <sup>c</sup>	∆D <sub>o</sub> (eV) <sup>d</sup>
	CF <sub>4</sub>	<0.01	<1	<0.01	1.3(KF + CF <sub>3</sub> )
4p <sup>2</sup> P	сн <sub>4</sub>	<0.01	<1	<0.01	-0.9(KH + CH <sub>3</sub> )
4p <sup>2</sup> P	H <sub>2</sub> O	0.024±0.008	1.5	0.025	0.0(KOH + H)
5p <sup>2</sup> P	CF4	0.21±0.04	26	0.16	2.8(KF + CF <sub>3</sub> )
5p <sup>2</sup> P	сн <sub>4</sub>	0.77±0.08	60	0.45	0.6(KH + CH <sub>3</sub> )
5p <sup>2</sup> P	H <sub>2</sub> O	1.0±0.1	85	0.71	1.5(KOH + H)

Table I. Measured quenching rate constants, k<sub>q</sub>, at 890°K.

<sup>a</sup>The uncertainty given for each  $k_q$  value is simply the standard deviation provided by a least squares fit of the pressure dependence of the fluorescence intensity to Eq. (1). The true error in  $k_q$  should be somewhat larger.

<sup>b</sup>Quenching cross section, calculated as  $k_q = Q_q < g >$ , where < g > is the thermal average relative speed.

<sup>c</sup>As discussed in the text,  $k_q/k_D$  is a measure of the quenching efficiency. <sup>d</sup> $\Delta D_o$  is the excergicity of the possible reaction; data from Ref. 5. -LEGAL NOTICE-

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