Lawrence Berkeley National Laboratory

Recent Work

Title

HEAVY-ION FUSION SYSTEM ASSESSMENT PROJECT Quarterly Status Report, April 1 - June 30, 1985.

Permalink <https://escholarship.org/uc/item/7zk3w78q>

Authors

Lee, E.P. Hovingh, J. Faltens, A.

Publication Date

1985-06-01

195-20759

Lawrence Berkeley Laboratory

UNIVERSITY OF CALIFORNIA

H.t:..Lt:.JVc_,_ **LAWRENCE**

Accelerator & Fusion Research Division

• _': I 1986

"

'SRARY AND

HEAVY-ION FUSION SYSTEM ASSESSMENT PROJECT: Quarterly Status Report, April 1-June 30, 1985

E.P. Lee, J. Hovingh, and A. Faltens

June 1985

TWO-WEEK LOAN COP

This is a Library Circulating Copy which may be borrowed for two week

Prepared for the U.S. Department of Energy under Contract DE-AC03-76SF00098

DISCLAIMER

This document was prepared as an account of work sponsored by the United States Government. While this document is believed to contain correct information, neither the United States Government nor any agency thereof, nor the Regents of the University of California, nor any of their employees, makes any warranty, express or implied, or assumes any legal responsibility for the accuracy, completeness, or usefulness of any information, apparatus, product, or process disclosed, or represents that its use would not infringe privately owned rights. Reference herein to any specific commercial product, process, or service by its trade name, trademark, manufacturer, or otherwise, does not necessarily constitute or imply its endorsement, recommendation, or favoring by the United States Government or any agency thereof, or the Regents of the University of California. The views and opinions of authors expressed herein do not necessarily state or reflect those of the United States Government or any agency thereof or the Regents of the University of California.

HEAVY-ION FUSION SYSTEM ASSESSMENT PROJECT*

Quarterly Status Report, Apr.-Jun. 1985

E. P. Lee, J. Hovingh, and A. Faltens

University of California Lawrence Berkeley Laboratory One Cyclotron Road Berkeley, CA 94720

*Work performed under the auspices of the U.S. Department of Energy by Lawrence Berkeley Laboratory under contract DE-AC03-76SF00098. Funded by the Office of Program Analysis.

ACCELERATOR SYSTEM MODELING

1. The cost Optimization Code LIACEP

1.1 Code Modifications

.. '

忥

'

The cost optimization code LIACEP has been modified to operate on the VAX(8600) computer. The code requires from 3% to 5% minutes of CPU time per design case. This compares with CPU times of 1 to 1% minutes per design case on the CDC 7600.

The tune-shift tables in LIACEP have been replaced by the analytic approximation formulated by Lee et al. (1) A comparison of the accumulative costs and efficiency of a 3 MJ drive, based on code runs using the analytical approximation vs. the tune-shift tables, is given in Table 1. The analytic approximation overestimates the cost by about 4.3% at 50 MV and underestimates the cost by 0.8% at 2500 KV. The total cost is underestimated by only 0.08% using the analytic approximation for the tune-shift tables.

More than 100 runs have been made with LIACEP since these changes were made. Most of these runs have been in support of the Heavy Ion Fusion System Assessment Project .

Table I. Comparison of a Cost and Efficiency Analysis of an Induction Linear Accelerator Using an Analytic Approximation with Tune-Shift Tables.

> $A = 200$ amu $\epsilon_n = 1.17 \times 10^{-5} \text{ m-radius}$ $N = 4$ beams $q = {}^+1$ $f = 1$ hertz $\sigma_{0} = 60^{\circ}, \quad \sigma = 24^{\circ}$

V

Superconducting quadrupoles, amorphous iron cores.

 $It/beamlet = 75\mu$ Coulombs

Ratio of cumulative costs and efficiency: (Analytic/Tables)

1.2 Calculations for the Heavy Ion Fusion System Assessment.

An early assessment of the optimized cost of the high energy end of induction linear accelerator was made for 200 amu ions with a charge state of +1, an input ion energy of SO MeV and an output ion energy of 10 GeV. The accelerator had a normalized emittance of 1.17 x 10^{-5} m-radians for each of the 4 beamlets. The focussing element consisted of superconducting quadrupoles and the acceleration cores are of amorphous iron. The initial tune for these cases is 60°, while the depressed tune is 24°. For a pulse repetition frequency of 5 hertz, the cost is well approximated by a constant plus a term which varies linearly with the beam energy.*

$$
C \n f=5,E
$$
 α E + 2.9 $1 \le E \le 10$ MJ

For an output beam energy of 3 MJ, the cost varies similarly with the pulse repetition frequency

> C_{n} α f + 112 1 \le f \le 10 hertz $E=3, f$

Thus the cost of this heavy ion induction linear accelerator is a much stronger function of the output beam energy than the pulse repetition frequency. The efficiency of the accelerator defined as the ratio of the output beam power, P, to the electrical power into the accelerator is given as

> $n \approx 5.1 \text{ P}^{1/2}$ $P \le 11.8 \text{ MW}_{B}$ $n \approx 10.7 \text{ P}^{1/5}$ P > 11.8 MW_p

ਨਾ

where the output beam power $P = EF$, where f is the pulse repetition frequency in hertz and E is the total beam output energy in MJ. We also varied the

* This scaling of cost with energy appears more adverse than it actually is in an optimized driver (where $C \propto E^{-4}$) due to the constraint on normalized emittance.

number of beamlets, and found the minimum cost case was between 4 and 16 beamlets, as shown in Table II. The minimum cost case at 3 beamlets is only 3.5% less than the cost of 4 beamlets.

 $\overline{6}$

Table II. Cost and Efficiency as a Function of Number of Beamlets for 3 MJ Driver.

> $A = 200$ amu $\epsilon_n = 1.17 \times 10^{-5} \text{ m- radians}$ $q = +1$ $f = 5 \text{ hertz}$ $\sigma_{0} = 60^{\circ}, \quad \sigma = 24^{\circ}$ Superconducting quadrupoles, amorphous iron cores Initial Ion Energy = 50 MeV Final Ion Energy = 10 GeV

tz

×,

TABLE III. Accelerator Parameter Space Investigated for

Heavy-Ion Fusion System Assessment.

Ion Mass 130, 160, 190, 210 amu

Ion Kinetic Energy 5, 10, 15, 20 GeV

Beam Energy 1, 2, 3, 5, 10 MJ

Emittance 1.5×10^{-5} , $(3 \times 10^{-5}) \times \text{m-radians}$

Pulse Repetition Frequency 5, (10), (15), (20)* hertz

Number of Beamlets $4, (8), (16)*$

Ion Charge State $+1$

Tune: 60°, Depressed Tune: 24°**

Initial Ion Kinetic Energy 50 MeV

.R

 $\mathcal{S}_\mathbf{d}$

"' ,.

Focussing System: Superconducting Quadrupoles

Core Material: Amorphous Iron

 $*()$ - not completed ** Recent experiments show that a depressed tune of 8° can be achieved. This will lead to cost savings. $\mathcal{L}^{\text{max}}_{\text{max}}$, $\mathcal{L}^{\text{max}}_{\text{max}}$

 $\sim 10^7$

The parameter space investigated to date for the Heavy Ion System Assessment Project is given in Table III. The emittance of 1.5 x 10^{-5} ~radians was selected to give a spot radius on target of 1 mm, due to emittance, and allowing an effective spot radius of 2 mm due to the additional effects of momentum spread, and various aberrations in the final transport and focussing system. The selection of an initial tune of 60° and a depressed tune of 24° is arbritrary, as larger initial tunes and smaller depressed tunes have been achieved in the laboratory.

r

Ţ4

The cost and efficiencies of the accelerator system with an initial ion energy of 50 MeV are given in Tables IV through VII for various ion energies and beam energies. The costs are shown in Figures 1 through 4. All the costs have been normalized because the current costs for material and labor have not yet been included in LIACEP. In addition, an appropriate parameter space has not yet been identified on which to focus further analysis. Upon identification of an appropriate parameter space, a point design will be made, including a layout of the accelerator modules, which will suggest cost saving techniques and better cost estimates. However, the present analysis is adequate for determining a design parameter space.

For a given total beam energy W, there is an ion energy for which the normalized cost of the accelerator system is minimized as shown in Tables IV through VII and in Figures 1 through 4. The accelerator efficiency is maximized at the same points the cost is minimized. The reason for the cost minimum is that the accelerator cost increases with voltage and increases with charge, but the charge decreases with an increase in voltage for a given total beam energy. In addition, for a fixed output emittance, the normalized emittance of the accelerator increases with output voltage, and the cost of the accelerator with a fixed output voltage and beam energy and species will decrease with an increase of the normalized emittance.

TABLE IV. Normalized Cost & Efficiency of A=130 anu, 4 Beamlet Accelerator System with f=5 hertz, $\epsilon = 1.5 \times 10^{-5}$ m-radians, $\ddot{}$ $q = +1$, and a Tune of $60^{\circ}-24^{\circ}$, with Superconducting Quads and Amorphous Iron Cores.

ε _j , GeV	5	10	15	20
ϵ_n , m-rad W, MJ	4.34 $\times 10^{-6}$	6.20×10^{-6}	7.67×10^{-6}	8.94×10^{-6}
1	0.9415	$0.7945 \div$	0.8697	0.9925
$\mathbf{2}$	1.550	1.114	1.094 \leftarrow	1.171
$\mathbf{3}$	2.091	1.419	1.308 \leftarrow	1.335
5	3.277	1.993	1.699	1.645 \leftarrow
10	6.045	3.322	2.604	2.342 \leftarrow
		$($ ϵ Minimum)		

Normalized Cost

 $($ Maximum $)$

Normalized Cost & Efficiency of A=160 anu, 4 Beamlet TABLE V. Accelerator System with f=5 hertz, ϵ = 1.5 x 10⁻⁵ m-radiars, $q = +1$, and a Tune of $60^{\circ}-24^{\circ}$, with Superconducting Quads and Amorphous Iron Cores.

Þ.

77

Ļ,

Normalized Cost

 $($ Maximum $)$

 Λ

TABLE VI. Normalized Cost & Efficiency of A=190 anu, 4 Beamlet Accelerator System with f=5 hertz, ϵ = 1.5 x 10⁻⁵ m-radians, $q = +1$, and a Tune of $60^{\circ}-24^{\circ}$, with Superconducting Quads and Amorphous Iron Cores.

ϵ_i , GeV	5	10	15	20
ϵ_n , m-rad W, MJ	3.58 $\times 10^{-6}$	5.09 \times 10 ⁻⁶	6.28 \times 10 ⁻⁶	7.30×10^{-6}
$\mathbf{1}$	1.109	0.884 \Leftarrow	0.9390	1.043
$\mathbf{2}$	1.878	1.282	1.216 ⇚	1.267
$\mathbf{3}$	2.587	1.645	1.471 \Leftarrow	1.467
5	4.052	2.383	1.946	1.841 \Leftarrow
10	7.293	4.068	3.109	2.687 \Leftarrow
		$($ ϵ Minimum)		
		Efficiency, %		
ϵ_i , GeV	5	10	15	20
ϵ_n , m-rad W, MJ	3.58 $\times 10^{-6}$	5.09 x 10^{-6} 6.28 x 10^{-6}		7.30×10^{-6}
$\mathbf{1}$	7.52	9.75 \Leftarrow	8.30	6.85
$\boldsymbol{2}$	8.18	12.52	12.96 \Leftrightarrow	11.31
3	8.72	13.65	14.76 $\pmb{\Leftarrow}$	14.58
5	8.38	15.61	17.27	18.30 ¢
10	8.93	16.38	19.64	22.17 \blacklozenge

Normalized Cost

 \mathbf{c}

 \mathcal{N}

 \mathcal{P}_X

 \mathcal{L}

 $($ Maximum $)$

 $\boldsymbol{9}$

TABLE VII. Normalized Cost & Efficiency of A=210 amu, 4 Beamlet Accelerator System with f=5 hertz, ϵ = 1.5 x 10⁻⁵ m-radiars, $q = +1$, and a Tune of $60^{\circ}-24^{\circ}$, with Superconducting Quads and Amorphous Iron Cores.

Normalized Cost

 $($ ϵ Maximum)

r.

v

N)

Accelerator Cost as a Function of Energy for Various 160 amu Ion Energies

Accelerator Cost as a Function of Energy for Various 190 amu Ion Energies

Figure 4.

Accelerator Cost as a Function of Energy for Various 21.0 amu Ion Energies

1.3 Post-Processor Code

'

,

A post-processor code has been developed to place LIAC'EP result: in an accelerator-target design parameter space. For a given ion species and kinetic energy, the code calculates the range based on the range-energy curves of Bangerter, Mark, and Thiessen.⁽²⁾ The spot radius is calculated from the beam output emittance, angle of convergence and a factor which accounts for the momentum spread and chromatic aberration in the beam and final focussing lenses. The spot size r and the range R formed into a $r^{3/2}$ R parameter and the total beam energy W are used to calculate the target gain based on the Lind1-Mark target gain curves⁽³⁾ Using the results from LIACEP, the fusion power per unit accelerator cost and the ratio of fusion power to power into the accelerator can be determined and compared to other ion species, kinetic energy, and beam energy cases.

The post processor code results can then be plotted on the cost-beam energy parameter space given in Figures 1 through 4. An example is shown in Figure 5 for $A = 210$, $\epsilon = 1.5 \times 10^{-5}$ m-radians, and a spot size of 2 mm radius, where the constant ratio of fusion power to power into the accelerator (nG) of 15 is plotted as a function of ion kinetic energy and beam energy. For the simultaneous satisfaction of $nG = 15$ and yield of 600 MJ, an ion kinetic of about 6 GeV and a beam energy of 4.3 MJ or an ion kinetic energy of 21 GeV and a beam energy of 6.4 MJ is required, with the latter case being the low cost option. The most cost effective case for this example, for a target yield of 600 MJ, is for an ion kinetic energy of about 10 GeV and a beam energy of about 4.5 MJ, which results in an nG of 18.5. Similar plots have been made for other ion species.

1.4 Current Activities

We are currently making more LIACEP runs at an emittance of 3×10^{-5} m-radians in support of the HIFSA project. In addition, the post processor

Figure 5.

A Coupled Accelerator - Target Design Parameter Space for $A = 210$ amu with an unnormalized emittance of 1.5 x 10^{-5} m-radians and a spot radius of 2 mm; target yield of 600 MJ and $\eta G = 15$.

code is being modified for use as a pre-processor code to allow a nore intelligent selection of the LIACEP input parameters that considers the target requirements for a given performance. We are in the process of modifying LIACEP to examine lower initial ion kinetic energies (fronl 50 MeV to 10 MeV). We are also studying the effects of changing the tune on the accelerator cost and efficiency.

References

'

- 1. E. P. Lee, T. J. Fessenden, and L. J. Laslett, "Transportable Charge in a Periodic Alternating Gradient System," (to be published) IEEE Trans. Nucl. Sci. (October, 1985).
- 2. R. 0. Bangerter, J. W-K. Mark, and A. R. Thiessen, Phys. Lett. A, 88, 225 (1982).
- 3. J. D. Lindl and J. W-K. Mark, "Revised Gain Curves for Single-Shell Ion-Beam Targets," 1982 Laser Program Annual Report, Lawrence Livermore National Laboratory, Livermore, CA, 94550, UCRL-50021-82, pp. 3-19 (1983).

This report was done with support from the Department of Energy. Any conclusions or opinions expressed in this report represent solely those of the author(s) and not necessarily those of The Regents of the University of California, the Lawrence Berkeley Laboratory or the Department of Energy.

•

 $\,$, $\,$, $\,$

Reference to a company or product name does not imply approval or recommendation of the product by the University of California or the U.S. Department of Energy to the exclusion of others that may be suitable.

 χ^2

 $\mathcal{L}^{\text{max}}_{\text{max}}$ and $\mathcal{L}^{\text{max}}_{\text{max}}$ and $\mathcal{L}^{\text{max}}_{\text{max}}$

 \sim \sim

LAWRENCE BERKELEY LABORATORY *TECHNICAL INFORMATION DEPARTMENT UNIVERSITY OF CALIFORNIA BERKELEY, CALIFORNIA 94720*

 $\mathcal{L}^{\text{max}}_{\text{max}}$, $\mathcal{L}^{\text{max}}_{\text{max}}$

and the control of the control of

 \mathcal{A}^{max} and \mathcal{A}^{max}

 ~ 100

 α

 $\sim 10^7$