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7Li + 7Li REACTION STUDIES LEADING TO MULTI-NEUTRON FINAL STATES

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 7 Li + 7 Li REACTION STUDIES LEADING TO MULTI-NEUTRON FINAL STATES

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September 1974

Boron and carbon reaction products have been observed in the bombardment of 7 Li by 79.6 MeV ⁷Li. Comparisons of the 7 Li(⁷Li, 11 B)t and 7 Li(7 Li, 11 C) 3n channels and a study of the 7 Li(7 Li, 10 C) 4n reaction are presented.

Although there has been extensive historical interest in questions of the possible stability of $\frac{3}{n}$ or $\frac{4}{n}$, and of the location of unbound resonances in these systems, no bound states nor uncontroversial multi-neutron resonance effects have so far been established in either of these systems [see ref. $[1]$ for a review of the 3n system; ref. [2] , for the 4n]. Nonetheless, since certain heavy-ion reactions observing neutron-deficient reaction products afford a new look at these (and other [3]) multi-neutron final states, we have investigated one of the simplest of these systems, that of $1u + 7u + 12c + 2n$, 11_C + 3n, and 10_C + 4n. By also measuring the energy spectra and cross-sections of the boron isotopes in the better-established 12 B + d, 11 B + t and 10 B + 4 H

* Work performed under the auspices of the U. S. Atomic Energy Commission. [†]Present address: Nuclear Physics Laboratory, University of Oxford, England.

channels (but ones in which the light product nuclei have lower T_{7}), one can hope to obtain some criteria by which to evaluate the yield in the carbon exit channels. Four of these reactions are discussed below; unfortunately, reactions on target contaminants precluded useful analysis of the 7 Li(7 Li, 12 C)2n and 7 Li(7 Li, 10 B) 4 H results.

A beam of 79.6 MeV $1.1+2$ (~ 150 nA) from the Lawrence Berkeley Laboratory 88-inch cyclotron was used to bombard a 110 μ gm/cm 2 7 Li target. Reaction products were observed in two similar counter telescope systems placed at opposite sides of the beam. The data reported below came from the system placed at 7.4° (lab) with a 0.086 msr solid angle; it consisted of two transmission (ΔE) detectors, 18 and 14 μ m thick (the first with subnanosecond pile-up rejection $[4]$), a 190 μ m E detector, and a reject detector. Although equivalent results were obtained with the second system, which was placed at 9.6° , they were of poorer quality. Other experimental details were similar to those described previously [5]; a comparison of two particle identification signals was employed to reduce background, with a stringent comparison rejecting \sim 50% of the events traversing the telescope. Electronic and beam energy stability were monitored continuously, and the absolute beam energy was determined using a precision analyzing magnet.

Figure 1 presents particle identification spectra from ⁷Li bombardment of 16 and 16 (as SiO₂) targets. Since reactions on carbon and oxygen were a severe background problem, the latter spectrum is shown to permit comparison of relative isotopic yields from one of the major target contaminants. As can be seen from the figure, 11 C, 11 B and 12 B were relatively strongly produced in the 7 Li + 7 Li reaction. This high 11 C, but low 10 C, yield from 7 Li led to poor 10 _C - 11 _C separation, so that only those 10 _C events whose particle identification signal fell in the lower half of its nominal spectral position were accepted for energy analysis.

The energy spectrum of the 7 Li(7 Li, 12 B)d reaction is shown in fig. 2. Moderate population of the (unresolved) bound states of 12 B can be seen in this $spectrum, with a composite cross-section of 4 $\mu b/sr$ c.m. Experimental observation$ of the strength of the two-neutron final state interaction in the 7 Li(7 Li, 12 C)2n data would be of great interest, but,as noted above, could not be observed due to interference by contaminant reactions.

Results from the 7 Li(7 Li, 11 B)t and 7 Li(7 Li, 11 C) 3n reactions are compared in fig. 3(a} and 3(b-c}, respectively. Transitions to a number of the bound 11 B final states can be seen which are substantially stronger than those in the 12 B + d data; in particular the ground state transition possesses a cross section of 23 μ b/sr c.m. However, the 7 Li(7 Li, 11 C) 3n data per se in figs. 3 (b) and 3(c) present no discernible structure. At this small forward angle the 11 C energy region that would correspond to a bound 3n system is free from reactions on target contaminants, and an upper limit of 70 nb/sr c.m. can be set for production of a bound $\frac{3}{n}$. Two imperfect comparisons are available: this limit is a factor of \sim 300 less than the yield of the 11 B g.s. + t channel, and is a factor of \sim 12 less than the average yields of the 16 O(7 Li, 11 C) 12 B g.s. and 12 C(7 Li, 11 C)⁸Li g.s. reactions at forward angles (obtained from separate experiments}. with regard to those transitions corresponding to an unbound 3n system, one sees in fig. 3(c) that the 11 C energy spectrum encompassing up to \sim 7 MeV excitation of three neutrons (before the bulk of the transitions from target contaminants begins) is well fit by four-body phase space.

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Figure 4 presents an energy spectrum from the attempted three-proton transfer 7 Li(7 Li, 10 C)4n reaction. Independent experiments on the 16 O(7 Li, 10 C)¹³B and 12 C(⁷Li,¹⁰C)⁹Li reactions successfully observed the transfer of three protons with comparable ground state cross sections, averaging \sim 450 nb/sr c.m. Peaks from reactions on these target contaminants account for the observed structure in the 4n continuum region of fig. 4; the underlying background appears to be adequately fit by five-body phase space. Again, at this forward angle, contaminant reactions do not interfere in the region of the 10^c energy spectrum corresponding to transitions leading to a bound $\frac{4}{n}$ [the known mass of $\frac{8}{n}$ He sets an upper limit to the total binding energy of $\frac{4}{n}$ (see ref. 2)]. The very minor background observed in this region arises from the 11 C "leak-through" remaining in this energy spectrum; however, it is still possible to set an upper limit of 30 nb/sr c.m. for the cross-section of this reaction leading to a bound 4n system. The only available comparison is to note that this limit is a factor of \sim 15 less than the yield of the observed three-proton transfer reactions on 12 C and 16 O.

These results set stringent limits $[1,2]$ in failing to observe transitions to a bound $\begin{smallmatrix}3&&4\1\end{smallmatrix}$ n; further, no resonance structure was evident in these heavy-ion studies of the unbound 3n and 4n systems. Better data on the unbound 4n system (requiring rigid maintenance of the $7L$ i target purity) would permit an interesting comparison with the 4 He(π , π ⁺)4n studies [6], in which a possible final state interaction· is observed between one neutron pair in the exit channel. Clearly the above approach can also be extended to search for bound or unbound structure in higher neutron configurations.

REFERENCES

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FIGURE CAPTIONS

Fig. 1. Particle identification spectra arising from 7 Li reactions on 7 Li and 16 ¹⁶0 targets. The relative intensities of all boron peaks are low by a factor of \sim 0.25. Certain weak groups such as 14 B in the 7 Li + 7 Li data must arise from'target contaminants.

Fig. 2. An energy spectrum from the 7_{Li} $(^{7}_{Li}$, $^{12}_{B}$)d reaction at 79.6 MeV and

- 7.4⁰. Dashed arrows denote the expected locations of the indicated transitions. Fig. 3. Spectra from the $\frac{7}{4}$ Li + $\frac{7}{4}$ Li reaction at 79.6 MeV.
	- (a) 7 Li(7 Li, 11 B)t. Dashed arrows denote the expected location of contaminant reactions.
	- (b) 7 Li(7 Li, 11 C)3n. See (a). An arrow with an asterisk denotes the location of a known contaminant reaction. Also indicated is the 11 C energy that would correspond to transitions to a three neutron system with zero binding energy (B.E.).

(c) A detail of the high-energy part of (b).

Fig. 4. An energy spectrum from the $7_{Li} (7_{Li} 10_{C}) 4n$ reaction at 79.6 MeV and 7.4^o. Known contaminant reactions are indicated either explicitly or by an arrow with an *asterisk.*

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Fig. 1

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 $Fig. 2$

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Fig. 4

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