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FINITE ELEMENT FORMULATION
AND SOLUTION OF CONTACT-IMPACT
PROBLEMS IN CONTINUUM MECHANICS

by

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Introduction

In this report we consider the general problem of contact and impact between two bodies. The report is divided into three basic parts. These parts describe: (I) The general theory of contact-impact problems, (II) A numerical scheme for the analysis of contact-impact problems, and (III) The description of computer program FEAP 74 for the solution of contact-impact problems. In an appendix we include the program subroutines and general input description for FEAP 74.

In Sections 1 to 6, Part I, we deal with spatial aspects of the theory and in Section 7, Part I, we deal with temporal aspects. This splitting of the theory is motivated by the way we intend to numerically solve the equations, i.e., the finite element method spatially and a finite difference method temporally.

Part II considers a numerical implementation of the theory given in Part I. Section 9 deals with spatial notions of the numerical problem and Section 10 the temporal. The solution scheme for the resulting algebraic problem is discussed in Sections 11 and 12.

The computer program FEAP 74 was modified to incorporate the numerical contact-impact model. The program modifications and capabilities together with two numerical examples are contained in Part III.

Finally, in the appendix we give listings for the contact subroutines together with the data input instructions.

PART I

VARIATIONAL FORMULATION OF CONTACT-IMPACTPROBLEMS IN CONTINUUM MECHANICS1. Preliminaries

Our conventions on indices are as follows:

Superscripts indicate to which body an entity pertains. Summation is to take place only when explicitly indicated.

Latin subscripts range over 1,2,3, while Greek subscripts range over 1,2. The summation convention is assumed to hold for both.

A body \mathcal{B} is a nice connected region of \mathbb{R}^3 with a piecewise smooth boundary $\partial\mathcal{B}$. A contact^{*} problem is a boundary value problem, or an initial-boundary value problem, in which two bodies, \mathcal{B}^1 and \mathcal{B}^2 , interact according to the principles of mechanics. Thus the primary kinematic axiom of a contact problem is that configurations \mathcal{C}^1 and \mathcal{C}^2 , of \mathcal{B}^1 and \mathcal{B}^2 , respectively, do not penetrate each other, i.e.,

$$\begin{aligned} (\mathcal{C}^1)^\circ \cap \mathcal{C}^2 &= \emptyset, \\ \mathcal{C}^1 \cap (\mathcal{C}^2)^\circ &= \emptyset, \end{aligned} \tag{1}$$

where $(\mathcal{C}^\alpha)^\circ$ denotes the interior of \mathcal{C}^α , $\alpha = 1,2$.

On the other hand the unique condition which characterizes contact problems is that material points on the boundaries of \mathcal{B}^1 and \mathcal{B}^2 may coalesce during the motion of the bodies. Thus we say \mathcal{B}^1 and \mathcal{B}^2 are in contact if $\partial\mathcal{B}^1 \cap \partial\mathcal{B}^2 \neq \emptyset$, and we define the contact surface e by

* It is usual for the term contact to have a static connotation while the term impact has a dynamic connotation. We shall use contact in the general sense to include static as well as dynamic phenomena.

$$\mathfrak{e} = \partial\mathcal{B}^1 \cap \partial\mathcal{B}^2 . \quad (2)$$

If \mathcal{B}^1 and \mathcal{B}^2 are never in contact then $\mathfrak{e} = \emptyset$ for all configurations \mathcal{B}^1 and \mathcal{B}^2 , and in this case an initial-boundary value problem for \mathcal{B}^1 and \mathcal{B}^2 reduces to one in which \mathcal{B}^1 and \mathcal{B}^2 may be treated separately. Thus a non-trivial contact problem is one in which $\mathfrak{e} \neq \emptyset$ for at least one instant during the motion of \mathcal{B}^1 and \mathcal{B}^2 . The picture (Fig. 1) illustrates these notions.

Equation (1) implies that \mathfrak{e} is a material surface with respect to both bodies, i.e., one which is not crossed by material particles. From this we may deduce the interface conditions on \mathfrak{e} .

Let \mathfrak{x} be a persistent point of \mathfrak{e} (one at which joining or releasing of the bodies is not instantaneously occurring) and \mathfrak{v} be the velocity of \mathfrak{x} ($\mathfrak{v} = \dot{\mathfrak{x}}$). Note that only the normal part of \mathfrak{v} is independent of the parametrization of \mathfrak{e} . Let \mathfrak{v}^1 and \mathfrak{v}^2 be the velocities of the material particles located at the points \mathfrak{x}^1 and \mathfrak{x}^2 , contained in $\partial\mathcal{B}^1$ and $\partial\mathcal{B}^2$, respectively, such that $\mathfrak{x} = \mathfrak{x}^1 = \mathfrak{x}^2$ at the present instant. Then since \mathfrak{e} is material and \mathfrak{x} is persistent

$$\mathfrak{v} \cdot \mathfrak{n} = \mathfrak{v}^1 \cdot \mathfrak{n} = \mathfrak{v}^2 \cdot \mathfrak{n} , \quad (3)$$

where \mathfrak{n} is a unit normal vector to \mathfrak{e} at \mathfrak{x} . From this it follows that a necessary condition for momentum to be balanced at \mathfrak{x} is that

$$(\mathfrak{t}^1 + \mathfrak{t}^2) \cdot \mathfrak{n} = \mathfrak{0} , \quad (4)$$

where \mathfrak{t}^m is the Cauchy traction vector with respect to $\partial\mathcal{B}^m$.

In addition we assume that no tensile tractions can occur on \mathfrak{e} ,

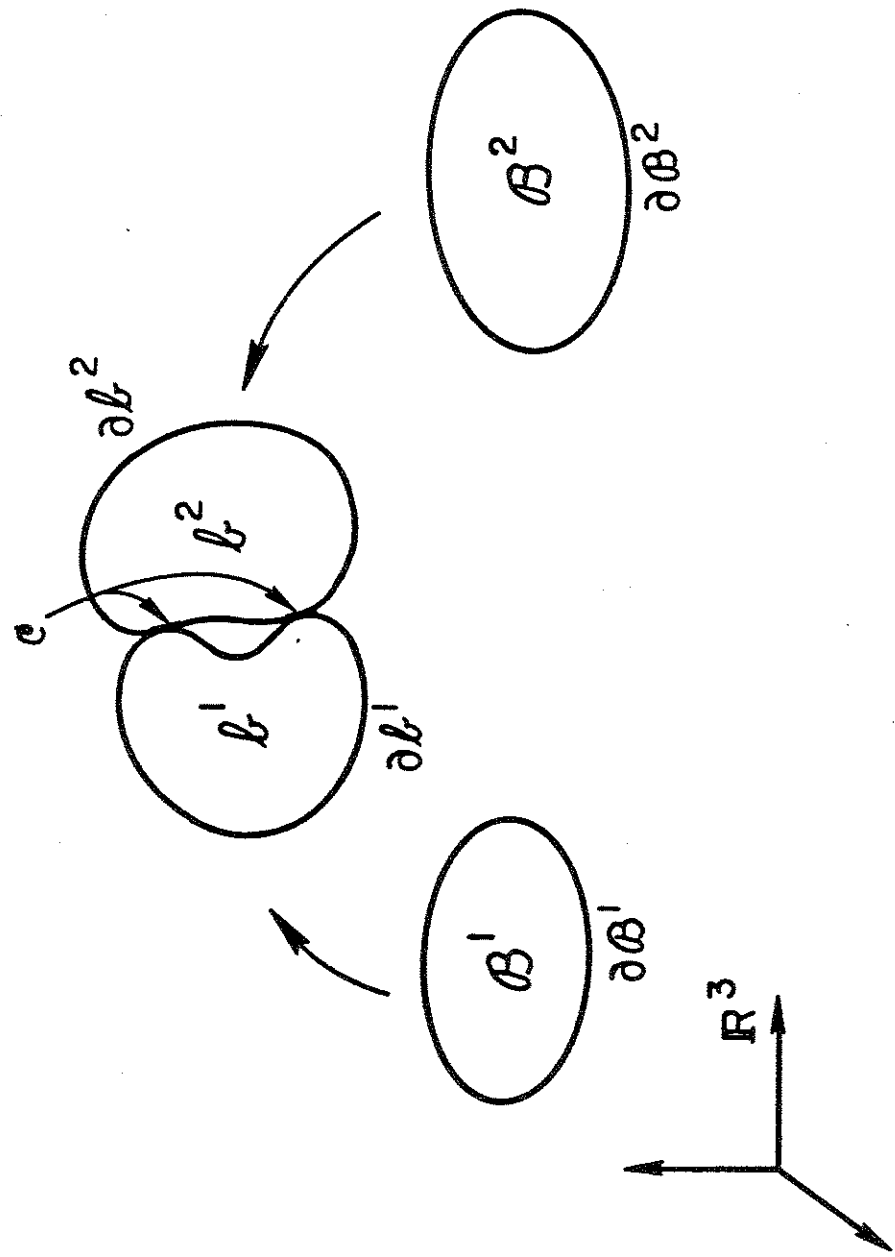


Figure 1

$$\underline{t}^{\alpha} \cdot \underline{n}^{\alpha} \leq 0, \quad (5)$$

where \underline{n}^{α} is the outward unit normal vector to $\partial \mathcal{B}^{\alpha}$. This condition excludes the possibility of the two bodies being glued together. Conditions (1-5) characterize our notion of a contact problem.

Note that thus far we have said nothing about the tangential parts of \underline{v}^{α} and \underline{t}^{α} . These remaining conditions are determined by the frictional nature of the contact. We shall study two simple cases.

Case I: If we assume that points, once in contact, move with \mathcal{C} until released, we have that

$$\underline{v}^1 = \underline{v}^2, \quad (6)$$

and therefore

$$\underline{t}^1 + \underline{t}^2 = \underline{0}. \quad (7)$$

For this model we say that a no-slip, or perfect friction, condition is achieved on \mathcal{C} . Thus condition (5) and equations (6) and (7) are the interface conditions for this case.

Case II: We may create the interface conditions for a frictionless, sliding contact by asserting that the tangential part of each \underline{t}^{α} is identically zero,

$$\underline{t}^{\alpha} - (\underline{t}^{\alpha} \cdot \underline{n}^{\alpha}) \underline{n}^{\alpha} = \underline{0}. \quad (8)$$

Eq. (8), along with (3-5), are the interface conditions for this case.

2. Variational Theorems

We will formulate a variational theorem for the contact problem of finite elastodynamics. We point out, however, that our treatment is entirely general and could be used in conjunction with any field theory, as the only unique feature of the formulation involves the handling of interface conditions. At the same time finite elastodynamics, though lending itself to a clean and simple variational statement, is a case of wide practical interest.

We shall first obtain a variational theorem for the usual initial-boundary value problem of finite elastodynamics by a trivial generalization of some work done by S. Nemat-Nasser [1].

For notational simplicity let \mathcal{Q} denote $\partial\mathcal{B}$, and let $d\mathcal{Q}$ and $d\mathcal{B}$ denote area and volume forms for \mathcal{B} and \mathcal{Q} , respectively. Let $\mathcal{Q}_T \subset \mathcal{Q}$ be that part of \mathcal{Q} where surface tractions are prescribed, and denote by $\bar{\mathbf{T}}$ the Piola - Kirchhoff traction vector representing these prescribed tractions. Call ρ_0 the density of \mathcal{B} in the initial configuration, $\underline{\mathbf{F}}$ the extrinsic body force vector and let $\underline{\mathbf{x}} = \underline{\mathbf{x}}_*(\underline{\mathbf{X}})$ represent the position at time t of the material particle located at $\underline{\mathbf{X}}$ in the initial configuration. For convenience we take \mathcal{B} to be the initial configuration. We denote by $\partial\underline{\mathbf{x}}/\partial\underline{\mathbf{X}}$ the deformation gradients and by $\Phi(\partial\underline{\mathbf{x}}/\partial\underline{\mathbf{X}})$ the strain energy density. Then if $\underline{\mathbf{x}}$ satisfies the kinematic boundary conditions

$$\underline{\mathbf{x}} = \bar{\underline{\mathbf{x}}} \quad (9)$$

on $\mathcal{Q}_x \subset \mathcal{Q}$, where

$$\begin{aligned} \mathcal{Q}_x \cup \mathcal{Q}_T &= \mathcal{Q}, \\ \mathcal{Q}_x \cap \mathcal{Q}_T &= \emptyset, \end{aligned}$$

the functional \mathbb{II} defined by

$$\begin{aligned} \Pi(\underline{x}) = \int_0^t \left\{ \int_{\mathcal{B}} \left(\Phi(\partial \underline{x} / \partial \underline{X}) - \rho_0 \dot{\underline{x}} \cdot \dot{\underline{x}} / 2 \right. \right. \\ \left. \left. - \rho_0 \underline{F} \cdot \underline{x} \right) d\mathcal{B} - \int_{\mathcal{A}_T} \underline{x} \cdot \underline{T} d\mathcal{A} \right\} dt, \end{aligned} \quad (10)$$

is stationary, i.e., its first variation vanishes

$$\begin{aligned} 0 = \delta \Pi(\underline{x}, \delta \underline{x}) = \int_0^t \left\{ \int_{\mathcal{B}} \left(\rho_0 (\ddot{\underline{x}} - \underline{F}) - \text{DIV } \underline{P} \right) \cdot \delta \underline{x} d\mathcal{B} \right. \\ \left. + \int_{\mathcal{A}_T} (\underline{T} - \bar{\underline{T}}) \cdot \delta \underline{x} d\mathcal{A} \right\} dt, \end{aligned} \quad (11)$$

subject to the constraint on variations $\delta \underline{x}_t = \delta \bar{\underline{x}}_t = 0$, ^{on \mathcal{A}_x} (12)
if and only if the equations of motion and traction boundary conditions are satisfied

$$\rho_0 (\ddot{\underline{x}} - \underline{F}) = \text{DIV } \underline{P}, \quad \text{in } \mathcal{B}, \quad (13)$$

$$\underline{T} = \bar{\underline{T}}, \quad \text{on } \mathcal{A}_T, \quad (14)$$

where $\underline{P} = \partial \Phi / \partial (\partial \underline{x} / \partial \underline{X})$ is the first Piola - Kirchhoff stress tensor, $\underline{T} = \underline{N} \cdot \underline{P}$ is the Piola - Kirchhoff traction vector, and \underline{N} is the outward unit normal vector to \mathcal{A} . The solution to the initial-boundary value problem must also satisfy the given initial conditions

$$\left. \begin{aligned} \underline{x} &= \underline{x}_0 \\ \dot{\underline{x}} &= \dot{\underline{x}}_0 \end{aligned} \right\} \text{ in } \mathcal{B}. \quad (15)$$

To interpret this variational theorem for two (non-interacting) bodies set

$$\begin{aligned} \mathcal{B} &= \mathcal{B}^1 \cup \mathcal{B}^2, \\ \mathcal{A} &= \mathcal{A}^1 \cup \mathcal{A}^2, \quad \text{etc.}, \end{aligned}$$

and write

$$\mathbb{I}(\underline{x}) = \mathbb{I}^1(\underline{x}^1) + \mathbb{I}^2(\underline{x}^2) .$$

The next step is to add to \mathbb{I} terms manifesting the interface conditions on \mathcal{C} and to stipulate the constraints under which the vanishing of the first variation of the appended functional corresponds to a solution of the contact problem. To do this we must consider further the kinematics and geometry of \mathcal{C} .

Define two piecewise smooth, invertible maps $\underline{x}^1, \underline{x}^2$ by the condition

$$(\underline{x}^\alpha)^{-1} : \mathcal{C} \longrightarrow \mathcal{C}^\alpha \subset \mathcal{Q}^\alpha , \quad (16)$$

where each \underline{x}^α identifies points on the boundary of the initial configuration \mathcal{B}^α which map into the contact surface \mathcal{C} at each instant of time. If $\underline{x} \in \mathcal{C}$, then $\underline{x}^1 = (\underline{x}^1)^{-1}(\underline{x})$ and $\underline{x}^2 = (\underline{x}^2)^{-1}(\underline{x})$ are the positions of particles in \mathcal{Q}^1 and \mathcal{Q}^2 , respectively, which have coalesced at $\underline{x} \in \mathcal{C}$. It is clear what the \underline{x}^α 's really are, viz., if $\underline{x}^\alpha = \underline{x}_c^\alpha(\underline{X}^\alpha)$, for all $\underline{X}^\alpha \in \mathcal{B}^\alpha$, represents the motion of body \mathcal{B}^α from the original configuration \mathcal{B}^α to the present one \mathcal{B}^α , then \underline{x}^α is the restriction of \underline{x}^α to \mathcal{C}^α ,

$$\underline{x}^\alpha(\underline{x}^\alpha) = \underline{x}^\alpha(\underline{X}^\alpha) , \quad (17)$$

for each $\underline{X}^\alpha \in \mathcal{C}^\alpha$, $\alpha = 1, 2$. For the time being we consider the \underline{x}^α 's as maps defined independently of the \underline{x}^α 's and consider (17) a constraint on possible motions.

We are interested in to what extent the relation

$$\underline{x} = \underline{x}^1(\underline{X}^1) = \underline{x}^2(\underline{X}^2) , \quad (18)$$

is smooth in time and analogously under what circumstances the variations of the \underline{x}^{α} 's are equal. In general the \underline{x}^{α} 's will not even be continuous in time since contact surfaces can be instantaneously created or destroyed. If we eliminate such exceptional instants and consider only persistent points, the bodies still may slide with respect to each other, as depicted in Fig. 2. Thus tangential velocities are seen to be unequal in general. However, when \underline{x} is persistent, the impenetrability condition (1) forces the normal velocity components to be equal, and concomitantly the normal components of variations of the \underline{x}^{α} 's are also equal

$$\delta \underline{x}^1 \cdot \underline{n} = \delta \underline{x}^2 \cdot \underline{n} \quad (19)$$

For sliding contact (Case II), Eq. (19) characterizes the constraint on variations of the \underline{x}^{α} 's equivalent to the velocity constraint (3).

For no-slip contact (Case I),

$$\delta \underline{x}^1 = \delta \underline{x}^2 , \quad (20)$$

is easily seen to be the condition on variations equivalent to Eq. (6). We shall see that Eqs. (19) and (20) lead to the proper interface conditions in the variational theorems.

Introduce vector valued Lagrange multipliers \underline{C}^{α} , and add

$$\mathcal{X} = - \sum_{\alpha=1}^2 \int_0^t \int_{\underline{C}^{\alpha}} \underline{C}^{\alpha} \cdot (\underline{x}^{\alpha} - \underline{x}^{\alpha}) d\underline{C}^{\alpha} dt , \quad (21)$$

to the functional Π (Eq. 10). Note that when $\underline{e} \neq \emptyset$,

$$\underline{a}^{\alpha} = \underline{a}_x^{\alpha} \cup \underline{a}_t^{\alpha} \cup \underline{e}^{\alpha} ,$$

and assume for consistency's sake that

$$\mathcal{C}^- \subset \mathcal{A}_+^-,$$

$$\underline{\mathbb{T}} = \underline{0} \quad \text{on} \quad \mathcal{C}^-. \quad (22)$$

This condition will preclude the ambiguous circumstance of non-zero tractions being specified on the contact area. Upon taking variations of

$\mathcal{J} = \mathbb{I} + \mathcal{X}$ we get Eqs. (13), (14) and,

$$\begin{aligned} 0 = & - \sum_{\alpha=1}^2 \int_0^t \int_{\mathcal{C}^-} \left\{ \delta \underline{\mathcal{C}}^\alpha \cdot (\underline{x}^\alpha - \underline{\mathcal{Z}}^\alpha) + \right. \\ & \left. + \delta \underline{x}^\alpha \cdot (\underline{\mathcal{C}}^\alpha - \underline{\mathbb{T}}^\alpha) - \delta \underline{\mathcal{Z}}^\alpha \cdot \underline{\mathcal{C}}^\alpha \right\} d\mathcal{C}^- dt + \\ & + \text{transversality condition.} \end{aligned} \quad (23)$$

The transversality condition is the classical terminology for variations associated with the domain \mathcal{C}^- .

The first summand of (23) gives us (17) which insures that the \underline{x}^α 's map into \mathcal{C} properly. The second summand identifies $\underline{\mathcal{C}}^\alpha$ as the Piola - Kirchhoff traction vector $\underline{\mathbb{T}}^\alpha$ on \mathcal{C}^- . Let us investigate the third summand.

Consider first Case I and define

$$\delta \underline{\mathcal{Z}} = \delta \underline{\mathcal{Z}}^\alpha, \quad \alpha = 1, 2, \quad (24)$$

which makes sense because of Eq. (20). This condition is equivalent to insisting

$$\dot{\underline{\mathcal{Z}}} = \dot{\underline{\mathcal{Z}}}^\alpha, \quad \alpha = 1, 2,$$

thus the first summand of (23) also implies (6) holds whenever we have a

persistent point. Let j^{α} denote the Jacobian determinant associated with χ^{α} ,

$$de = j^{\alpha} d\mathcal{C}^{\alpha}. \quad (25)$$

Notice then that since $\underline{\mathcal{T}}^{\alpha}$ is the Piola - Kirchhoff traction vector, $(1/j^{\alpha}) \underline{\mathcal{T}}^{\alpha}$ is the corresponding Cauchy traction vector. With these we have for the third summand,

$$0 = \sum_{\alpha=1}^2 \int_{\mathcal{C}^{\alpha}} \delta \chi^{\alpha} \cdot \underline{\mathcal{T}}^{\alpha} d\mathcal{C}^{\alpha} = \int_{\mathcal{E}} \delta \chi \cdot ((1/j^1) \underline{\mathcal{T}}^1 + (1/j^2) \underline{\mathcal{T}}^2) de, \quad (26)$$

which in words means the Cauchy traction vectors are in equilibrium. Thus the momentum balance, Eq. (7), is satisfied on \mathcal{E} .

In Case II we only have that (19) holds, so define

$$\delta \chi^{(\alpha)} = \delta \chi^{\alpha} \cdot \underline{n}, \quad \alpha = 1, 2. \quad (27)$$

This requirement also insures that,

$$\dot{\chi}^1 \cdot \underline{n} = \dot{\chi}^2 \cdot \underline{n},$$

thus the first summand of (23) implies (3). For this case the third summand takes the form,

$$0 = \sum_{\alpha=1}^2 \int_{\mathcal{C}^{\alpha}} \delta \chi^{\alpha} \cdot \underline{\mathcal{T}}^{\alpha} d\mathcal{C}^{\alpha} = \int_{\mathcal{E}} \delta \chi^{(\alpha)} ((1/j^1) \underline{\mathcal{T}}^1 \cdot \underline{n} + (1/j^2) \underline{\mathcal{T}}^2 \cdot \underline{n}) de + \sum_{\alpha=1}^2 \int_{\mathcal{C}^{\alpha}} (\delta \chi^{\alpha} - \delta \chi^{(\alpha)} \underline{n}) \cdot \underline{\mathcal{T}}^{\alpha} d\mathcal{C}^{\alpha}. \quad (28)$$

The integral over \mathcal{E} gives us Eq. (4). The significance of the second integral hinges on the observation that $(\delta \chi^{\alpha} - \delta \chi^{(\alpha)} \underline{n})$ is a tangent vector to \mathcal{C} for each α . Thus the tangential part of each $\underline{\mathcal{T}}^{\alpha}$ is identically zero, which is equivalent to the shear free condition, Eq. (8),

which we require for Case II.

A standard calculation enables us to write the transversality condition as,

$$0 = \sum_{\alpha=1}^2 \int_{\partial C^{\alpha}} (\delta X^{\alpha} \cdot \underline{T}^{\alpha} (\underline{c}^{\alpha} \cdot (\underline{x}^{\alpha} - \underline{x}'^{\alpha}))) d(\partial C^{\alpha}), \quad (29)$$

where the transversal \underline{T}^{α} is a unit vector field tangent to C^{α} , and perpendicular and pointing outward with respect to ∂C^{α} , Fig. 3. Thus (29) implies that

$$\underline{c}^{\alpha} \cdot (\underline{x}^{\alpha} - \underline{x}'^{\alpha}) = 0 \quad \text{on } \partial C^{\alpha}, \quad \alpha=1,2 \quad (30)$$

Assuming continuity of the integrands of (21) on the closure of C^{α} , condition (30) is already implied by the first summand of (23). This assumption precludes \underline{c}^{α} taking the form of a δ -distribution on ∂C^{α} .

Although this assumption is warranted here it may not be true when one employs certain approximate theories in mechanics. For instance consider the case where a Bernoulli-Euler beam is uniformly loaded and sits on a rigid parabolic surface (Fig. 4). At the contact points a, a' , concentrated reactions must exist to balance shear forces. This example is actually from a completely different class of contact problems in that contact is made along a part of the interior rather than the boundary. Such problems as the contact of plates and shells also fall into this class. We could summarize such situations by the description -- m -dimensional contact of m -dimensional bodies, e.g., for the beam $m=1$, and for plates and shells $m=2$. The case under investigation in this paper ($m=3$) is an example of the $(m-1)$ -dimensional contact of m -dimensional bodies.

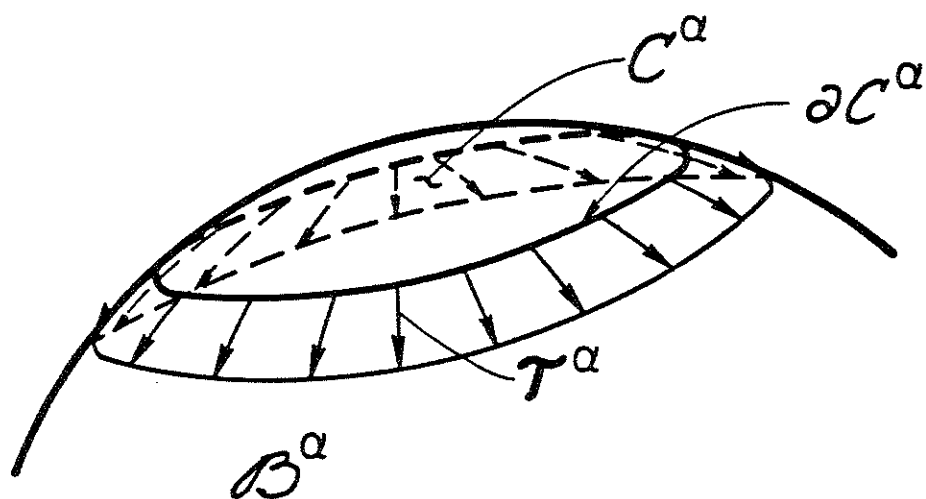


Figure 3

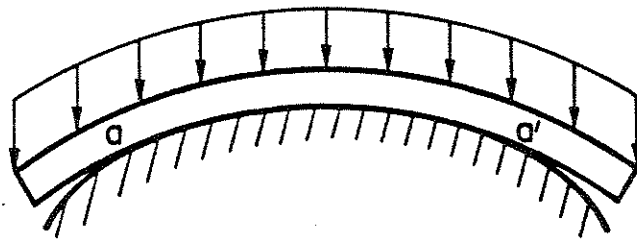


Figure 4

It is good to keep in mind cases such as that illustrated in Fig. 4 when considering specific boundary value problems.

A further point worth mentioning here is that the transversality condition will in general be an independent one in a numerical algorithm. For example, if the fields in the integrand of (21) are approximated by a family of trial functions, Eq. (23) only implies that some weighted integrals over the C^α 's vanish. The condition (29) requires that weighted integrals over the δC^α 's also vanish.

We now summarize our results in the following theorems:

Theorem I: Let (1), (2), (5), (9), (12), (15), and (20) hold. Then x is a solution to the no-slip contact problem (Case I), that is, (6), (7), (13), (14) and (17) also hold if, and only if, $\delta J = 0$ for arbitrary variations of x^α , ψ^α and τ^α , $\alpha = 1, 2$.

Theorem II: Let (1), (2), (5), (9), (12), (15), and (19) hold. Then x is a solution to the sliding contact problem (Case II), that is, (3), (4), (8), (13), (14) and (17) also hold if, and only if, $\delta J = 0$ for arbitrary variations of x^α , ψ^α and τ^α , $\alpha = 1, 2$.

3. Consideration of Theorems I and II as Computational Tools

Theorems I and II may be employed to generate numerical algorithms for the solution of contact problems. The basic idea is to represent \underline{x}^n , $\underline{\chi}^n$ and \underline{c}^n as the product of known functions on \mathbb{R}^3 with unknown parameters depending on time. Then Theorems I and II provide us with a method for generating an approximate system of equations (e.g., by the classical Ritz-Galerkin technique) in terms of these unknown parameters, which then can be solved incrementally and/or iteratively, subject to the side conditions of the theorems. The constraints (1) and (5) will both take the form of inequalities in actual computations, thus the ideas of optimization theory will probably be useful in the actual construction of a numerical algorithm.

The finite element method is a powerful technique for obtaining a system of approximate equations, and it is of interest to find out how amenable are Theorems I and II to a finite element formulation. Unfortunately the term $\underline{\chi}$ would result in a terrible mess if the integrand was represented by typical finite element functions. This is because the boundaries of the \underline{c}^n 's are unknown and thus a parametric integration would bury the defining parameters of the \underline{c}^n 's in the arguments of Heaviside functions representing the supports of the elements. Note that a classical Ritz-Galerkin approximation would not be subject to this pitfall, since the associated trial functions could be chosen to be real analytic and thus easily integrated parametrically to a relatively simple form. However, such a formulation is restricted to a geometrically simpler class of problems. Thus it is desirable to seek a generalization that will lend itself cleanly to a finite element formulation.

4. Variational Theorems Without Transversality Conditions

Let \tilde{E}^n be a fixed part of A_T^n such that

$$\tilde{E}^n \supset E^n, \quad (31)$$

and

$$\bar{I} = 0 \quad \text{on} \quad \tilde{E}^n \sim E^n. \quad (32)$$

Define a scalar valued function n^n on \tilde{E}^n such that

$$n^n(x^n) = 0 \quad \text{if} \quad x^n \in \tilde{E}^n \sim E^n. \quad (33)$$

Let $\tilde{\varepsilon} \supset \varepsilon$, and define the maps \tilde{x}^n by the condition

$$(\tilde{x}^n)^{-1} : \tilde{\varepsilon} \longrightarrow \tilde{E}^n,$$

where, as before, \tilde{x}^n represents x^n on E^n ; but on $\tilde{E}^n \sim E^n$ we place no physical interpretation on \tilde{x}^n . Thus on \tilde{E}^n we will always have that,

$$n^n(\tilde{x}^n - \tilde{x}^n) = 0, \quad (34)$$

since $\tilde{x}^n = \tilde{x}^n$ on E^n and $n^n = 0$ on the relative complement $\tilde{E}^n \sim E^n$.

Introduce vector valued Lagrange multipliers σ^n and let $\mathcal{L} = \mathbb{I} + \mathcal{M}$ where

$$\mathcal{M} = - \sum_{a=1}^2 \int_0^t \int_{\tilde{E}^n} \sigma^n \cdot n^n(\tilde{x}^n - \tilde{x}^n) d\tilde{E}^n dt. \quad (35)$$

We require that the variations of \tilde{x}^n satisfy the same conditions as before, but now for all \tilde{E}^n :

$$\left. \begin{array}{l} \text{Case I:} \quad \delta \tilde{x}^n \stackrel{\text{def.}}{=} \delta \tilde{x}^n \\ \text{Case II:} \quad \delta \tilde{x}^n(n) \stackrel{\text{def.}}{=} \delta \tilde{x}^n \cdot n \end{array} \right\} \quad \text{on} \quad \tilde{E}^n \quad (36)$$

where \underline{n} is a unit normal vector to \tilde{E} . Computing the first variation of \mathcal{L} we have the usual conditions emanating from Π and

$$\begin{aligned}
 0 = & - \sum_{\alpha=1}^2 \int_0^t \int_{\tilde{E}_\alpha} \left\{ \delta \underline{\sigma}^\alpha \cdot (\underline{\eta}^\alpha (\underline{x}^\alpha - \underline{z}^\alpha)) \right. \\
 & + \delta \underline{\eta}^\alpha (\underline{\sigma}^\alpha \cdot (\underline{x}^\alpha - \underline{z}^\alpha)) \\
 & + \delta \underline{x}^\alpha \cdot (\underline{\eta}^\alpha \underline{\sigma}^\alpha - \underline{I}^\alpha) \\
 & \left. - \delta \underline{z}^\alpha \cdot (\underline{\eta}^\alpha \underline{\sigma}^\alpha) \right\} d\tilde{E}^\alpha dt .
 \end{aligned} \tag{37}$$

The first summand gives us (34) and we define

$$\mathcal{C}^\alpha = \{ \underline{x}^\alpha \in \tilde{E}^\alpha : \underline{x}^\alpha(\underline{X}^\alpha) = \underline{z}^\alpha(\underline{X}^\alpha) \} . \tag{38}$$

The third summand defines $\underline{\eta}^\alpha \underline{\sigma}^\alpha$ as the Piola - Kirchhoff traction vector. Note that this insures that $\underline{I}^\alpha = 0$ on $\tilde{E}^\alpha \sim \mathcal{C}^\alpha$ since $\underline{\eta}^\alpha = 0$ there. The fourth summand gives us the appropriate Cauchy traction condition across \mathcal{C} for each case of (36). The second summand is identically satisfied on \mathcal{C}^α since $\underline{x}^\alpha = \underline{z}^\alpha$. On $\tilde{E}^\alpha \sim \mathcal{C}^\alpha$ it tells us that $\underline{\sigma}^\alpha$ is orthogonal to $\underline{x}^\alpha - \underline{z}^\alpha$, but this is of no physical interest.

Thus we can state the following theorems:

Theorem I': Let (1), (2), (5), (9), (12), (15), (31), (32) and (36)₁ hold. Then \underline{x} is a solution to the no-slip contact problem (Case I), that is, (6), (7), (13), (14) and (17) also hold where \mathcal{C}^α is defined by (38), if $\delta \mathcal{L} = 0$ for arbitrary variations of \underline{x}^α , \underline{z}^α , $\underline{\eta}^\alpha$ and $\underline{\sigma}^\alpha$, $\alpha = 1, 2$.

Theorem II': Let (1), (2), (5), (9), (12), (15), (31), (32) and (36)₂ hold. Then \underline{x} is a solution to the sliding contact problem (Case II), that is, (3), (4), (8), (13), (14) and (17) hold where \mathcal{C}^α is defined by (38), if $\delta \mathcal{L} = 0$ for arbitrary variations of \underline{x}^α , \underline{z}^α , $\underline{\eta}^\alpha$ and $\underline{\sigma}^\alpha$, $\alpha = 1, 2$.

The important feature of these theorems is that the regions \bar{E}^n are fixed. Thus transversality conditions are absent, and the theorems may be applied to finite element formulations. In fact one would naturally take \bar{E}^n to be a union of elements in \mathcal{a}^n , large enough to contain E^n throughout the motion.

Thus far our considerations have been quite general and, in fact, more general than would be required for the solution of particular classes of contact problems. In the next section we illustrate the many simplifications which can be made in the application of the preceding theorems to a class of problems of wide practical interest.

5. Hertzian Contact Problems

We wish to characterize contact problems in which the contact surface is approximately planar and the bodies have undergone small deformations in the neighborhood of the contact surface.

Assume the following:

(1) $\underline{n} \stackrel{\text{def}}{=} n_i \underline{e}_i \approx \underline{e}_3$ on \mathcal{C} , where the n_i indicate components with respect to the standard basis $\{\underline{e}_i\}_1^3$ for \mathbb{R}^3 ,

(see Fig. 5).

(2) $j^\alpha \approx 1$, $\alpha = 1, 2$, thus $\underline{t}^\alpha \approx \underline{T}^\alpha$ on \mathcal{C}^α .

Assumptions (1) and (2) together imply that,

$$t_3^\alpha \approx \underline{t}^\alpha \cdot \underline{n} \approx \underline{T}^\alpha \cdot \underline{n} \approx T_3^\alpha, \quad ,$$

and that,

$$(\underline{t}_1^\alpha, \underline{t}_2^\alpha, 0) \approx \underline{t}^\alpha - (\underline{t}^\alpha \cdot \underline{n}) \underline{n} \approx \underline{T}^\alpha - (\underline{T}^\alpha \cdot \underline{n}) \underline{n} \approx (T_1^\alpha, T_2^\alpha, 0).$$

(3) Material points which eventually contact have, to the first order, the same initial coordinates \underline{z}_1 and \underline{z}_2 . Explicitly we manifest this idea by requiring that the χ^α 's satisfy

$$\chi^1(\underline{z}_1, \underline{z}_2, X_3^1(\underline{z}_1, \underline{z}_2)) = \chi^2(\underline{z}_1, \underline{z}_2, X_3^2(\underline{z}_1, \underline{z}_2)). \quad (39)$$

This is depicted in Fig. 6. Since X_3^α are given functions which define the surfaces \mathcal{C}^α , it follows from (39) that,

$$\delta \chi^1 = \delta \chi^2.$$

We term problems for which these assumptions hold Hertzian, since these assumptions are implicit in Hertz' classical theory [2] (see

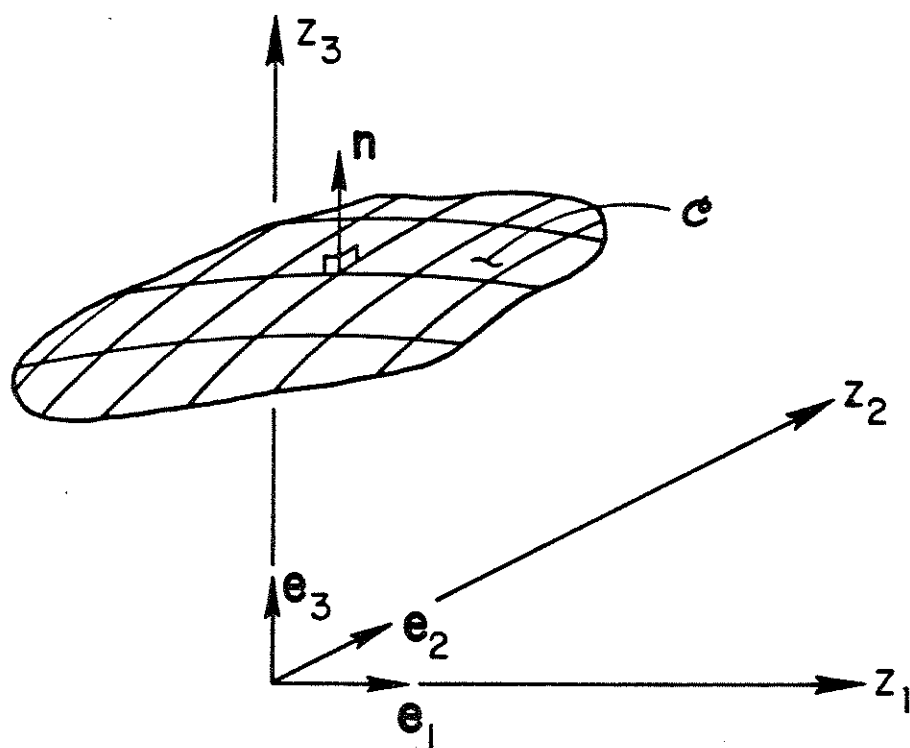


Figure 5

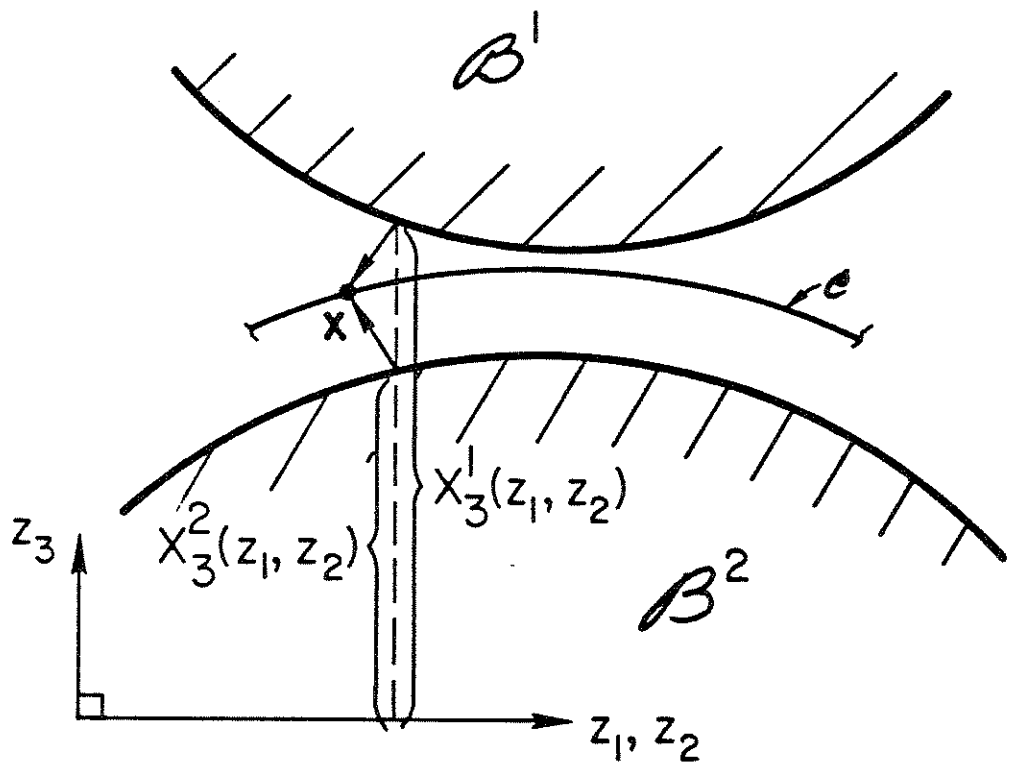


Figure 6

Goldsmith [3] for an excellent exposition of this work and also many applications of Hertz' theory to impact problems). It should be pointed out that the formulation we are about to give is still considerably more general than those to which Hertz' theory applies.

We now show how these assumptions allow us to make simplifications in the preceding theorems.

Theorems I and II:

Due to assumption (3) the term \mathcal{K} can be replaced by an integral over a region in the z_1, z_2 -plane. This region, say c , is the projection of \mathcal{C} onto the z_1, z_2 -plane, and due to assumption (2) it coincides, to the first order, with the projections of the \mathcal{C}^n 's. Thus \mathcal{K} can be written

$$\mathcal{K} = - \sum_{n=1}^2 \int_0^t \int_c \underline{\tau}^n \cdot (\underline{x}^n - \underline{z}^n) \, dc \, dt . \quad (40)$$

Since, for Case I, we know that the momentum balance on \mathcal{C} requires that

$$\underline{\tau}^1 + \underline{\tau}^2 = \underline{0} ,$$

we may make use of this relation immediately. Thus define

$$\underline{\tau} = \underline{\tau}^1 = -\underline{\tau}^2 ,$$

and substitute into (40). Employing (39), the integrand simplifies to

$$\underline{\tau} \cdot (\underline{x}^2 - \underline{x}^1) . \quad (41)$$

The analog of (23) becomes

$$\begin{aligned}
0 = \int_0^t \int_C \left\{ \delta \tau \cdot (\underline{x}^2 - \underline{x}^1) + \right. \\
\left. + \delta x^1 \cdot (\underline{T}^1 - \tau) + \delta x^2 \cdot (\underline{T}^2 + \tau) \right\} dc dt \\
(42) \\
+ \text{transversality condition.}
\end{aligned}$$

Thus the same conclusions of Theorem I can be drawn. However, from a numerical standpoint things are considerably different. First of all, since the \underline{x}^{α} 's are absent in this formulation, we do not get a uniquely defined τ ; \underline{x}^1 and \underline{x}^2 will not in general be the same pointwise. If the graph of τ is important it could be constructed by averaging \underline{x}^1 and \underline{x}^2 , which, if the solution is any good, should be reasonably close pointwise. On the other hand, the \underline{x}^{α} 's being absent engenders a considerable saving in the number of equations to be solved and in their complexity.

The analogous case for Theorem II is constructed simply by setting

$$\tau_1 = \tau_2 = 0, \quad \tau \stackrel{\text{def.}}{=} \tau_3.$$

Then the integrand of \mathcal{K} becomes

$$\tau (x_3^2 - x_3^1) \quad (43)$$

and (42) reduces to

$$\begin{aligned}
0 = \int_0^t \int_C \left\{ \delta \tau (x_3^2 - x_3^1) + \right. \\
+ \delta x_3^1 (T_3^1 - \tau) + \delta x_3^2 (T_3^2 + \tau) \\
+ \delta x_\alpha^1 T_\alpha^1 + \delta x_\alpha^2 T_\alpha^2 \left. \right\} dc dt \\
(44) \\
+ \text{transversality condition.}
\end{aligned}$$

Hence the conclusions of Theorem II hold.

Thus in the case of Hertzian contact we can add the simplifications manifested in (41) and (43) to the conditions of Theorems I and II, respectively, and still garner the same conclusions.

Theorems I' and II':

For these cases \mathcal{M} can be written as an integral over $\tilde{\mathcal{E}}$, the projection of \mathcal{E} :

$$\mathcal{M} = - \sum_{\alpha=1}^2 \int_0^t \int_{\tilde{\mathcal{E}}} \underline{\sigma}^{\alpha} \cdot \underline{n}^{\alpha} (\underline{x}^{\alpha} - \underline{x}^{\alpha}) d\tilde{\mathcal{E}} dt .$$

Due to the present geometric situation, it is appropriate to take

$$\underline{n}^1 = \underline{n}^2 ,$$

and thus define

$$\underline{n} = \underline{n}^{\alpha} , \quad \alpha = 1, 2 .$$

Analogous to the considerations for Theorems I and II, the momentum balance across \mathcal{E} motivates the simplification

$$\underline{\sigma}^{\text{def.}} \underline{\sigma}^1 = - \underline{\sigma}^2 .$$

With these and (39), the integrand of \mathcal{M} can be written

$$\underline{\sigma} \cdot \underline{n} (\underline{x}^2 - \underline{x}^1) .$$

A further simplification can be made by setting*

$$\underline{\sigma}_3 = - \underline{n} .$$

This eliminates one unknown function and, as we shall see, has the effect of satisfying (5) naturally. Thus the integrand of \mathcal{M} becomes

* This is a standard ploy of optimization theory, see p. 82, [4].

$$\sigma_{\alpha} \eta (x_{\alpha}^2 - x_{\alpha}^1) - (\eta)^2 (x_3^2 - x_3^1) , \quad (45)$$

and the analog of (23) is

$$\begin{aligned} 0 = \int_0^t \int_{\mathcal{E}} \{ & \delta \eta (\sigma_{\alpha} (x_{\alpha}^2 - x_{\alpha}^1) - 2\eta (x_3^2 - x_3^1)) \\ & + \delta \sigma_{\alpha} (\eta (x_{\alpha}^2 - x_{\alpha}^1)) \\ & + \delta x_{\alpha}^1 (T_{\alpha}^1 - \eta \sigma_{\alpha}) + \delta x_{\alpha}^2 (T_{\alpha}^2 + \eta \sigma_{\alpha}) \\ & + \delta x_3^1 (T_3^1 + (\eta)^2) + \delta x_3^2 (T_3^2 - (\eta)^2) \} d\mathcal{E} dt . \end{aligned} \quad (46)$$

Summand two tells us that either $\eta = 0$ or $x_{\alpha}^1 = x_{\alpha}^2$, on \mathcal{E} . Suppose $\eta \neq 0$, then $x_{\alpha}^1 = x_{\alpha}^2, \alpha=1,2$. Summand one then gives us that $x_3^1 = x_3^2$ on \mathcal{E} . Thus we have

$$\eta (x^2 - x^1) = 0 , \quad \text{on } \mathcal{E} ,$$

as required, and \mathcal{E} is defined as the subset of \mathcal{E} where $x^1 = x^2$.

The last four summands give the momentum balance conditions, as usual, and, in addition, the last two summands imply that the normal tractions are compressive (since $(\eta)^2 \geq 0$). Thus we have the conclusions of Theorem I' and condition (5).

The analogous set up for Theorem II' is accomplished by setting $\sigma_{\alpha} = 0$ in (45) yielding

$$-(\eta)^2 (x_3^2 - x_3^1) \quad (47)$$

for the integrand of \mathcal{M} . With this Eq. (46) becomes

$$\begin{aligned}
0 = \int_0^t \int_{\bar{c}} \{ & -2 \delta \eta (\eta (x_s^2 - x_s^1)) + \\
& + \delta x_a^1 T_a^1 + \delta x_a^2 T_a^2 \\
& + \delta x_s^1 (T_s^1 + (\eta)^2) + \delta x_s^2 (T_s^2 - (\eta)^2) \} d\bar{c} dt.
\end{aligned}$$

In this case we achieve the conclusion of Theorem II' and condition (5).

Thus to Theorems I' and II' we can delete condition (5), add the simplifications manifested in (45) and (47), and achieve the conclusions of Theorems I' and II', respectively, plus condition (5).

6. Contact Problems for One, Two and Three-dimensional Bodies

The previous work needs only trivial modification to be made applicable to contact problems involving bodies of different dimensions. There are many cases of considerable interest which fall into this category. For example, models consisting of a shell and a plate, or a solid and a plate, are useful for the study of head impact. The modifications necessary are essentially interpretative. An example illustrates this assertion.

Consider the frictionless Hertzian contact of a three-dimensional solid and a two-dimensional plate. Let \mathcal{B}^1 represent the solid and \mathcal{B}^2 the plate. In evaluating Π , the \mathcal{B}^1 part is as before while the \mathcal{B}^2 part would manifest the particular plate theory used. The contact term \mathcal{K} (or \mathcal{M}) would be exactly as before. However note that c (or \bar{c}) is, in this case, also identifiable with part of the two-dimensional "volume" of the plate, rather than its boundary. Taking variations, everything is as before except that the term $\tau \delta x_3^2$ (or $-(\tau)^2 \delta x_3^2$) contributes to the transverse momentum equation of the plate, rather than to its boundary conditions. The interpretation of τ (or $-(\tau)^2$) is thus two-fold, i.e., it is the normal component of the traction vector with respect to \mathcal{B}^1 , as before, and it is also the equivalent normal "body force" with respect to \mathcal{B}^2 , manifested by the interaction with \mathcal{B}^1 (Fig. 7).

This interpretation is general, namely, for one and two-dimensional bodies the contact force is an equivalent "body force" which contributes to the momentum equations, rather than the boundary conditions. With this interpretation in mind, the construction of variational theorems, analogous to the ones constructed in Sections 2, 4 and 5, for the class of one, two and three-dimensional contact problems, is just a formal deductive exercise involving only appropriate definitions for Π .

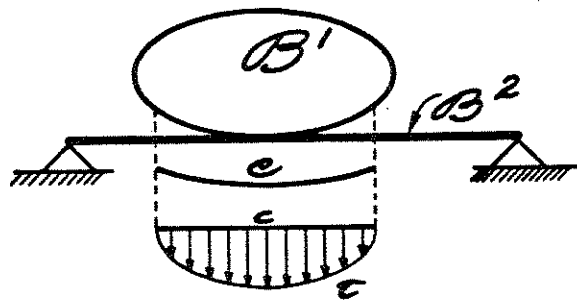


Figure 7

7. Impact

The previous sections deal with spatial aspects of contact problems. In this section we investigate temporal considerations, i.e., those phenomena which are unique to dynamic contact or impact. To manifest the problem encountered in such situations consider the following hypothetical situation. Assume that we are in the process of numerically solving some impact problem and suppose that it is discovered as we monitor the motion of the bodies that they impact somewhere in the time interval (t_1, t_2) . At time t_1 we know the states of both bodies and we know that somewhere between t_1 and t_2 they have coalesced over a portion of their boundaries. Assume for the moment we know the geometry of the contact surface e . The question which arises then is what is the state of e at time t_2 , i.e., what are the velocity and traction vectors on e ? It is necessary to know this information to carry forth the step forward time integration. The question though seems improperly posed without specifying considerable data about the nature of the impact. To get a handle on things, we will initially formulate a simple one-dimensional problem involving the impact of two elastic rods. Although this problem is trivial, it provides considerable insight into the general nature of impact of continuum bodies. Since we are interested in the state of e (in this case a point) immediately after impact, whether the rods are finite or semi-infinite is immaterial.

Assume that the pre-impact states of the two bodies are given by the following data:

$$\mathcal{B}^{\alpha} : \quad \underline{v}^{\alpha}, \quad (\partial x / \partial X)^{\alpha}, \quad P^{\alpha}; \quad \alpha = 1, 2. \quad (48)$$

At impact the rods coalesce at t , and for some finite time interval thereafter (at least) $x \in t$ is persistent. At the moment of impact shock waves begin to propagate in each body. The space-time picture is depicted in Fig. 8. As discussed in section 1, since t is material and x is persistent, we have

$$v \stackrel{\text{def.}}{=} v_+^1 = v_+^2, \quad P \stackrel{\text{def.}}{=} T_+^1 = -T_+^2, \quad (49)^*$$

for the post-impact state (t_+). In addition to (49), the well known shock conditions must hold across the wave fronts:

$$\begin{aligned} [v^\alpha] + U^\alpha [(\partial x / \partial X)^\alpha] &= 0, \\ \rho_0 U^\alpha [v^\alpha] &= [P^\alpha], \end{aligned} \quad (50)$$

where U^α is the material velocity of the shock in B^α , and $[\]$ is the wave-front jump operator which assigns to a function the difference in its values behind and in front of the wave, i.e., $[f(X,t)] = f(X^-,t) - f(X^+,t)$ where X is a material point denoting the location of the wave-front. As can be deduced from Fig. 8, the states into which the shocks initially propagate are the pre-impact states given by (48), and the state at t , immediately after the shocks pass, is given by the post-impact state (49). These observations in conjunction with (50) yield,

$$\begin{aligned} v_-^\alpha - v + U^\alpha \{ (\partial x / \partial X)_-^\alpha - (\partial x / \partial X)_+^\alpha \} &= 0, \\ \rho_0 U^\alpha (v_-^\alpha - v) + P_-^\alpha - P &= 0. \end{aligned} \quad (51)^{**}$$

*For convenience we choose the initial state to be the pre-impact state, thus we need not distinguish between Cauchy and Piola tractions.

**A consistency condition for these equations is that $v_-^1 - v_-^2 > 0$. Otherwise the impact would not occur.

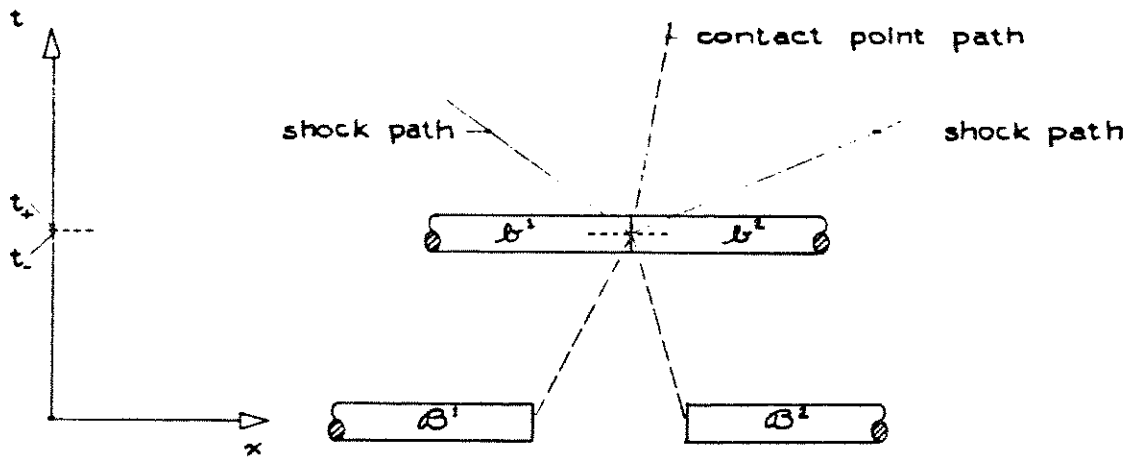


Figure 8

The four Eqs. (51) and constitutive equations relating P^α to $(\partial x / \partial X)^\alpha$ yield a formally deterministic system of six equations in the six unknowns v , P , U^α , $(\partial x / \partial X)^\alpha$. Thus we see that the desired quantities v and P depend on the pre-impact states and material properties of both B^α . The precise form of this relationship depends upon the constitutive equations of the bodies. As a simple example, assume we have linear constitutive equations $P^\alpha = E^\alpha \{ (\partial x / \partial X)^\alpha - 1 \}$, E^α constant, and let the pre-impact state be given by

$$\begin{aligned} v_-^\alpha &= V^\alpha, \\ (\partial x / \partial X)_-^\alpha &= 1, \\ P_-^\alpha &= 0. \end{aligned} \quad (52)$$

These conditions, when inserted in Eqs. (51), lead to:

$$\begin{aligned} v &= \frac{e_0^2 U^2 V^2 - e_0^1 U^1 V^1}{e_0^2 U^2 - e_0^1 U^1}, \\ P &= \frac{V^2 - V^1}{\left(\frac{U^2}{E^2} - \frac{U^1}{E^1} \right)}, \\ (U^\alpha)^2 &= E^\alpha / e_0^\alpha. \end{aligned} \quad (53)$$

Note that the denominators in Eqs. (53)_{1,2} present no problems since $e_0^\alpha > 0$, $E^\alpha > 0$ and $+1 = \text{sgn } U^2 = -\text{sgn } U^1$.

This result is also appropriate whenever the intensity of the impact is small enough such that the non-linear constitutive equation can be replaced by its linear approximation about the pre-impact state. In this case E^α is a tangent modulus evaluated at the pre-impact strain

$\{ (\partial x / \partial X)_-^\alpha - 1 \} = 0$. To further simplify, consider the case when both rods have identical properties (i.e., $e_0 = e_0^\alpha$, $E = E^\alpha$, $\alpha = 1, 2$). Then

$$\begin{aligned}
 v &= \frac{V^1 + V^2}{2} \quad , \\
 P &= \rho U (V^2 - V^1) / 2 \quad , \\
 (U)^2 &= E / \rho_0 .
 \end{aligned}
 \tag{54}$$

In Eqs. (54) U is positive, and since consistency requires $V^1 - V^2 > 0$, P is compressive.

Thus for the one-dimensional case at least the problem of computing the post-impact state is easily achieved. The solution of (51) for the fully non-linear case can be automated as part of a numerical algorithm. Although this problem is trivial, it serves to indicate that the post-impact problem, the solution of which is essential in a numerical algorithm, is one of wave propagation.

In the analysis of higher dimensional bodies the solution of the post-impact problem becomes greatly complicated due to the geometric variety of impact conditions. However, considerable simplifications can be taken advantage of if one keeps in mind the nature of the discrete problem. For instance, if a certain portion of the boundaries of two bodies have coalesced in ϵ , each interior point of ϵ , at which the tangent plane is well defined, may be treated, to the first order, as a point on the mating surface of two impacting half-spaces. As long as time steps are kept small enough, the local behavior is well represented. The post-impact problem for the general case, analogous to (51), can be automated as part of the numerical algorithm, and for many simple cases can be solved explicitly.

With these notions in mind, let us return to the case of main interest in this report, namely three-dimensional continuum bodies. We shall consider only the case of a frictionless contact surface (Case II), and leave the solution of the post-impact problem for the no-slip case (Case I), which is more difficult, for future work. With the proper interpretations, the one-dimensional rod formulation (Eqs. (48-54)) suffices to completely characterize this case. This is so because no tangential motions or stresses may be communicated across a frictionless surface, and thus we need only consider the configuration of normal incidence. In this case the requisite constitutive functionals in (51) would be those relating P^α , the normal Piola stress, to the normal component of strain, holding all other components of strain fixed at the pre-impact values. For example, in the linear isotropic case, E^α (Young's modulus) in Eqs. (53,54) would be replaced by $\lambda^\alpha + 2\mu^\alpha$ ($\lambda^\alpha, \mu^\alpha$ are the Lamé and shear moduli, respectively) and the propagation velocity would be that of dilatational waves.

PART II

A NUMERICAL SCHEME FOR ANALYSIS OF
CONTACT-IMPACT PROBLEMS8. Numerical Solution of Contact-Impact Problems

In performing numerical computations based on the above described variational formulation for contact-impact problems we have employed three distinct levels of approximation: (1) a spatial discretization of the bodies and contact surfaces, (2) a temporal discretization to determine the response of the discretized bodies, and (3) a numerical solution for the resulting system of nonlinear algebraic equations.

In the following sections we shall restrict our attention to the Hertzian contact problem described in Section 5. Significant numerical difficulties are encountered in the solution of impact problems; to complicate the problem further by introducing the additional steps necessary to determine the contact surface maps for the full kinematically nonlinear case is left for a future study. While this is a simple impact problem in terms of determining the contact surface and the full power of the preceding theory is neither necessary nor exploited in its solution, many of the features of the general problem are employed here.

9. Spatial Discretization of the Bodies and Contact Surface

The bodies \mathcal{B}^1 and \mathcal{B}^2 are discretized using standard finite element methods, (e.g., see [5]). In order to facilitate the computation of a discrete Hertzian contact surface the nodes of \mathcal{B}^1 are arranged so that they align with the nodes of \mathcal{B}^2 . This is consistent with the notions of condition 3 of Section 5 and ensures that during determination of the approximation to the contact surface contiguous nodes of the two bodies will meet. Thus, the simulation of the contact surface is trivial. The development of a numerical model for Hertzian contact problems is based upon the form of Theorem II' which uses (47) for the integrand of \mathcal{M} . For numerical computations we introduce the displacement vector \underline{u} such that

$$\underline{x} = \underline{X} + \underline{u} \quad (55)$$

For a compatible finite element displacement field the integrand of \mathcal{M} can be approximated by taking $\mathcal{M}^2(\underline{x}, t)$ as the product of $E^2(t)$ and $\delta(\underline{x} - \underline{x}_i)$ (i.e., Dirac delta functions in space). This corresponds to taking \mathcal{M}^2 as "concentrated nodal loads" which are the generalized forces of the contact pressure. With this discretization we can describe pseudo contact elements between each pair of candidate contact nodes. Let these nodes be denoted as $(\cdot)_i^{\leftarrow}$ and the generalized force as $(E_i)^{\leftarrow}$; then

$$\mathcal{M} = \int_0^t \sum_i (E_i(t))^2 (u_{z_i}^2(t) - u_{z_i}^1(t) + X_{z_i}^2 - X_{z_i}^1) dt \quad (56)$$

where $\{i\}$ are the set of candidate contact nodes which span $\tilde{\mathcal{E}}$; $u_{z_i}^{\leftarrow}$ are the nodal displacements in the z_3 direction and $X_{z_i}^{\leftarrow}$ are the nodal coordinates of the candidate contact nodes.*

*We assume here that 3 is the direction nominally normal to the contact surfaces, e.g., see Fig. 5.

Use of the finite element method in Theorem II' with \mathcal{M} given by (56) produces a set of nonlinear second order ordinary differential equations which together with the impenetrability conditions define the discretized contact impact problem. These equations take the form:

$$\underline{\underline{M}} \ddot{\underline{u}} + \underline{\underline{K}}(\underline{u}) = \underline{\underline{R}} \quad , \quad (57)$$

where $\underline{\underline{M}}$ is the usual finite element mass matrix, $\underline{\underline{K}}$ represents the elastic stiffness forces together with the contact terms, $\underline{\underline{R}}$ is the set of generalized forces resulting from boundary tractions and \underline{u} is the set of time dependent nodal displacements (which also include the $(\epsilon_i)^2$). For inelastic materials Theorem II' can be extended by treating the first variation as a Galerkin method (principle of virtual work) and replacing the elastic constitutive model by more general theories, e.g., viscoelastic, elastoplastic, viscoplastic, etc. In this case

$$\underline{\underline{K}}(\underline{u}) \rightarrow \underline{\underline{K}}(\underline{u}, \dot{\underline{u}}) \quad (58)$$

in (57).

10. Temporal Discretization

A temporal discretization of the second order ordinary differential equations which result from a finite element spatial discretization of the contact-impact problem is accomplished herein by using the Newmark family of methods [6]. The Newmark family of methods is a one-step integration method with two free parameters which can be used to control stability and numerical damping. The method is essentially a difference method in time. The behavior of the method for linear elasto-dynamics problems is discussed in [6,7]. The algorithm is given by

$$\underline{u}_{n+1} = \underline{u}_n + \Delta t \dot{\underline{u}}_n + \left(\frac{1}{2} - \beta\right) \Delta t^2 \ddot{\underline{u}}_n + \beta \Delta t^2 \ddot{\underline{u}}_{n+1}, \quad (59)$$

$$\text{and } \dot{\underline{u}}_{n+1} = \dot{\underline{u}}_n + (1 - \gamma) \Delta t \ddot{\underline{u}}_n + \gamma \Delta t \ddot{\underline{u}}_{n+1},$$

where $\underline{u}_n = \underline{u}(t_n)$, $\Delta t = t_{n+1} - t_n$, and β, γ are the two parameters. For linear problems $\gamma = .5 + \delta = .5$ produces no artificial viscosity and $\beta \geq \frac{1}{4} (1 + \delta)^2$ produces unconditional stability (i.e., the method is stiffly stable). Such generalization is not possible for nonlinear problems and during solution it may be necessary to monitor the solution for any signs of instability. In (59) $\beta = 0$ produces an explicit method for \underline{u}_{n+1} , and if \underline{M} is diagonal (lumped mass) with \underline{K} and \underline{R} independent of $\dot{\underline{u}}$ the solution can be advanced without solving a large set of simultaneous equations; for all other cases the method is implicit and equations must be solved. Solution of (59)₁ for $\ddot{\underline{u}}_{n+1}$ in terms of the solution at t_n and \underline{u}_{n+1} gives

$$\ddot{\underline{u}}_{n+1} = \frac{1}{\beta \Delta t^2} (\underline{u}_{n+1} - \underline{u}_n) - \frac{1}{\beta \Delta t} \dot{\underline{u}}_n - \left(\frac{1-2\beta}{2\beta}\right) \ddot{\underline{u}}_n \quad (60)$$

which can also be used in $(59)_2$ to express the velocity in terms of the solution at t_n and \underline{u}_{n+1} . Since in this process we divide by β and Δt it is no longer possible to consider zero β or zero time steps.

11. Solution of the Nonlinear Algebraic Problem

Use of the Newmark method in (57) (including (58)) yields the set of nonlinear algebraic equations:

$$\frac{1}{\beta \Delta t^2} \underline{M} \underline{u}_{n+1} + \underline{K}(\underline{u}_{n+1}, \underline{u}_n, \dot{\underline{u}}_n, \ddot{\underline{u}}_n) = \underline{R}_{n+1} + \underline{M} \underline{A}_n, \quad (61)$$

where

$$\underline{A}_n = \frac{1}{\beta \Delta t^2} \underline{u}_n + \frac{1}{\beta \Delta t} \dot{\underline{u}}_n + \left(\frac{1-2\beta}{2\beta} \right) \ddot{\underline{u}}_n.$$

A Newton-Raphson iterative solution to this set of equations can formally be constructed, giving:

$$\left(\frac{1}{\beta \Delta t^2} \underline{M} + \partial_u \underline{K} - \partial_u \underline{R} \right) \Delta \underline{u}^{(i)} = \underline{R} - \underline{K}(\underline{u}_{n+1}^{(i)}, \underline{u}_n, \dot{\underline{u}}_n, \ddot{\underline{u}}_n) - \underline{M} \ddot{\underline{u}}_{n+1}^{(i)}, \quad (62)$$

where $\partial_u \underline{R}$ is the effect of loads varying with the deformation and

$$(\partial_u \underline{K})_{ij} = \partial K_i / \partial u_j, \quad (63)$$

is the tangent stiffness matrix. The coefficient to $\Delta \underline{u}^{(i)}$ is generally called the Jacobian matrix of the Newton-Raphson iteration. The solution is advanced by taking

$$\underline{u}_{n+1}^{(i+1)} = \underline{u}_{n+1}^{(i)} + \Delta \underline{u}^{(i)}, \quad (64)$$

and iterating until a norm of the solution satisfies

$$\|\Delta \underline{u}^{(i)}\| \leq \epsilon \|\underline{u}_{n+1}^{(i)}\|, \quad (65)$$

where ϵ is some small positive error tolerance. In the work reported here the norm $\| \cdot \|$ is taken as the Euclidian norm

$$\| \underline{x} \| = \left(\sum_i x_i^2 \right)^{1/2}, \quad (66)$$

and the load vector \underline{R} is assumed to be independent of \underline{u} . For stable elastic materials the resulting tangent stiffness is then symmetric and positive definite, consequently, standard direct solution methods normally employed in the solution of linear finite element problems can be used. For inelastic materials or deformation dependent loads the tangent stiffness may be asymmetric. In these cases some special methods may be necessary to effect a solution.

12. Discretized Impact Conditions

In the previous numerical development \tilde{c} has been defined by discrete points which correspond to nodes along the boundaries of \mathcal{B}^1 and \mathcal{B}^2 . When, during the course of advancing the solution in time, any one of these points violates the impenetrability condition a re-solution must be obtained in which the $(\epsilon_i)^2$ are now non-zero and the u_3^* satisfy the impenetrability condition. Some control and monitoring are required to effect this in a computer program. In addition to satisfying these conditions, the impact relations denoted in Section 7 must be invoked. In the present study these conditions are applied to the solution at the end of a time step in which points first go into contact. Accordingly we compute from (50)*

$$\dot{u}_+ = \frac{e_o^2 U^2 \dot{u}_-^2 - e_o^1 U^1 \dot{u}_-^1}{e_o^2 U^2 - e_o^1 U^1} \quad , \quad (67)$$

and assign this value to the appropriate node of \mathcal{B}^1 and \mathcal{B}^2 .

To determine the solution vector \underline{u} at t_{n+1} , we have solved the set of equations (61). As described above the shock conditions are then used to determine the value of the velocity at time t_{n+1} for all points which have come into contact during the time interval. In order to get a consistent solution at these points we must modify the accelerations and contact force to reflect the shock conditions. This is accomplished by re-solving the equilibrium conditions of \mathcal{B}^1 and \mathcal{B}^2 at point i . The expanded forms of the appropriate equations are:

*The $()_-$ denotes a value which is computed before impact, whereas $()_+$ denotes the value after impact.

$$M^1 \ddot{u}_-^1 + K^1(\underline{u}) + (\mathcal{E}_i)_-^2 = R^1, \quad (68)$$

and

$$M^2 \ddot{u}_-^2 + K^2(\underline{u}) + (\mathcal{E}_i)_-^2 = R^2.$$

For nodes which have come into contact we must enforce the condition on acceleration

$$\ddot{u}_+^1 = \ddot{u}_+^2 = \ddot{u}_+, \quad (69)$$

and compute the contact force $(\mathcal{E}_i)_+^2$. The solution for these is obtained from

$$M^1 \ddot{u}_+ + K^1(\underline{u}) + (\mathcal{E}_i)_+^2 = R^1,$$

and

$$M^2 \ddot{u}_+ + K^2(\underline{u}) + (\mathcal{E}_i)_+^2 = R^2.$$

These are two equations in two unknowns which can be solved for the \ddot{u}_+ and $(\mathcal{E}_i)_+^2$. If $K^{\alpha}(\underline{u})$ is independent of velocity the stiffness forces and R^{α} will remain unchanged during the impact, hence we can solve the simpler problem

$$M^1 \ddot{u}_+ + (\mathcal{E}_i)_+^2 = M^1 \ddot{u}_-^1 + (\mathcal{E}_i)_-^2$$

$$M^2 \ddot{u}_+ - (\mathcal{E}_i)_-^2 = M^2 \ddot{u}_-^2 - (\mathcal{E}_i)_-^2$$

whose solution is

$$\ddot{u}_+ = \frac{M^1 \ddot{u}_-^1 + M^2 \ddot{u}_-^2}{M^1 + M^2},$$

and

$$2(\mathcal{E}_i)_+^2 = 2(\mathcal{E}_i)_-^2 + M^1 (\ddot{u}_-^1 - \ddot{u}_+) - M^2 (\ddot{u}_-^2 - \ddot{u}_+).$$

This completes the numerical specification of the solution at t_{n+1} ; this solution process is now repeated for each of the succeeding time steps.

At this point it is important to compare the solution procedure for impact of a continuum discretized by a finite element method with the solution procedure for a physically discrete body, i.e., a body composed of mass points joined by massless elastic springs. Both problems may be described by algebraic equations of the form of (57). The impenetrability condition is also identical. The impact conditions, however, are different. For the discretized continuum the procedure is described above. The study of the impact of mass points is considered in elementary mechanics books, e.g. [8]. The impact of two mass points is described by impulsive motion such that at t_- the velocities of the two mass points are V_-^1 and V_-^2 ; after impact at time t_+ , the two points have velocities V_+^1 and V_+^2 . The two points will not in general stay in contact (i.e., $V_+^1 \neq V_+^2$) but will rebound. The conditions used to compute the V_+^1 and V_+^2 are:

Balance of Momentum*

$$M^1 \{V^1\} + M^2 \{V^2\} = 0, \quad (71)$$

and use of an equation involving the "coefficient of restitution", e :

$$\frac{V_+^2 - V_+^1}{V_-^1 - V_-^2} = e. \quad (72)$$

For $e=1$ energy is conserved whereas for $e=0$ the points "stick" and energy is dissipated. We must comment in passing that (72) is the energy

* $\{f(\epsilon)\} = f(t_+) - f(t_-)$.

equation in disguise. To see this we can write the jump conditions for energy as

$$\frac{1}{2} M^1 \{(\dot{v}^1)^2\} + \frac{1}{2} M^2 \{(\dot{v}^2)^2\} = \{v\} \quad . \quad (73)$$

The term $\{v\}$ can exist only if other energies are dissipated during the jump. We rewrite (73) by using

$$\frac{1}{2} \{(\dot{v}^1)^2\} = [v^1] \langle \dot{v}^1 \rangle \quad ,$$

where

$$\langle \dot{v}^1 \rangle = \frac{1}{2} (\dot{v}_+^1 + \dot{v}_-^1) \quad . \quad (74)$$

Use of the momentum equation (71) then gives, after dividing by $M^1 [v^1]$

$$\langle \dot{v}^1 \rangle - \langle \dot{v}^2 \rangle = \frac{\{v\}}{M^1 [v^1]} \quad ,$$

or after recollecting terms and dividing by $(\dot{v}_-^1 - \dot{v}_-^2)$ we obtain:

$$\frac{\dot{v}_+^2 - \dot{v}_+^1}{\dot{v}_-^1 - \dot{v}_-^2} = 1 - \frac{\{v\}}{M^1 [v^1] (\dot{v}_-^1 - \dot{v}_-^2)} \quad . \quad (75)$$

The significance of the coefficient of restitution then is associated with the right hand side of (75).

It is clear from the above developments that the numerical simulation of the discretized continuum and the physically discrete system involve two distinct methods for treating the impact conditions. It is imperative then to associate the correct method for the problem at hand. In the

present study we are interested in the impact of continua, and in this case we shall employ the discrete shock condition to effect the solution. This a priori assumes that the response we are computing involves a time scale associated with wave propagation problems. Consequently, we cannot expect the computation procedure for advancing the solution in time to be accurate if we take time steps greatly in excess of transit times through each body. In this context it may be important to consider an "explicit" time integration procedure in future work. The stability restrictions may be too severe to make this feasible.

PART III

FEAP 74 - A COMPUTER PROGRAM FOR
SOLUTION OF CONTACT-IMPACT PROBLEMS13. Development of a Contact-Impact Model for FEAP

In order to incorporate an ability to compute solutions to contact-impact problems using a finite element method as described above it is necessary to have available a computer program which can solve the nonlinear equations of motion given by (61). The computer program FEAP is a general program to solve finite element problems. The program has a capability of solving both quasistatic and dynamic problems and can incorporate several types of elements simultaneously. The nonlinear capabilities required for the solution of contact-impact problems have been incorporated into FEAP and currently includes the user options (see Input Instructions):

- (1) Selection of quasistatic or dynamic option: The dynamic option will integrate the equations of motion using the one-step Newmark method to advance the solution in time. Quasistatic analysis is accommodated by any one-step algorithm. The algorithm employed is incorporated into each element routine and thus is defined by the developer of each element. Impact problems require description of the contact surface and wave speeds.
- (2) Selection of the nonlinear method to advance the solution: Options include:
 - (a) No iterations in each time step. Unbalanced forces at each time are added to the next time step.
 - (b) Iterations in each time step to achieve a balance of force within each time step. In this option the user can select to reform

the Jacobian matrix for each iteration or only at the first iteration in each time step.

In the impact problems solved to date it has been necessary to use the general form of the Newton-Raphson algorithm. This includes a complete forming and factoring of the Jacobian matrix for each iteration of each time step in the analysis. If the method described herein is to become computationally effective improvements in the computer program are paramount. Undoubtedly the most important aspect in reducing computer times is to introduce a substructuring system so that the highly nonlinear equations in the vicinity of the contact surface can be isolated from the remainder of the bodies. This will normally involve only a small number of equations in the total system of (62). The solution of a large finite element problem will generally concentrate the computer solution time in the forming and factoring of the tangent stiffness matrix. The fewer times that it is necessary to perform this costly step the more efficient the solution algorithm. Substructuring can be used then to restrict the part of the equations which must be formed and factored often, and thus greatly reduce the computer costs in analyzing impact problems.

The version of FEAP which can currently be used to analyze contact-impact problems includes, in addition to the nonlinear Newton-Raphson iterative algorithm, a new special contact-impact element and a new subroutine to describe impact surfaces and the discrete shock conditions described in Section 12. These are described in the following sections.

14. Contact Element for Hertzian Contact

The contact-impact element which has been developed is called ELMT05 and can be used along any coordinate direction. As developed it cannot be used along normals which are in non-coordinate directions. The development of the contact element assumes that within the framework of linear elasticity theory a node on \mathcal{B}^1 will impact on a node of \mathcal{B}^2 . In using this contact element we shall assume that the contact surface on \mathcal{B}^2 is located at larger coordinate values than the contact surface of \mathcal{B}^1 . The contact element is described by three nodes. Node 1 is associated with \mathcal{B}^1 , Node 3 is associated with \mathcal{B}^2 , and Node 2 is used as storage for the contact force $(\mathcal{E})^2$. The user can select the direction of contact motion by specifying the degree of freedom of the nodal unknowns to which the contact is to be measured; this must agree with the physical direction of the element (see Fig. 9). The degree of freedom for the contact element is specified during the MATERIAL data input and consists of a single card in I5 format. The element nodes are described along with all other elements according to Section 4 of the Input Instructions. The Node order as shown in Fig. 9 must be observed.

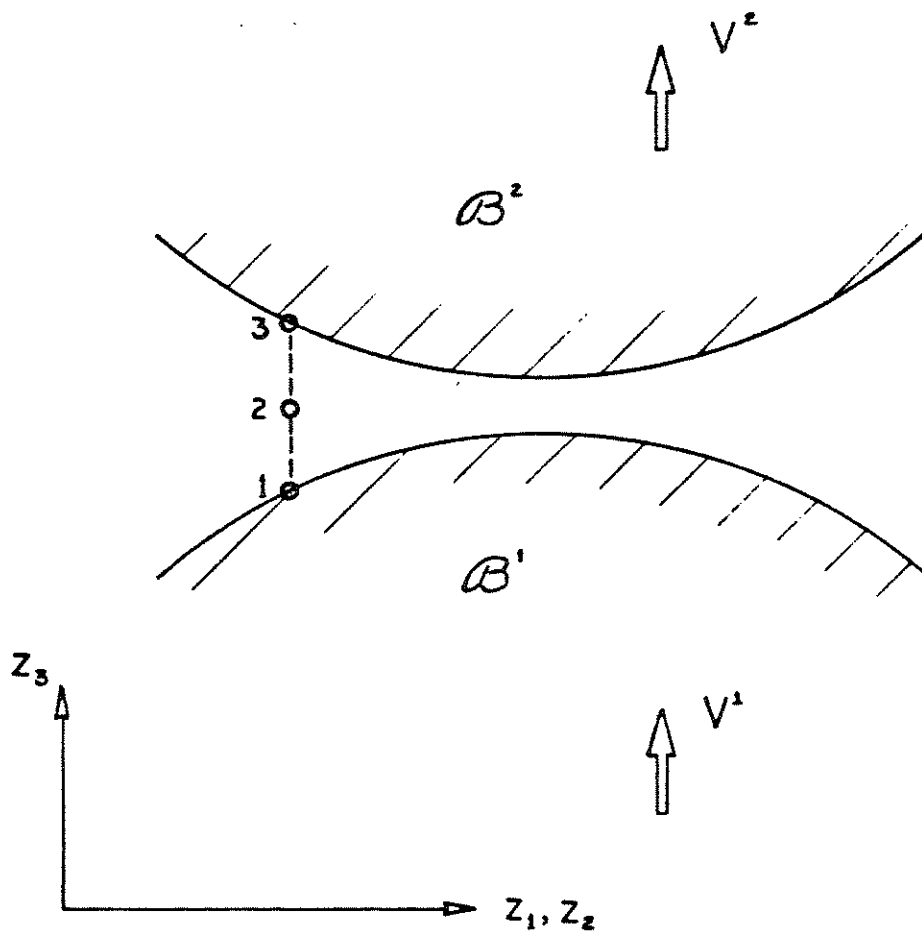


Figure 9

15. Impact Surface Description

The definition of the impact surface includes a list of all elements on the contact surface together with the degree of freedom describing the direction of contact motion (as described above). In addition, the product of mass density and wave velocity (always a positive number) for each body is input. This assumes, currently, that (1) each contact surface belongs to a linear material, and (2) the same material exists along all of the contact surface. This data need be prescribed only for impact problems, quasistatic contact problems do not require this data since no velocity or acceleration computations are performed in this class of problems. Data to be input for the impact surface is given in Table I.

Table I - Impact Surface Data

| | | |
|--------------------|----------|--|
| CARD 1) (6X,A6) | | |
| COL. 7 to 12 | | Must contain CONTAC |
| CARD 2) (2F10.0) | | |
| COL. 1 to 10 | ρU | of body 1 |
| COL. 11 to 20 | ρU | of body 2 |
| CARD 3) (I5) | | |
| COL. 1 to 5 | | NLIST, number of elements on contact surface |
| CARD 4) (2I5) | | |
| Repeat NLIST times | | |
| COL. 1 to 5 | | Contact element number |
| COL. 6 to 10 | | Degree of freedom of this contact element |

16. Example Problems

Two example problems are included to illustrate the characteristics of the methodology and the associated computer program described above for Hertzian contact problems. The first problem is a quasistatic contact problem which is used to demonstrate the ability of the computer program to compute an evolving contact surface. The second problem will demonstrate the ability of the program to properly model the temporal response of an impact problem.

To model a problem in which a contact surface will change under different load levels we consider two beams with an initial parabolic curvature. A symmetric configuration is analyzed and the resulting finite element model is shown in Fig. 10. Each element is nominally one unit by one unit. The gap at the load end is initially 0.5 units. The material properties used are $E = 500$ and $\nu = 0$. The load P is applied as shown and allowed to increase linearly in time. The problem then is to determine the contact surface at various load levels. In order to eliminate a singularity in the system of equations it was necessary to permanently attach the two nodes at the symmetry axis of the contact surface. All other nodes along the boundaries between the two bodies are assumed to be possible contact points and contact elements are assigned between vertical nodal pairs. The load was varied from 0.2 to 0.8 in increments of 0.1 and the computed contact surface and forces were computed. These are given in Table II. The deformed shape at a load of 0.4 is also shown in Fig. 10 as a dotted form. The attached node at the center has influenced the solution at loads above 0.3 since the contact pressure there is tensile (negative). The force is small and should not greatly affect the actual contact region computed. As the load increases the

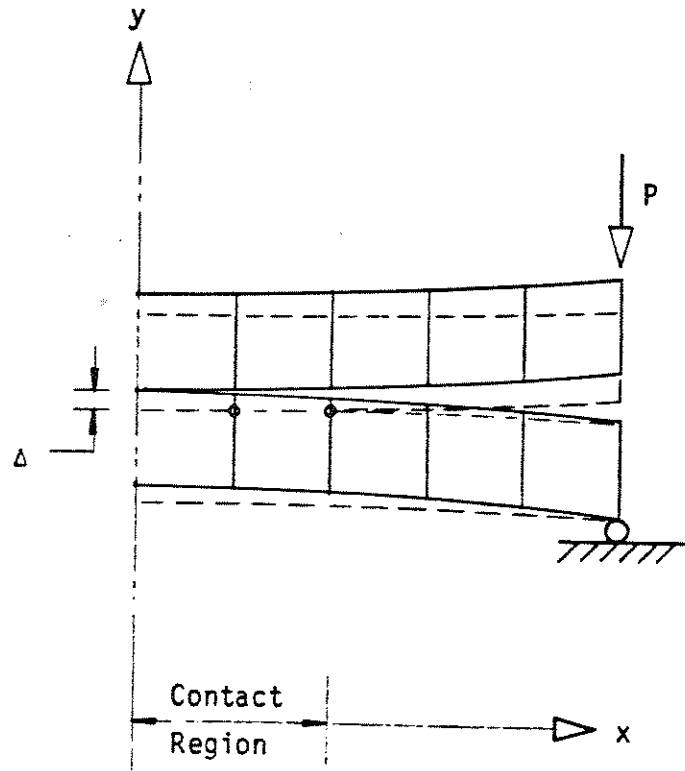


Figure 10

contact surface moves toward the load. This is conceptually correct since if the beams were modeled according to Euler-Bernoulli theory the contact force would be a point load which gradually moves from the center to the outer edge according to the relationship (using the above values for sizes and material properties)

$$X = 5 \left(1 - \frac{1}{6P} \right) .$$

This relation predicts that the contact point will be non-zero only after P exceeds $1/6$. The finite element model is in qualitative agreement with this beam theory, but since shear deformations are included the finite element solution gives a distributed load on the contact surface. It is interesting to also note that the contact force over the center of the beams is zero, just as in the beam theory.

Table II - Contact Forces

| LOAD | X-COORDINATE | | | | | | BEAM THEORY-X |
|------|--------------|-----|-----|-----|-----|---|------------------|
| | 0 | 1 | 2 | 3 | 4 | 5 | |
| 0.2 | 0.2 | - | - | - | - | - | 0.83 |
| 0.3 | .07 | .23 | - | - | - | - | 2.22 |
| 0.4 | -.01 | .07 | .34 | - | - | - | 2.92 |
| 0.5 | -0.00 | .02 | .23 | .25 | - | - | 3.33 |
| 0.6 | -.01 | - | .09 | .51 | .01 | - | 3.61 |
| 0.7 | -.01 | - | .07 | .45 | .19 | - | 3.81 |
| 0.8 | -.01 | - | .05 | .38 | .38 | - | 3.96 |

This problem demonstrates that the computer program can model the evolution of a contact surface. Of particular importance is to note that as the load increases the program can both attach and detach a contact point. This is an essential requirement for the analysis of

impact problems as is shown in the next problem.

As a simple example we consider the impact against a rigid wall of a finite, linear elastic rod traveling at constant velocity. The rod has a modulus of elasticity E of 100, and a mass density ρ of 0.1. The arrival velocity is taken to be 0.1 (units may be assigned in any convenient system). The rod is taken to be 10 units long and is divided into 10 elements plus one contact element as shown in Fig. 11. At time zero the rod is just arriving at the wall. The exact solution predicts a contact duration of 0.2 time units. This corresponds to the time required for a wave to travel from the contact point to the left end and back to the contact point at which time the rod will part from the wall. The problem was analyzed using FEAP with time steps of 0.01 unit (transit time across an element) and the rod remains in contact until time 0.20 units and has rebounded at time 0.21. Thus the program can predict accurately the contact duration of the rod. The finite element solution obtained is compared with the exact solution in Fig. 12. The agreement of stresses and contact force is good. The largest discrepancy exists in defining the shock front, which is "smeared" by the finite element method and ordinary differential equation solution method used here. This is the same type of solutions which are commonly obtained with numerical solutions of this type even without impact. Solutions such as the impact shocks generated are probably one of the most difficult responses to accurately calculate by a finite element method.

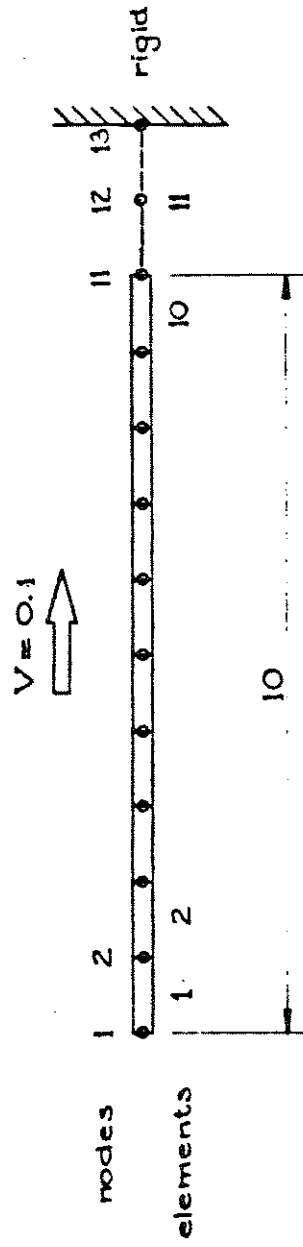
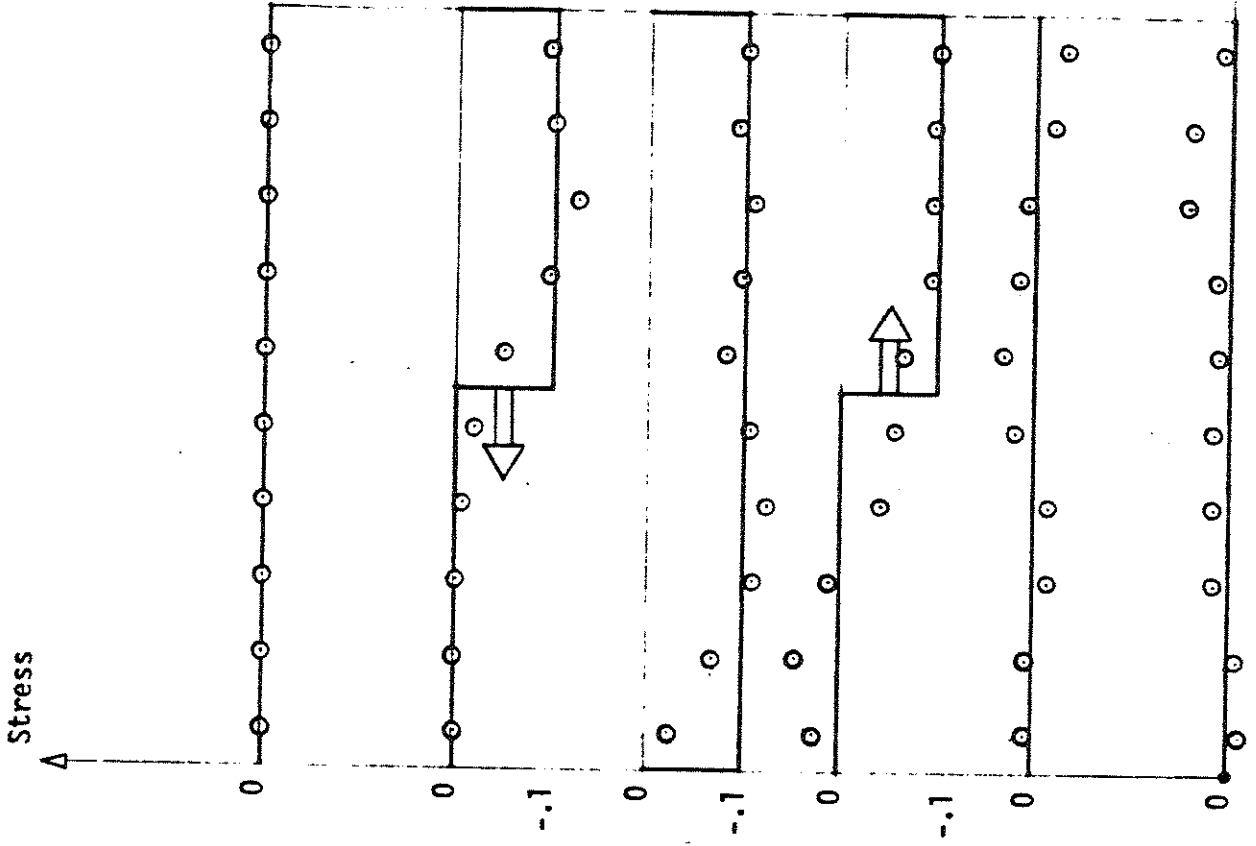
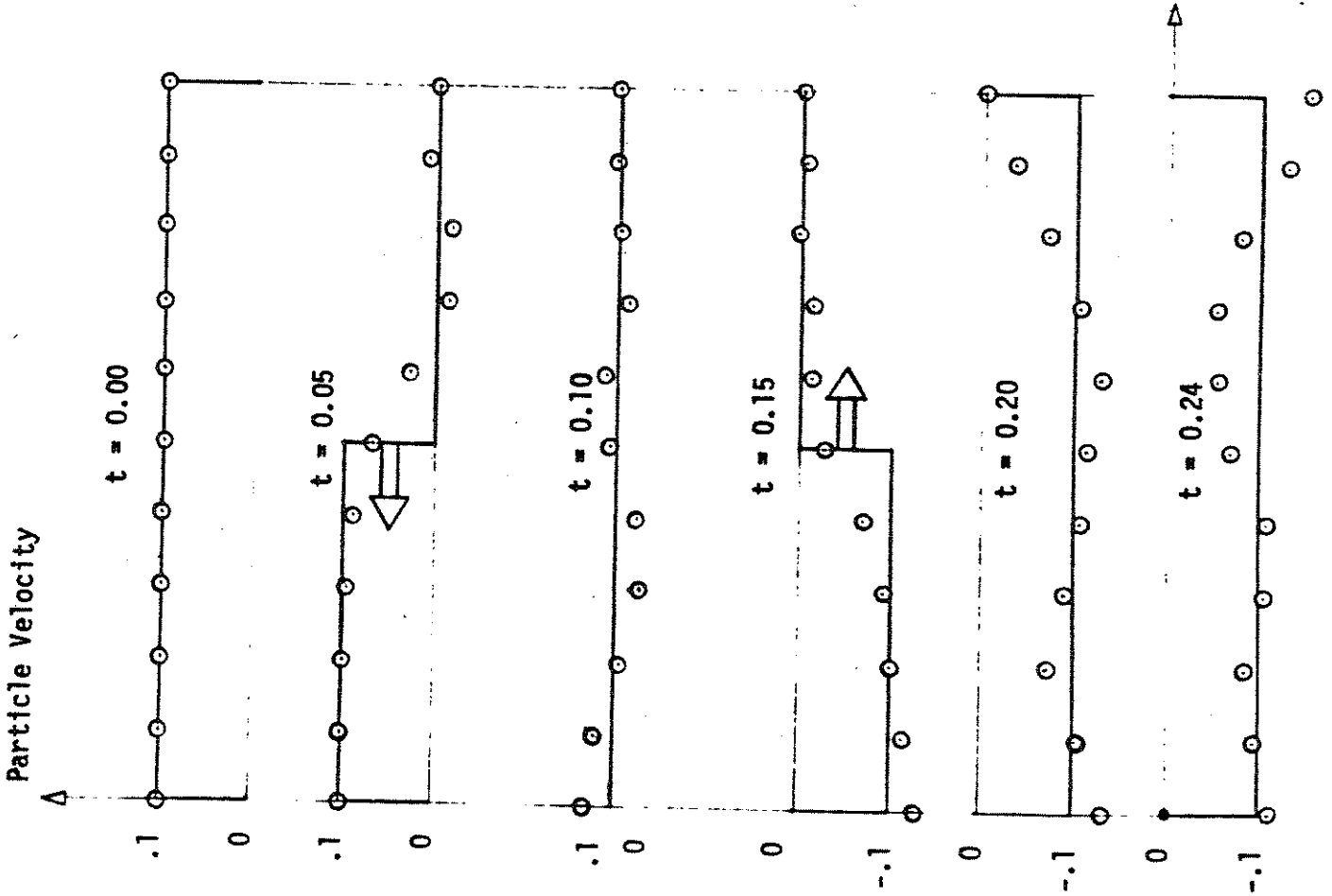


Figure 11

Comparison of Numerical Data with Exact Solution for Bar Impacting Rigid Wall



17. Closure and Recommendations for Future Work

In the preceding sections we have presented a theory for contact-impact problems together with the numerical development of a Hertzian contact-impact model. The computer program FEAP 74 has been modified to include the model and has successfully solved a contact problem and an impact problem. The work reported herein must be considered initiatory; the general theories and their numerical implementation have not been completed. The problem is of such a complicated nature and the literature existing prior to this study was so meager that we consider it fitting to document the work completed thus far.

We have attempted to qualify each stage of the development throughout the report, however, it may be fitting to reiterate future work which we consider to be essential for numerical models to be effective and efficient tools for predictive analyses.

- (1) The restriction of Hertzian type contact must be removed. This involves the non-trivial task of finding appropriate numerical methods to handle the χ^* maps.
- (2) Improved methods for solving the set of nonlinear algebraic equations must be found. We have suggested two methods which should be considered: (a) Substructure the problem about the contact regime so that a more efficient forming and factoring of the tangent stiffness can be performed; and (b) Since the impact problem is a wave propagation problem an explicit time integration of the equations of motion should be explored. In complex situations the explicit integration method may have severe stability limitations which could make it unacceptable.

- (3) Methods of utilizing the shock conditions need to be explored further. We have noted some peculiar anomalies when the bodies separate. These appear to be caused by a shock like separation phenomena.
- (4) When the wave propagation property of the impact problem is ignored by taking time steps greatly in excess of the transit times in a body the computed response is meaningless. Under such situations the bodies rebound within a single time step. Currently the rebound velocity is much too large. When the shock conditions are used for a class of problems where the response desired is in the target instead of in the impactor, it may be expedient to take a large time step. Methods should be explored to accomplish this capability.

The above recommendations for future work should in no way minimize what has been accomplished by the present study. For the first time a contact-impact theory in the form of a variational problem has been presented in a general form. This formulation was motivated by the fact that numerical solutions would be obtained by a finite element method. In addition the necessary foundation for the numerical solution has been thought out and within this context a computer program has been developed for Hertzian contact-impact problems.

The implementations considered here have produced results which are hopeful signs for the eventual success of the more general impact problems.

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APPENDICES

A. Input Instructions for Contact/Impact Problems

In order to analyze contact/impact problems in FEAP, users must prepare the data for a time dependent analysis. This will include the following Data Type Identification Cards (see Section 1, Appendix B):

FEAP 74

MATERIAL

NODAL

ELEMENT

CONTACT (for impact problems only)

loadings

INITIAL CONDITIONS (if non-zero)

and

VISCOE (for quasi-static contact problems)

or

IMPLICIT (for impact problems)

In performing the necessary solution to (62) the full Newton-Raphson method must be employed; this is controlled by the data in col. 76-80 of the first card following the VISCOE or IMPLICIT card, and consists of a negative number (negative uses full Newton-Raphson iteration with the absolute value of the number giving the number of iterations to be performed before going to the next time step). As an example of the required input data Table A shows the input data used for the impact problem reported in Section 16.

14.19.14 OUTPUT
FEAP74 * * ONE DIMENSIONAL BAR CONTACT PROBLEM

1 X
1 U
3
13 NODAL
1 0.
11 10.
13 10.
11 ELEMENTS 111111
1 1 1 1
1 2
11 12 13

3 INITIAL CONDITIONS
1 DISPLACEMENT
1 VELOCITY
1 ACCELERATIONS
11 0
11 .1
11 .1

2 MATERIALS
1 ELM03
100. 1. .01
2 ELM05
1 CONTACT
1 1.0 1.0E+30
11 1

1 IMPLICIT
1.0 5 1 1 13 1 11 0 0 0.25
.01 25 1 1 13 1 11 0 0 0.25
STOP
.01 10 1 1 13 9 11 0 0 0.25

-3
-3
-3

B. Input Instructions for FEAP 74

The input instructions for the description of a finite element mesh, together with the initial and boundary conditions, is described by subroutine MANUAL listed on the following pages. The input of the contact surface for impact problems is described in Table I of this report.

The description of material properties for the contact element is described in Section 14. For material properties for other elements in FEAP special input instructions must be supplied.

SUBROUTINE MANUAL

C ***** USER INSTRUCTIONS AND INPUT FORMATS FOR FEAP-74 *****

FINITE ELEMENT ANALYSIS PROGRAM

***** I N D E X *****

| SECTION | WORD | C O N T E N T | MAN | IC |
|---------|--------|--------------------------------------|---------|----|
| 1.0 | | DATA DIRECTORY WORDS | MAN 2C | |
| 2.0 | FEAP74 | PROBLEM INITIATION AND CONTROL CARDS | MAN 3C | |
| 2.1 | REMARK | REMARK CARD | MAN 4C | |
| 2.2 | TITLE | TITLE CHANGE | MAN 5C | |
| 2.3 | STOP | EXECUTION TERMINATION | MAN 6C | |
| 3.0 | MATERI | MATERIAL CHARACTERIZATION | MAN 7C | |
| 4.0 | NODAL | SEQUENTIAL NODAL GENERATION | MAN 8C | |
| 4.1 | GENERA | NON-SEQUENTIAL NODAL GENERATION | MAN 9C | |
| 4.2 | BOUNDA | BOUNDARY CODE INPUT | MAN 10C | |
| 4.3 | POLAR | POLAR TO CARTESIAN CONVERSION | MAN 11C | |
| 5.0 | ELEMEN | ELEMENT GENERATION | MAN 12C | |
| 5.1 | BLOCK | BLOCK MESH GENERATION | MAN 13C | |
| 6.0 | VECTOR | USER VECTOR INPUT | MAN 14C | |
| 6.1 | INITIA | INITIAL CONDITION | MAN 15C | |
| 7.0 | FORCE | GENERALIZED NODAL FORCE | MAN 16C | |
| 7.1 | BLOADS | SURFACE LOADS | MAN 17C | |
| 7.2 | ELOADS | ELEMENT LOADS | MAN 18C | |
| 7.3 | | PROPORTIONAL LOADS | MAN 19C | |
| 8.0 | SOLVE | INITIATION OF STATIC SOLUTION | MAN 20C | |
| | RESOLV | STATIC SOLUTION | MAN 21C | |
| 8.1 | EXPLIC | EXPLICIT DYNAMIC INTEGRATION | MAN 22C | |
| 8.2 | IMPLIC | IMPLICIT INTEGRATION | MAN 23C | |
| | VISCOE | IMPLICIT TIME INTEGRATION | MAN 24C | |
| 8.3 | MESH | MESH CHECK | MAN 25C | |
| | PLOT | PLOT MESH | MAN 26C | |
| 8.4 | FOURIE | FOURIER SERIES HARMONICS | MAN 27C | |
| 9.0 | | OUTPUT CONTROL | MAN 28C | |

FEAP74 IS A GENERAL (F)INITE (E)LEMENT (A)NALYSIS (P)ROGRAM WHICH FURNISHES TO THE USER MESH INPUT/OUTPUT, ELEMENT ASSEMBLY AND SOLUTION OF EQUATIONS (LINEAR, IMPLICIT AND EXPLICIT TIME DEPENDENT, NONLINEAR), PRESCRIBED GENERALIZED NODAL FORCES, PRESCRIBED NODAL AND ELEMENT DATA, AND OUTPUT OF THE GENERALIZED DISPLACEMENTS AND FORCES. ELEMENT MATRICES FOR TWO AND THREE DIMENSIONAL LINEAR ELASTICITY, SHELLS, PLATES, AND FIELD (LAPLACE EQUATION) PROBLEMS ARE AVAILABLE. ALTERNATIVELY USERS MAY SUPPLY THEIR OWN ELEMENT LIBRARY BY PROVIDING A SUBROUTINE CALLED ELMTNN, WHERE NN IS A TWO DIGIT NUMBER (01-10), IDENTIFYING THE ELEMENT SUBROUTINE. EACH ELEMENT SUBROUTINE HAS AT LEAST FOUR BASIC FUNCTIONS WHICH ARE DELINEATED BY A SWITCHING PARAMETER, ISW, IN THE SUBROUTINE.

MAN 29C
MAN 30C
MAN 31C
MAN 32C
MAN 33C
MAN 34C
MAN 34C
MAN 35C
MAN 35C
MAN 36C
MAN 36C
MAN 37C
MAN 38C
MAN 39C
MAN 40C
MAN 41C
MAN 42C
MAN 43C
MAN 44C
MAN 45C
MAN 46C
MAN 47C
MAN 48C
MAN 49C
MAN 50C

SEE SECTION 7.1 FOR DATA INPUT DETAILS.

INTEGRATION TABLE IS ACCESSED BY THE CALL

CALL INTEGL(LIM,NCI,NDIM,LINT,STUW)

STUW(4,M) INTEGRATION POINTS AND WEIGHTS.
 NOTE M MUST BE SET EXPLICITLY AND BE LARGER THAN OR EQUAL TO LINT.
 LINT - RETURNS WITH NUMBER INTEGRATION POINTS.
 NCI = 0 RETURNS GAUSS POINTS AND WEIGHTS IN STUW.
 LIM = 1 TO 5 IS NUMBER OF GAUSS POINTS/DIR-ECTION.
 NCI = 1 RETURNS A SPECIAL 3-D GAUSS FORMULA.
 SET LIM = 1 FOR 6 PT. CUBIC ACCURACY
 SET LIM = 2 FOR 14 PT. QUINTIC ACCURACY.
 NCI = 2 RETURNS TRIANGULAR INTEGRATION FORMULA
 SET LIM = 1 FOR 1 PT. LINEAR ACCURACY.
 SET LIM = 2 FOR 3 PT. QUADRATIC ACCURACY.
 SET LIM = 3 FOR 7 PT. QUARTIC ACCURACY.

1.) DATA TYPE IDENTIFICATION CARDS (15,IX,12A6).

EACH DATA SEGMENT IS PRECEDED BY A CARD WHICH IDENTIFIES THE TYPE OF DATA AND LIMITS ON THE AMOUNT OF DATA WHICH IMMEDIATELY FOLLOWS THE CARD. EXCEPT AS NOTED THE DATA SEGMENTS MAY APPEAR IN ANY ORDER. THE IDENTITY CARDS MAY ALSO AID THE USER IN INTERPRETTING THE INPUT DATA CARDS. AS SUPPLIED THERE ARE TWENTY-FIVE DIFFERENT DATA IDENTIFICATION CARDS. THESE ARE

COL 7 TO 12 IDENTITY(RESTRICTIONS)

FEAP74 START OF EACH PROBLEM (MUST PRECEDE ALL OTHER DATA).

TITLE CHANGE OUTPUT PAGE HEADINGS

REMARK COMMENTS ON OUTPUT

MATERI MATERIAL CHARACTERIZATION.

NODAL NODAL CARDS

POLAR POLAR CONVERSION. (PRECEDE BY NODAL, GENERA, OR BLOCK)

ELEMEN ELEMENT CONNECTION CARDS.

GENERA GENERATE NODES IN A LINEAR PATH BY ANY INCREMENT

BLOCK GENERATE ALL MESH DATA (BOTH NODAL AND ELEMENT)

BOUNDARY CODE PRESCRIPTION (PRECEDE BY NODAL OR GENERATE OR BLOCK)

BOUNDARY CODE PRESCRIPTION (PRECEDE BY NODAL OR GENERATE OR BLOCK)

VECTOR NODAL, OR GENERA OR ELEMENT DATA (PRECEDE BY NODAL, OR GENERA AND ELEMEN, OR BLOCK)

FORCE NODAL GENERALIZED FORCES (PRECEDE BY NODAL OR GENERA OR BLOCK).

- MAN104C
- MAN105C
- MAN106C
- MAN107C
- MAN108C
- MAN109C
- MAN110C
- MAN111C
- MAN112C
- MAN113C
- MAN114C
- MAN115C
- MAN116C
- MAN117C
- MAN118C
- MAN119C
- MAN120C
- MAN121C
- MAN122C
- MAN123C
- MAN124C
- MAN125C
- MAN126C
- MAN127C
- MAN128C
- MAN129C
- MAN130C
- MAN131C
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- MAN136C
- MAN137C
- MAN138C
- MAN139C
- MAN140C
- MAN141C
- MAN142C
- MAN143C
- MAN144C
- MAN145C
- MAN146C
- MAN147C
- MAN148C
- MAN149C
- MAN150C
- MAN151C
- MAN152C
- MAN153C
- MAN154C
- MAN155C
- MAN156C
- MAN157C

BLOADS SURFACE LOADINGS (SAME AS FORCE).
 ELOADS ELEMENT LOADINGS (SAME AS FORCE).
 MESH CHECK CONSISTENCY OF MESH ONLY (SAME AS SOLVE)
 PLOT PLOT MESH (SAME AS SOLVE)
 INITIA INITIAL CONDITION PRESCRIPTION FOR DYNAMIC
 SOLVE ANALYSIS (PRECEDE BY NODAL, GENERA OR BLOCK)
 COMPLETE FORMULATION AND SOLUTION FROM ELEMENTS
 (PRECEDE BY MATERI, NODAL OR GENERA, AND ELEMEN
 OR PRECEDE BY MATERI AND BLOCK)
 RESOLV USE PREVIOUS PROBLEM DESCRIPTION WITH NEW LOAD
 ONLY (PRECEDE BY SOLVE AND NEW LOADING CARDS).
 EXPLIC DYNAMIC SOLUTION BY EXPLICIT INTEGRATION, (SAME
 AS SOLVE)
 IMPLIC IMPLICIT INTEGRATION OF DYNAMIC PROBLEMS
 (PRECEDE BY SAME DATA AS FOR SOLVE)
 VISCOE QUASI-STATIC LINEAR VISCOELASTIC INTEGRATION
 (PRECEDE BY SAME DATA AS FOR SOLVE)
 FOURIE FOURIER COMPOSITION (SAME AS SOLVE)
 ADDUP ACCUMULATE FOURIER SOLUTION (AFTER FOURIE)
 STOP NORMAL EXIT (MUST FOLLOW ALL DATA)

NOTE EACH IDENTIFIER IS PUNCHED STARTING IN COL 7 (LEFT
 JUSTIFIED).

EXCESS CARDS MAY EXIST BETWEEN EACH SECTION OF DATA, HOWEVER,
 THE DATA TO BE USED MUST IMMEDIATELY FOLLOW THE TYPE CARD AND
 MUST BE IN PROPER ORDER. NO PARTICULAR ORDER OF THE TYPE
 CARDS IS NECESSARY EXCEPT THAT THE FEAP74 CARD MUST ALWAYS BE
 THE FIRST CARD IN EACH SET OF DATA, AND RESTRICTIONS MUST BE
 OBSERVED.

2.) PROBLEM INITIATION AND CONTROL CARDS

CARD 1. (6X,12A6)

COL 7 TO 12 MUST CONTAIN WORD FEAP74
 COL 13 TO 78 OUTPUT PAGE HEADER

CARD 2. (15,1X,3A6)

COL 1 TO 5 NDIM - SPATIAL DIMENSION OF PROBLEM (1 TO 3)
 COL 7 TO 12 NAMES TO BE PRINTED AS OUTPUT HEADERS TO
 COL 13 TO 18 COORDINATES - IF BLANK SET TO 1,2,3 AS NEEDED.
 COL 19 TO 24

CARD 3. (15,1X,6A6)

COL 1 TO 5 NDF - NUMBER OF UNKNOWN'S PER NODE (1 TO 6)
 COL 7 TO 12 NAMES TO BE PRINTED AS OUTPUT HEADERS OF THE
 COL 13 TO 18 GENERALIZED DISPLACEMENTS AND FORCES - IF
 BLANK SET TO 1,2,3,4,5,6 AS NECESSARY

COL 37 TO 42

MAN158C
 MAN159C
 MAN160C
 MAN161C
 MAN162C
 MAN163C
 MAN164C
 MAN165C
 MAN166C
 MAN167C
 MAN168C
 MAN169C
 MAN170C
 MAN171C
 MAN172C
 MAN173C
 MAN174C
 MAN175C
 MAN176C
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 MAN178C
 MAN179C
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 MAN258C
 MAN259C
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 MAN261C
 MAN262C
 MAN263C
 MAN264C
 MAN265C

CARD 4. (615,5F10.0)
 COL 1 TO 5 MEN - MAXIMUM NUMBER OF NODES CONNECTED TO ANY
 ELEMENT (1 TO 20).
 COL 6 TO 10 NEXTRA - INCREASES ELEMENT MATRIX SIZE FROM
 NDF*NEN TO NDF*NEN + NEXTRA
 COL 11 TO 15 IREC - COMPUTE GENERALIZED FORCE CHECK IF
 NONZERO (FOR TIME INVARIANT ANALYSIS ONLY)
 COL 16 TO 20 MBAN - MAXIMUM EXPECTED BANDWIDTH, DEFAULT IS
 SET TO 100. USED AS AN ERROR CHECK TO PREVENT
 RUNNING WITH AN OBVIOUS ERROR.
 COL 21 TO 25 IBUF - BUFFER SIZE FOR STORAGE OF HISTORY
 EFFECTS IN TIME DEPENDENT ANALYSIS. DEFAULT IS
 IBUF = ISZDT/20
 COL 26 TO 30 NC1 - USER INTEGER CONSTANT
 COL 31 TO 40 CON1 - USER DEFINED CONSTANT
 COL 41 TO 50 CON2 - USER DEFINED CONSTANT
 COL 51 TO 60 CON3(1) - USER DEFINED CONSTANT
 COL 61 TO 70 CON3(2) - USER DEFINED CONSTANT
 COL 71 TO 80 CON3(3) - USER DEFINED CONSTANT

2.1) *REMARK* USER COMMENTS ON OUTPUT. (6X,12A6)

SUBSEQUENT CARDS

COL 7 TO 12 MUST CONTAIN REMARK
 COL 13 TO 78 STATEMENTS TO BE OUTPUT, USE AS MANY REMARK
 CARDS AS DESIRED. INSERT BEFORE ANY TYPE CARD.

2.2) TITLE CHANGE ON OUTPUT (6X,12A6)

COL 7 TO 12 MUST CONTAIN TITLE
 COL 13 TO 78 NEW TITLE DESCRIPTOR

2.3) EXECUTION TERMINATION (6X,A4)

COL 7 TO 10 MUST CONTAIN STOP. INSERT AFTER LAST PROBLEM.

3.) MATERIAL CHARACTERIZATION (15,1X,12A6)

COL 1 TO 5 NUMMAT - NUMBER OF DIFFERENT MATERIAL CHARACT-
 ERIZATIONS TO FOLLOW.
 COL 7 TO 12 MUST CONTAIN WORD MATERI

THE FOLLOWING CARDS ARE SUPPLIED FOR EACH MATERIAL TO BE CHARAC
 TERIZED. (MUST BE EXACTLY NUMMAT SETS OF CARDS)

CARD 1.) ELEMENT SELECTOR CARD (15,1X,A5,11A6)

COL 1 TO 5 MATERIAL NUMBER (1 TO NUMMAT)
 COL 7 TO 11 ELMNN - WHERE NN IS NUMBER OF ELEMENT CLASS (01
 TO 10) TO WHICH THE CHARACTERIZATION BELONGS.
 COL 12 TO 77 ALPHANUMERIC INFORMATION TO BE OUTPUT.

CARD 2.), ETC. ** USER DEFINED FOR EACH ELEMENT TYPE PROVIDED.
 EXCESS BLANK CARDS ARE PERMISSIBLE BETWEEN EACH MATERIAL SET.
 4.) NODAL CARDS (15,1X,A6)

COL 1 TO 5 NUMNP - NUMBER OF NODAL POINTS
 COL 7 TO 12 MUST CONTAIN NODAL

SUBSEQUENT CARDS LAST NODAL CARD MUST NOT BE GENERATED.
 (15,115,3F10,0)

COL 1 TO 5 NODE NUMBER
 COL 15 1 IF 1 DISPLACEMENT IS SPECIFIED
 COL 16 1 IF 2 DISPLACEMENT IS SPECIFIED
 COL 17 1 IF 3 DISPLACEMENT IS SPECIFIED
 COL 18 1 IF 4 DISPLACEMENT IS SPECIFIED
 COL 19 1 IF 5 DISPLACEMENT IS SPECIFIED
 COL 20 1 IF 6 DISPLACEMENT IS SPECIFIED
 COL 21 TO 30 1 COORDINATE VALUE
 COL 31 TO 40 2 COORDINATE VALUE * AS REQUIRED
 COL 41 TO 50 3 COORDINATE VALUE

NODAL CARDS MUST BE IN ORDER. MISSING NODES ARE INTERPOLATED LINEARLY FROM INPUT NODES. IF SUCCEEDING CARDS HAVE IDENTICAL BOUNDARY CODES, THIS BOUNDARY CODE WILL BE ASSIGNED TO THE INTERVENING NODES. IN ALL OTHER CASES THE BOUNDARY CODE IS SET TO ZERO *TERMINATE ON NODE NUMNP OR A BLANK CARD*

4.1) NON SEQUENTIAL NODAL GENERATOR OPTION. (15,1X,A6)

COL 1 TO 5 NUMBER OF NODAL POINTS
 COL 7 TO 12 MUST CONTAIN GENERA

SUBSEQUENT CARDS (215,110,3F10,0)

COL 1 TO 5 NODE-NUMBER
 COL 6 TO 10 NODE-NUMBER-INCREMENT WHICH WILL BE SUCCESSIVELY ADDED TO NODE-NUMBER UNTIL SUM IS GREATER THAN NODE-NUMBER ON FOLLOWING CARD (ALGEBRAIC).
 BOUNDARY CODE. SAME AS INPUT FOR NODAL.
 IF SUCCEEDING CARDS HAVE IDENTICAL BOUNDARY CODES, THIS BOUNDARY CODE WILL BE ASSIGNED TO THE INTERVENING NODES. IN ALL OTHER CASES THE BOUNDARY CODE IS SET TO ZERO.

COL 21 TO 30 1 COORDINATE VALUE *
 COL 31 TO 40 2 COORDINATE VALUE * AS REQUIRED *
 COL 41 TO 50 3 COORDINATE VALUE *

*TERMINATE WITH BLANK CARD *

4.2) BOUNDARY CODE PATCH-UP OPTION. (6X,A6)

- MAN266C
- MAN267C
- MAN267C
- MAN267C
- MAN269C
- MAN269C
- MAN270C
- MAN271C
- MAN272C
- MAN273C
- MAN274C
- MAN275C
- MAN276C
- MAN277C
- MAN278C
- MAN279C
- MAN280C
- MAN281C
- MAN282C
- MAN283C
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- MAN301C
- MAN302C
- MAN303C
- MAN304C
- MAN305C
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 MAN358C
 MAN359C
 MAN360C
 MAN361C
 MAN362C
 MAN363C
 MAN364C
 MAN365C
 MAN366C
 MAN367C
 MAN368C
 MAN369C
 MAN370C
 MAN371C

COL 7 TO 12 MUST CONTAIN BOUNDARY
 SUBSEQUENT CARDS. (815)
 COL 1 TO 5 N, NODE NUMBER TO HAVE REDEFINED BOUNDARY CODE.
 COL 6 TO 10 NX, GENERATOR INCREMENT TO BE ADDED ALGEBRAICALLY TO N,
 UNTIL SUM EXCEEDS (MAX OR MIN) THE VALUE OF THE FOLLOWING CARD.
 COL 11 TO 15 IBC(1), (I=1,2,...,NDF) CODE FOR SPECIFYING FORCE
 COL 16 TO 20 OR DISPLACEMENT BOUNDARY CONDITIONS,
 COL ... IBC(1) .EQ. 0, FORCE SPECIFIED.
 IBC(1) .GT. 0, DISPLACEMENT SPECIFIED, NO INTERVENING GENERATION.
 IBC(1) .LT. 0, DISPLACEMENT SPECIFIED, GENERATE BETWEEN MISSING NODES IN ALGEBRAIC INCREMENTS OF NX.

* TERMINATE WITH A BLANK CARD. *

4.3) POLAR OR CYLINDRICAL COORDINATE CONVERSION TO CARTESIAN COORDINATES (6X,A6)

COL 7 TO 12 MUST CONTAIN POLAR (LEFT JUSTIFIED)
 CARD 1. (315,5X,2F10.0)

COL 1 TO 5 N1, FIRST NODE TO BE CONVERTED
 COL 6 TO 10 N2, LAST NODE TO BE CONVERTED
 COL 11 TO 15 N3, INCREMENT ADDED (ALGEBRAICALLY), N1 TO N2
 COL 21 TO 30 X0, ORIGIN OF POLAR X-COORDINATE
 COL 31 TO 40 Y0, ORIGIN OF POLAR Y-COORDINATE

5.) ELEMENT CARDS (15,1X,A6)

COL 1 TO 5 NUMEL - NUMBER OF ELEMENTS
 COL 7 TO 12 MUST CONTAIN ELEMENT

SUBSEQUENT CARDS (415,20I3/20I4)

CARD 1.

COL 1 TO 5 ELEMENT NUMBER
 COL 6 TO 10 MATERIAL NUMBER
 COL 11 TO 15 NUMBER OF SUBSEQUENT ELEMENTS USING SAME STIFFNESS MATRIX * SAVES RECOMPUTATION OF SIMILAR MATRICES. ELEMENT MUST ALSO HAVE SAME ELEMENT FORCE VECTOR * IF THESE ARE IN THE STIFFNESS SUBROUTINE *
 COL 16 TO 20 PRINT ELEMENT MATRIX IF NONZERO.
 COL 21 TO 23 IXD(1) ELEMENT INCREMENT ARRAY ON NODE 1.
 COL 24 TO 26 IXD(2) * IF NOT INPUT IS SET AUTOMATICALLY UP TO
 COL 79 TO 80 IXD(20) FOR SERENDIPITY ELEMENTS * SEE REPORT

MAN479C
 MAN479C
 MAN490C
 MAN491C
 MAN492C
 MAN493C
 MAN494C
 MAN495C
 MAN496C
 MAN497C
 MAN498C
 MAN499C
 MAN500C
 MAN501C
 MAN502C
 MAN503C
 MAN504C
 MAN505C
 MAN506C
 MAN507C
 MAN508C
 MAN509C
 MAN510C
 MAN511C
 MAN512C
 MAN513C
 MAN514C
 MAN515C
 MAN516C
 MAN517C
 MAN518C
 MAN519C
 MAN520C
 MAN521C
 MAN522C
 MAN523C
 MAN524C
 MAN525C
 MAN526C
 MAN527C
 MAN528C
 MAN529C
 MAN530C
 MAN531C

CARD 1. (6X.2A6) REPEAT NICD TIMES
 COL 7 TO 18 DESCRIPTIVE TITLE FOR INITIAL CONDITIONS
 CARD 2. (215.7F10.0)
 COL 1 TO 5 POSITION NUMBER, AS IN VECTOR CARDS FOR IPICK=1
 COL 6 TO 10 GENERATOR INCREMENT
 COL 11 TO 20 INITIAL CONDITION 1
 COL 21 TO 30 INITIAL CONDITION 2
 COL AS REQUIRED FOR NICD INITIAL CONDITIONS

INTERPOLATION BETWEEN INPUT VALUES AS DESCRIBED IN VECTOR INPUT
 ***** NOTE ***** IF MISSING THE INITIAL CONDITIONS ARE SET ZERO

7.) FORCE CARDS (15.1X.A6)
 COL 1 TO 5 LAST NODE TO WHICH A FORCE IS TO BE SPECIFIED
 COL 7 TO 12 MUST CONTAIN FORCE
 SUBSEQUENT CARDS (15.5X.6F10.0)

THE FOLLOWING VALUES ARE EACH INTERPRETTED AS FORCES IF THE
 CORRESPONDING BOUNDARY CODE IS A 0 *ZERO* AND AS A DISPLACEMENT
 IF THE CORRESPONDING BOUNDARY CODE IS 1 *ONE*.

COL 1 TO 5 NODE TO WHICH FORCE OR DISPLACEMENT IS APPLIED
 COL 11 TO 20 VALUE OF 1 FORCE/DISPLACEMENT
 COL 21 TO 30 VALUE OF 2 FORCE/DISPLACEMENT * AS *
 COL 31 TO 40 VALUE OF 3 FORCE/DISPLACEMENT * REQUIRED *
 COL 41 TO 50 VALUE OF 4 FORCE/DISPLACEMENT
 COL 51 TO 60 VALUE OF 5 FORCE/DISPLACEMENT
 COL 61 TO 70 VALUE OF 6 FORCE/DISPLACEMENT

7.1) SURFACE LOAD CARDS (15.1X.A6)
 COL 1 TO 5 NUMBER OF LOADED FACE CARDS
 COL 7 TO 12 MUST CONTAIN BLOADS
 CARD 1. (15.1X.A5.14.815.813)

COL 1 TO 5 DIMENSION OF LOADING SURFACE, (1 OR 2).
 COL 7 TO 11 SLD(NN), ALPHA-NUMERIC NAME OF SURFACE LOADING
 SUBROUTINE (NN IS BETWEEN 1 AND 5)
 COL 12 TO 15 NRT, NUMBER OF ADDITIONAL ELEMENT LOAD
 SURFACES TO BE GENERATED FROM CURRENT MODEL.
 COL 16 TO 20 IPRES(N), NODE NUMBERS DEFINING LOADING SURFACE
 COL 21 TO 25 OF CURRENT ELEMENT.
 COL ... (IDENTIFY FROM 2 TO 8 AS REQUIRED)
 COL 51 TO 55 INC(N), INCREMENT VALUE ADDED TO IPRES(N) TO
 COL 56 TO 58 IDENTIFY NODE NUMBERS OF A GENERATED SEQUENCE.
 COL 59 TO 61

MAN532C
 MAN533C
 MAN534C
 MAN535C
 MAN536C
 MAN537C
 MAN538C
 MAN539C
 MAN540C
 MAN541C
 MAN542C
 MAN543C
 MAN544C
 MAN545C
 MAN546C
 MAN547C
 MAN548C
 MAN549C
 MAN550C
 MAN551C
 MAN552C
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 MAN576C
 MAN577C
 MAN578C
 MAN579C
 MAN580C
 MAN581C
 MAN582C
 MAN583C
 MAN584C
 MAN585C

COL ... (IDENTIFY FROM 2 TO 8 AS REQUIRED)
 CARD 2. (8F10.0)
 COL 1 TO 80 LOAD AT NODES GIVEN ON PREVIOUS CARD *
 MUST CORRESPOND IN SEQUENCE TO THE NODE NUMBERS

7.2) ELEMENT LOAD CARDS (15,1X,A6)

COL 1 TO 5 NLD, NUMBER OF ELEMNT LOAD CARDS.
 COL 7 TO 11 MUST CONTAIN ELOADS

SUBSEQUENT CARDS (15,1X,A5,14,15,6F10.0)

COL 1 TO 5 IEL, INITIAL ELEMENT OF A GENERATED SEQUENCE.
 COL 7 TO 11 ELM(NN), ALPHA-NUMERIC NAME OF ELEMENT
 SUBROUTINE WHERE ELEMENT LOADS ARE COMPUTED.
 USED AS CHECK TO INSURE IEL, ETC. ARE PROPER
 ELEMENTS.

COL 12 TO 15 INC, INCREMENT NUMBER IN A GENERATED SEQUENCE,
 (DEFAULT = 1).

COL 16 TO 20 JEL, TERMINAL ELEMENT NUMBER IN A GENERATED
 SEQUENCE, IF JEL = 0, ONLY IEL IS COUNTED.

COL 21 TO 80 PR-USER DEFINED VALUES FOR DETERMINING BODY
 LOADS IN THE ISU=5 PORTION OF ELM(NN).

NOTE, USER MUST PROVIDE COMPUTATION OF LOADS IN ELMTNN.
 PR IS TRANSFERRED TO SUBROUTINE ELMTNN IN THE U VECTOR,
 WHEN ISU =5, ONLY.

7.3) PROPORTIONAL LOADS FOR TIME DEPENDENT ANALYSIS

TRANSFER TO THIS OPTION OCCURS ONLY FOR TIME ANALYSES.

ONE CARD FOR EACH PROPORTIONAL LOAD REQUIRED

COL 1 TO 5 PROPORTIONAL LOAD TYPE, 1,2 OR 3
 COL 6 TO 10 K, TABLE CONSTANT
 COL 11 TO 20 TMIN, SMALLEST TIME LOADING IS VALID
 COL 21 TO 30 TMAX, LARGEST TIME LOADING IS VALID
 COL 31 TO 40 A0
 COL 41 TO 50 A1
 COL 51 TO 60 A2
 COL 61 TO 70 A3
 COL 71 TO 80 A4

LOAD TYPE 1. T = TIME

PROP = A0 + A1*T + A2*T*T + A3*T*T*T + A4*T*T*T*T

LOAD TYPE 2.

C PROP = A0*(SIN(A1*T))*K + A2*(COS(A3*T))*K + A5
 C LOAD TYPE 3.
 C

C PROP = USER DEFINED FUNCTION FROM SUBROUTINE
 C EXPRD(PROP,T,A) WHERE A IS AN ARRAY WITH
 C INFORMATION FOR COLUMNS 6-80 OF DATA CARDS.
 C

C **NOTE** PROPORTIONAL LOADS CAN BE ACCUMULATED FROM DIFFERENT
 C TYPES AT THE SAME TIME.
 C

C 8.) INITIATION OF TIME INDEPENDENT SOLUTION (15,IX,A6)

C COL 1 TO 5 IOUT, OUTPUT CONTROL CODE,
 C

C IOUT .EQ. 0, ALL STRESSES AND DISP. PRINTED
 C IOUT .NE. 0, SELECTED PRINTOUT, MORE DATA INPUT
 C SEE SECTION 9 FOR DATA PREPARATION.
 C

C COL 7 TO 12 MUST CONTAIN SOLVE *INDICATES ALL DATA INPUT*
 C COMPLETE FORMULATION AND SOLUTION OF EQUATIONS.
 C COL 7 TO 12 MUST CONTAIN RESOLV TO OBTAIN SUBSEQUENT
 C SOLUTIONS WHERE BOUNDARY CODES DO NOT CHANGE
 C AND ALL PRESCRIBED DISPLACEMENTS ARE ZERO.
 C

C 8.1) INITIATION OF DYNAMIC SOLUTION BY EXPLICIT INTEGRATION.

C COL 1 TO 5 IPRT, OUTPUT CONTROL FOR DISPLACEMENT AND
 C STRESS PRINTOUT. SEE SECT. 9 FOR DATA INPUT.
 C IOUT = 1 - MIN(1,IPRT)
 C IF IOUT .NE. 0, THE SPATIAL CONTROL DATA
 C COMES AT THE END OF THE DYNAMIC SEGMENT.
 C MUST CONTAIN EXPLIC

C COL 7 TO 12

C SUBSEQUENT CARDS (215,2F10.0,215)

C COL 1 TO 5 NUMBER OF TIME STEPS
 C COL 6 TO 10 PRINT INTERVAL
 C COL 11 TO 20 TIME INCREMENT
 C COL 21 TO 30 NEWMARK DELTA-DAMPING TERM (GAMMA - .5)
 C COL 31 TO 35 NUMBER OF TIME EVOLUTION ELEMENT VARIABLE PLOTS
 C COL 36 TO 40 NPROP, NUMBER OF PROPORTIONAL LOADS TO BE INPUT
 C COL 41 TO 45 NFORC, LAST NODE ON WHICH A FORCE IS CHANGED
 C DURING EACH TIME STEP.
 C COL 46 TO 50 KKK, STABILITY CHECK OVERRIDE ** CAUTION USE
 C ONLY WHEN A BETTER ESTIMATE OF THE STABLE TIME
 C STEP IS AVAILABLE THAN CAN BE PERFORMED BY CODE
 C KKK ZERO, USES INTERNAL STABILITY CHECK.
 C KKK NONZERO, DISREGARDS STABILITY CHECK.
 C NONL, 0 FOR LINEAR
 C COL 51 TO 55 1 FOR NON LINEAR
 C COL 56 TO 60 LPR, 0 PRINTS LOADS
 C 1 SUPPRESS PRINT
 C

- MAN586C
- MAN587C
- MAN588C
- MAN589C
- MAN590C
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MAN676C
MAN677C
MAN688C
MAN689C
MAN690C
MAN691C
MAN692C
MAN693C
MAN694C
MAN695C
MAN696C
MAN697C
MAN698C

SUBSEQUENT CARDS (315) ONE FOR EACH STRESS PLOT.

COL 1 TO 5 ELEMENT NUMBER CONTAINING STRESS TO BE PLOTTED.
COL 6 TO 10 LOCAL COORDINATE POINT CODE, 1 TO 9, AS PATTERNED AFTER, COL 11 TO COL 19, IN SECT. 9.
COL 11 TO 15 PLOT COMPONENT CODE, 1 TO 6 FOR SIGMA(I,J), I.E., SIGMA(1,1)=1, SIGMA(1,2)=2, SIGMA(1,3)=3, SIGMA(2,2) = 4, SIGMA(2,3) = 5, SIGMA(3,3) = 6.

IF(NPROP,NE.0) READ PROPORTIONAL LOAD CARDS, SEE SECT. 7.3

IF(NFORC,NE.0) READ FORCE CARDS AT EACH TIME STEP. IF OUTPUT IS LIMITED BY IOUT NONZERO, THE FIRST FORCE CARD SET PRECEDES OUTPUT CARDS AND THE REMAINDER FOLLOW THE OUTPUT CARDS NO BLANK CARDS MAY BE USED BETWEEN SETS OF CARDS OTHER THAN THE USUAL BLANK TERMINATOR CARD FOR FORCE INPUT CARDS.

IF(IOUT,NE.0) DATA FOR SPATIAL PRINTOUT CONTROL, SEE SECT.9.

SPECIAL COMMENTS FOR DYNAMIC OPTION

- (1) ONLY COLUMNS 1 TO 66 ARE AVAILABLE FOR PAGE HEADING.
- (2) MAXIMUM ADVANTAGE OF ELEMENT REUSE OPTION SHOULD BE TAKEN.
- (3) INITIAL CONDITIONS FOR DISPLACEMENT AND VELOCITY VECTORS, AS WELL AS STORAGE FOR ACCELERATION VECTOR, MAY BE MADE THROUGH INPUT OF AN INITIAL CONDITION CARD SET, WITHOUT SPECIFIED INITIAL CONDITIONS THEY ARE AUTOMATICALLY SET ZERO.
- (4) SPATIAL LOADING IS INPUT THROUGH FORCE OR BOUNDARY PRESSURE CARDS.
- (5) EXTREME CAUTION ON ORDER OF DATA CARDS MUST BE OBSERVED. NO EXTRA CARDS ARE PERMITTED AND STRICT COUNTS ARE OBSERVED EXCEPT FOR THE NUMBER OF FORCE CARDS USED IN EACH TIME STEP.

8.2) INITIATION OF IMPLICIT TIME INTEGRATIONS (15,1X,A6)

COL 1 TO 5 NSEQ, NUMBER OF TIME SEQUENCES
 COL 7 TO 12 MUST CONTAIN VISCODE FOR LINEAR VISCOELASTIC QUASI-STATIC PROBLEMS (ONE INITIAL CONDITION ONLY MUST BE USED)
 COL 7 TO 12 MUST CONTAIN IMPLIC FOR DYNAMIC IMPLICIT INTEGRATIONS (THREE INITIAL CONDITIONS ARE REQUIRED, MORE CAN BE SPECIFIED WITHOUT ERROR)

SUBSEQUENT CARDS, ONE SET FOR EACH TIME SEQUENCE

CARD 1. (F10.0,0.815,2F10.0,2I5)

COL 1 TO 10 DT, TIME INCREMENT (NONZERO FOR IMPLIC)
 COL 11 TO 15 NTS, NUMBER OF TIME STEPS IN SEQUENCE
 COL 16 TO 20 INT, PRINT INTERVAL (DEFAULT 1)
 COL 21 TO 25 NNI, FIRST NODE PRINTED
 COL 26 TO 30 NNE, LAST NODE PRINTED
 COL 31 TO 35 NEI, FIRST ELEMENT STRESS TO BE PRINTED

C COL 11 TO 20 1 - FORCE/DISPL. MULTIPLIER
 C COL 21 TO 30 2 - FORCE/DISPL. MULTIPLIER
 C COL 31 TO 40 3 - FORCE/DISPL. MULTIPLIER
 C COL 41 TO 50 4 - FORCE USER MULTIPLIER
 C COL 51 TO 60 5 - FORCE USER MULTIPLIER
 C COL 61 TO 70 6 - FORCE USER MULTIPLIER
 C COL 71 TO 80 USER CONSTANT.

9.) OUTPUT CONTROL FOR LIMITED PRINTS

DISPLACEMENT OUTPUT CONTROL, IF IOUT .NE. 0.

CARD 1. (IS)

COL 1 TO 5 NUMDIS - NUMBER OF DISPLACEMENT PRINT CARDS

SUBSEQUENT CARDS (215) SKIP IF NUMDIS = 0

COL 1 TO 5 NODAL NUMBER TO BE OUTPUT.
 COL 6 TO 10 HIGHER NODE NUMBER OF A GENERATED SEQUENCE.
 IF ZERO JUST FIRST NODE IS COUNTED.
 COL 11 TO 15 INCREMENT TO GENERATOR, DEFAULT = 1
 *** REPEAT UNTIL NUMDIS CARDS HAVE BEEN READ

STRESS OUTPUT CONTROL, IF IOUT .NE. 0.

CARD 1. (15,5X,911)

COL 1 TO 5 NUMSTR - NUMBER OF STRESS OUTPUT CARDS
 COL 11 TO 19 NSIG(9) - PRINT PATTERN WITHIN AN ELEMENT.
 LOCAL POINTS OF EACH ELEMENT CAN BE
 SUPPRESSED BY NON-ZERO ENTRIES AS FOLLOWS,
 E.G.

COL 11 SUPPRESS PRINT AT LOCAL POINT 1, (0, 0, 0)
 COL 12 SUPPRESS PRINT AT LOCAL POINT 2, (-1, 0, 0)
 COL 13 SUPPRESS PRINT AT LOCAL POINT 3, (1, 0, 0)
 COL 14 SUPPRESS PRINT AT LOCAL POINT 4, (0,-1, 0)
 COL 15 SUPPRESS PRINT AT LOCAL POINT 5, (0, 1, 0)
 COL 16 SUPPRESS PRINT AT LOCAL POINT 6, (0, 0,-1)
 COL 17 SUPPRESS PRINT AT LOCAL POINT 7, (0, 0, 1)
 COL 18 SUPPRESS PRINT AT LOCAL POINT 8
 COL 19 SUPPRESS PRINT AT LOCAL POINT 9

SUBSEQUENT CARDS (215) SKIP IF NUMSTR = 0

COL 1 TO 5 ELEMENT NUMBER TO BE PRINTED.
 COL 6 TO 10 HIGHER ELEMENT NUMBER OF A GENERATED SEQUENCE.
 IF ZERO ONLY FIRST ELEMENT IS COUNTED.
 COL 11 TO 15 INCREMENT TO GENERATOR, DEFAULT = 1
 *** REPEAT UNTIL NUMSTR CARDS HAVE BEEN READ

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C*****

C. Listing of the Contact/Impact Subroutines Added to FEAP 74

The listings for the subroutines which are added to FEAP 74 for the contact/impact theory described herein are given below.

```

000014 SUBROUTINE CONTAC(ISW,IX,NEL1,NDF,NUMP,DU,U,UD,UDD)
000014 C*****CONTAC **** 04/29/74 *****
000014 LOGICAL CFLAG,FLAG
000014 DIMENSION IX(NEL1,NUMP),DU(NDF,NUMP),U(NDF,NUMP),UD(NDF,NUMP),
000014 UDD(NDF,NUMP)
000014 DIMENSION RMP(10)
000014 COMMON/CONTACT/FLAG,CFLAG,LIST,ICLIST(10),ICDEG(10),RU1,RU2,
000014 RM(2,10)
000014 COMMON/LABELS/HEAD(12),O,IPG,XHED(3),UHED(6),XH,FH,UH,NSTR,FLAG(7)
000014 COMMON/TAPES/ ITP5,ITP6
000014 DATA CFLAG,FLAG/,FALSE,...FALSE./
000022 GO TO (1,2,3),ISW
000034 READ(ITP5,1000) LIST,RU1,RU2
000036 RP = RU1+RU2
000036 IF(RP.EQ.0.0) RP = 1.
000044 READ(ITP5,1001) (ICLIST(L),ICDEG(L),L=1,LIST)
000070 WRITE(ITP6,2000) O,HEAD,IPG,RU1,RU2,(ICLIST(L),L=1,LIST)
000131 IPG = IPG + 1
000133 FLAGC = .TRUE.
000134 RETURN
000134 DO 200 L = 1,LIST
000141 N = ICLIST(L)
000143 IDEG = ICDEG(L)
000144 I = IX(L,N)
000150 RM(1,L) = DU(IDEG,I)
000155 I = IX(3,N)
000161 RM(2,L) = DU(IDEG,I)
000171 WRITE(ITP6,2002)
000175 DO 220 L = 1,LIST
000202 N = ICLIST(L)
000204 RMP(L) = RM(1,L) + RM(2,L)
000210 IF(RM(1,L).EQ.0.0.OR.RM(2,L).EQ.0.0) RMP(L) = RMP(L)*1.E+30
000222 WRITE(ITP6,2001) N,(RM(1,L),I=1,2),RMP(L)
000256 DO 210 N = 1,NUMP
000260 DO 210 I = 1,NDF
000261 DU(I,N) = 0.0
000272 RETURN
000274 DO 300 L = 1,LIST
000274 N = ICLIST(L)
000276 IDEG = ICDEG(L)
000277 K = IX(2,N)
000303 UD(IDEG,K) = 0.0
000310 UDD(IDEG,K) = 0.0
000313 IF(U(IDEG,K)) 300,300,310
000317 I = IX(1,N)
000323 J = IX(3,N)
000326 CU = (RU1*UD(IDEG,I)+RU2*UD(IDEG,J))/RP
000342 UD(IDEG,I) = CU
000347 UD(IDEG,J) = CU
000352 CU = (RM(1,L)*UD(IDEG,I)+RM(2,L)*UD(IDEG,J))/RMP(L)
000375 IF(RM(1,L).GT.RM(2,L))U(IDEG,K)=U(IDEG,K)-RM(2,L)*(UDD(IDEG,J)-CU)
000420 IF(RM(1,L).LE.RM(2,L))U(IDEG,K)=U(IDEG,K)+RM(1,L)*(UDD(IDEG,J)-CU)
000443 UDD(IDEG,I) = CU

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000451 UDD(1DEG.J) = CU
000454 300 CONTINUE
000457 RETURN
000457 1000 FORMAT(15,5X,2F10.0)
000457 1001 FORMAT(2I5)
000457 2000 FORMAT(A1,12A6,30X,4HPAGE,14//5X,25HCONTACT SHEET DESCRIPTION//
C 10X,13HBODY 1 RO*U =,E13.5//
C 10X,13HBODY 2 RO*U =,E13.5//
X 11X,7HELEMENT,1X,9HDIRECTION/(115,110))
000457 2001 FORMAT(110,6E13.4)
000457 2002 FORMAT(/1AH ELEMENT,13H BODY 1 MASS,13H BODY 2 MASS)
000457 END

```

ELM 1C

SUBROUTINE ELMT05(N,MA,ND,IM,NDF,NEL,NELI,NSTF,NSIZV,NVEC,MCT,DM,D,
X,XYZ,IX,F,FORCE,ESTIF,U,VECT,ISW)

C***** ELMT05 ***** 04/29/74 *****

```

000027 LOGICAL NPR
000027 DIMENSION ESTIF(NSTF,NSTF),FORCE(NSTF,2)
000027 DIMENSION U(NDF,1),IX(NELI,1),PLOT(8),THED(8)
000027 COMMON/LOCALS/ DUL(6,20),UL(6,20),UDL(6,20),UDDL(6,20)
000027 COMMON/PRTPLT/ NSIG(9),NPLT(9,2),NT,NSTEP,NUNPLT,NEDATA(20,3),NPR
000027 COMMON/SHAP/ XJAC,SHAPE(4,20),SG(3,3),SK(3,3),X(3,20)
000027 COMMON/TAPES/ ITP5,ITP6
000027 COMMON/TIMHIS/ TIME,DT,DTP,NH,ISZH,C0,C1,C2,C3,C4,C5,C6,NCT
000027 EQUIVALENCE (PLOT,TAU),(PLOT(2),ETA)
000027 DATA THED/3HTAU,3HETA,2HU1,3HJD1,4HDD1,2HUJ3,3HUD3,4HUDD3/
000027 GO TO (1,2,3,5,3), ISW
000041 READ(ITP5,1000) IDEG,NPS
000051 IDEG = MAX0(IDEG,1)
000055 WRITE(ITP6,2000) IDEG,NPS
000064 TOL = 1.E-06
000066 I = NDF + IDEG
000073 J = I + NDF
000074 RJ = 0.5
000075 RJ = 0.5
000076 C6 = 0.
000077 NV = NPS
000100 NA = NPS + NPS
000101 RETURN
000102 DM = X(IDEG,3) - X(IDEG,1)
000106 DM = DM + 1.E-25
000110 RETURN
000110 TAU = UL(IDEG,2)
000112 Dn = (X(IDEG,3)-X(IDEG,1)) + (UL(IDEG,3)-UL(IDEG,1))
000120 ETA = 0.
000121 IF(DD.LT.TOL) ETA = 1.
000125 IF(TAU.LT.0.0) ETA = 0.0
000130 IF(ISW.GT.3) GO TO 4
000134 ESTIF(IDEG,1) = ETA
000142 ESTIF(I,IDEG) = ETA
000145 ESTIF(I,J) = -ETA
000151 ESTIF(J,1) = -ETA
000155 ESTIF(I,1) = 1.0 - ETA
000163 RETURN
000163 IF(ISW.EQ.4) GO TO 6
000166 FORCE(I,1) = ETA*DD
000174 IF(ETA.EQ.0.0) FORCE(I,1) = -TAU
000203 TAU = AMAX(0.0,TAU)
000207 FORCE(IDEG,1) = -TAU
000215 FORCE(J,1) = TAU
000220 IF(.NOT.NPP) WRITE(ITP6,2001) N,DM,MA,TAU,ETA
000243 IF(.NOT.NPLT,LE.0) RETURN
000246 PLOT(3) = UL(IDEG,1)
000250 PLOT(4) = UDL(IDEG,1)
000252 PLOT(5) = UDDL(IDEG,1)
000254 PLOT(6) = UL(IDEG,3)

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000256 PLOT(7) = UDL(IDEQ,3)
000260 PLOT(8) = UDDL(IDEQ,3)
000262 DO 60 K = 1,NUMPLT
000263 KK = NPLT(K,2)
000265 IF (NPLT(K,1).GT.0) CALL PLDATA(NDIM,NPLT(K,1),THED(KK),X(1,2),
      C PLOT(KK))
000306 60 CONTINUE
000311 5 RETURN
000312 1000 FORMAT(2I5)
000312 2000 FORMAT(5X,2IHPPOINT CONTACT ELEMENT/
000312 C 5X,27HCCONTACT DEGREE OF FREEDOM = 13,5X,5HNFS = 15)
000312 X 5HETA =,F3.1)
000312 2001 FORMAT(5X,7HELEMENT,15,5X,A5,5X,8HMATERIAL,13,5X,5HTAU =,E15.5,5X,
000312 END

```

```

SUBROUTINE ELMT09(N,MA,NDIM,NDF,NELM,NEL1,NSTF,NSIZV,NVEC,NCT,DM,
X D,XYZ,IX,F,FORCE,ESTIF,U,VECT,ISW)
LOGICAL NOPRINT,NPR,NPL
DIMENSION ESTIF(NSTF,NSTF),D(3,2,1,1),IX(NEL1,1),U(NDF,1),
X V(3,20),DU(3),XX(3),FORCE(NSTF,2)
COMMON/GAUS/LIM,SGAUSS(5,5),WGAUSS(5,5)
COMMON/LABELS/HEAD(12),O,IPG,XHED(3),UHED(6),XH,FX,UH,NSTR,FLAG(7)
COMMON/LOCALS/DUL(6,20),UL(6,20),UDL(6,20),HDDL(6,20)
COMMON/PRTPLT/NSIG(9),NPLT(9,2),NT,NSTEP,NUMPLT,NEDATA(20,3),NPR
COMMON/SHAP/XJAC,SHAPE(4,20),SG(3,3),SK(3,3),X(3,20)
COMMON/TAPES/ITP5,ITP6
COMMON/TIMIS/TIME,DT,DTP,NH,ISZH,C0,C1,C2,C3,C4,C5,C6,NCT
DATA SH,EH,BL,6HSTRESS,6HSTRAIN,6H
GO TO (1,2,3,4,5,4),ISW
CONTINUE
C.....
BAR STIFFNESS CHARACTERIZATION
READ(ITP5,1000) E,AA,RO
WRITE(ITP6,2100) E,AA,RO
D(1,1,MA) = E
D(1,2,MA) = AA
D(1,3,MA) = RO
C6 = 0.
RETURN
CONTINUE
2
RETURN
CONTINUE
3
COMPUTATION OF BAR STIFFNESS NO BENDING * ISOPARAMETRIC 1ST AND 2N
C.....
CONSTRUCT STIFFNESS AT INTEGRATION POINTS
EA = D(1,1,MA)*D(1,2,MA)
RA = D(1,3,MA)*D(1,2,MA)
DO 250 II = 1,NELM
SS = SGAUSS(II,NELM)
CALL LINE(SS,NDIM,NELM)
WJ=WGAUSS(II,NELM)
DVOL=EA*WJ/(XJAC**XJAC**XJAC)
C..... COMPUTE A LUMPED MASS MATRIX
IU = 0
RAA=RA**XJAC**WJ
DO 217 I = 1,NELM
AA=RAA**SHAPE(2,I)
DO 215 KK = 1,NDIM
FORCE(IU+KK,2) = FORCE(IU+KK,2) + AA
IU = IU + NDF
DO 250 KK = 1,NDIM
DO 250 LL = KK,NDIM
SGD=SG(KK,LL)*DVOL
II = KK
DO 240 I = 1,NELM
SHI=SHAPE(1,I)*SGD
L = I
IF(KK.EQ.LL) L = I
J1 = LL + NDF*(L-1)
DO 230 J = L,NELM

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ELM 1C
ELM 2C
ELM 4C
ELM 5C
ELM 8C
ELM 10C
ELM 11C
ELM 13C
ELM 14C
ELM 17C
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ELM 22C
ELM 23C
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000214 ESTIF(I1,J1)=ESTIF(I1,J1)+SHI*SHAPE(I,J)
000225 J1 = J1 + NDF
000231 J1 = I1 + NDF
000234 CONTINUE
000241 CONTINUE
C.....
000243 KK = NELM*NDF
000245 LL = KK - 1
000247 DO 360 I=1,LL
000250 K = I + 1
000253 DO 360 J = K, KK
000254 ESTIF(I,J) = ESTIF(I,J) + ESTIF(J,I)
000266 ESTIF(J,I) = ESTIF(I,J)
000302 IF(C5.EQ.0.0) RETURN
000304 DO 370 I = 1, NSTF
000306 ESTIF(I,I) = ESTIF(I,I) + C6*FORCE(I,2)
000324 RETURN
000325 CONTINUE
000325 E = D(.,1,MA)
000331 NELS = NELS
DO 400 NN = 1, NELS
000332 NPRINT=NPR
000334 IF(NSIG(NN).GT.0) NPRINT=.TRUE.
000341 NPL=.TRUE.
000342 IF(NPL.NN,1).GT.0) NPL = .FALSE.
000345 SS = SGAUSS(NN,NELS)
000351 CALL LINE(SS,NDIM,NELM)
C.....
COMPUTE STRESS AND STRAIN
DO 100 KK = 1, NDIM
000353 XX(KK) = 0.0
000350 DU(KK) = 0.
000361 DO 100 LL = 1, NELM
000362 XX(KK) = XX(KK) + X(KK,LL)*SHAPE(2,LL)
000372 DU(KK) = DU(KK) +UL(KK,LL)*SHAPE(1,LL)
000404 EPS = 0.
000405 DO 200 KK = 1, NDIM
000406 EPS = EPS + SK(KK,1)*DU(KK)/XJAC/XJAC
000416 SIG = E*EPS
000420 IF(NPRINT) GO TO 350
000421 MCT = MCT - 1
000423 IF(MCT.GT.0) GO TO 300
MCT=50
000425 WRITE(ITP6,2000) HEAD, TIME, IPG, ( XHED(I),XH, I=1,NDIM),BL,SH,BL,EH
000426 IPG = IPG + 1
000465 DO 600 I = 1, NELM
000467 AF = UJ*SHAPE(2,I)
000523 IF(.NOT.NPL) CALL PLDATA(NDIM,IPLT(NN,1),5HAXIAL,XX,SIG)
000542 IF(ISW.EQ.4) GO TO 400
000545 WT = WGAUSS(NN,NELS)*D(1,2,MA)
000553 WJ = WT*D(1,3,MA)*XJAC
000557 IU = 0
000560 DO 600 I = 1, NELM
000562 AF = UJ*SHAPE(2,I)
000565 EB = SHAPE(1,I)*SIG*WT/XJAC
000572 DO 610 FK = 1, NDIM

```

ELM 48C
ELM 49C
ELM 50C
ELM 51C
ELM 52C
ELM 54C
ELM 55C
ELM 56C
ELM 57C
ELM 58C
ELM 59C

ELM 60C
ELM 61C
ELM 64C

ELM 71C
ELM 72C
ELM 73C
ELM 74C
ELM 75C
ELM 76C
ELM 77C

ELM 79C
ELM 80C
ELM 81C
ELM 82C

ELM 83C
ELM 84C

ELM 87C
ELM 88C

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000573 615 FORCE(KK+IU,I) = FORCE(KK+IU,I) - AA*UDDL(KK,I) - BB*SK(KK,I)
000615 600 IU = IU + NDF
000621 400 CONTINUE
000624 5 RETURN
000625 1000 FORMAT(3F10.0)
000625 2000 FORMAT(1H1,12A6,E13.5,17X,4HPAGE,13//10H ELEMENT,4X,8HMATERIAL,
X 3X,5(2A5))
000625 2001 FORMAT(110,15.5X,A5,5E12.4)
000625 2100 FORMAT(37H0LINEAR ELASTIC MATERIAL, BAR ELEMENT//
X 5X,3HE =,E15.5,5X,6HAREA =,E15.5,5X,9HDENSITY =E15.5/)
000625 4000 FORMAT(A5,15,1P3E12.5,24X,110)
000625 END

```

ELM 91C

ELM 93C

ELM 96C

ELM 97C

ELM 98C

ELM 99C

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000006 SUBROUTINE LINE(S,NDIM,NEL)
C***** LINE ***** 7/02/73 ***** ISOPARAMETRIC LINES *****
C.... SHAPE FUNCTION ROUTINE FOR 1ST AND 2ND ORDER ISOPARAMETRIC LINES
C..... COMMON/SHAP/ XJAC,SHAPE(4,20),SG(3,3),SK(3,3),X(3,20)
C..... FORM SHAPE FUNCTIONS
      SHAPE(1,1) = -0.5
      SHAPE(2,1) = 0.5*(1.0-S)
      SHAPE(1,NEL) = 0.5
      SHAPE(2,NEL) = 0.5*(1.0+S)
      IF(NEL.EQ.2) GO TO 350
      SHAPE(1,2) = -2.0*S
      SHAPE(2,2) = 1.0-S*S
      K = NEL-1
      DO 100 I = 1,NEL,K
        SHAPE(1,I) = SHAPE(1,1) - 0.5*SHAPE(1,2)
        SHAPE(2,I) = SHAPE(2,1) - 0.5*SHAPE(2,2)
C..... FORM JACOBIAN MATRIX AND DETERMINANT
      DO 360 I = 1,3
        SK(I,1) = 0.0
      DO 400 J = 1,NDIM
        DO 400 J = 1,NEL
          SK(I,J) = SK(I,1) + X(I,J)*SHAPE(I,J)
          XJAC = 0.0
      DO 500 I = 1,NDIM
        XJAC = XJAC + SK(I,1)*SK(I,1)
      DO 500 J = 1,NDIM
        SG(I,J) = SK(I,1)*SK(J,1)
      XJAC = SORT(XJAC)
      RETURN
      END
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000105

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SUBROUTINE NLNORM(ISJ,MDEG,NN,U,DU,DF,F,PROP,NCT,NTT,IDEST,IBLK)
C***** NLNORM ***** 12/14/73 *****
DIMENSION U(1),DU(MDEG),DF(1),F(MDEG),IDEST(1)
COMMON/NORMS/ UNORM,DNORM,ANORM,CS,CSP,DP,DNP,IFL,XFLAG,ERR
GO TO (100,200,300,400), ISJ
100 CS=1.0
CSP=1.0
DP=0.75
IFL=0
XFLAG=.FALSE.
ERR=1.0E-3
RETURN
200 UNORM = 0.
DNORM = 0.
ANORM=0.
DO 500 N = 1,MDEG
DUN = DU(N)
UN = U(N)
IF (ISJ.EQ.3) UN = UN + DF(N) + DUN
ANORM=ANORM+UN*DUN
UNORM=UNORM+UN*UN
DNORM=DNORM+DUN*DUN
UNORM=SQRT(UNORM)
DNORM=SQRT(DNORM)
IF (NN.EQ.1) DNP=UNORM
CS=1.0
IF (UNORM*DNORM.NE.0.) CS=ANORM/(UNORM*DNORM)
IF (IFL.EQ.1) DP=0.75
IF (IFL.EQ.1.AND.XFLAG) CSP=CS
IFL=0
IF (DNORM.LE.0.5*UNORM) GO TO 550
IFL=1
DP=0.5*UNORM/DNORM
550 IF (DP*DNORM.GT.DNP) DP=DNP/DNORM
IF (DP.EQ.0.0) DP = 1.0
IF (CS*CSP.LT.0.) DP=DP/2.
IF (CS*CSP.GT.0.) DP=1.25*DP
DNP=DP*DNORM
CSP=CS
IF (ISJ.EQ.2) RETURN
DO 600 N = 1,MDEG
DF(N) = DF(N) + DU(N)
IF (NCT.EQ.0) DF(N) = DU(N)
DU(N) = DF(N)
RETURN
600
END

```

NLN 1C
NLN 2C
NLN 3C
NLN 4C
NLN 5C
NLN 6C
NLN 7C
NLN 8C
NLN 9C
NLN 10C
NLN 11C
NLN 12C
NLN 13C
NLN 14C
NLN 15C
NLN 16C
NLN 17C
NLN 18C
NLN 19C
NLN 20C
NLN 21C
NLN 22C
NLN 23C
NLN 24C
NLN 25C
NLN 26C
NLN 27C
NLN 28C
NLN 29C
NLN 30C
NLN 31C
NLN 32C
NLN 33C
NLN 34C
NLN 35C
NLN 36C
NLN 37C
NLN 38C
NLN 39C
NLN 40C
NLN 41C
NLN 42C
NLN 43C
NLN 44C
NLN 45C

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SUBROUTINE TSOLVE(NUMNP,NUMEL,NUMMAT,NDIM,NDF,NEN,NEL1,NSTF,NVEC,
1 NSIZV,NSICD,NSICD,IBLK,ISZA,NEQB,MAXBAN,NDEG,ID1,M8,IBUF,DS,
2 TYPE,D,ICOD,XYZ,F,IX,IDEST,VECT,FORCE,ESTIF,LD,A,DU,MAXB,H,U,DF,
3 NSEQ,TYPE)
C***** TSOLVE *****03/06/74 *****
LOGICAL CFLAG,FLAGC
LOGICAL NPR,VFLAG,DFLAG
REAL LABL(12)
DIMENSION TYPE(1),ICOD(1),XYZ(NDIM,1),F(NDF,1),IX(NEL1,1),
1 IDEST(NDF,1),VECT(NSIZV,1),FORCE(NSTF,2),ESTIF(NSTF,NSTF),
2 A(NENB,1),DU(NDF,1),H(IBUF),U(NDF,1),DF(1),DS(M8),TYNE(7),LD(1)
COMMON/CONTACT/FLAGC,CFLAG,LIST,ICLIST(10),ICDEG(10),RUI,RUZ,
C RM(2,10)
COMMON/LABELS/HEAD(12),0,IPG,XHED(3),UHED(6),XH,FH,UH,NSTR,FLAG(7)
COMMON/LOCALS/ DUL(6,20),UL(6,20),UDL(6,20),UDDL(6,20)
COMMON/NORMS/ UNORM,DNORM,ANDRM,CS,CSP,DP,DNP,IFL,XFLAG,ERR
COMMON/PERTLT/ NSIG(9),NPLT(9,2),NT,NSTEP,HUMPLT,NEDATA(20,3),NPR
COMMON/SHAP/ XJAC,SHAPE(4,20),SG(3,3),SK(3,3),X(3,20)
COMMON/TAPES/ ITP5,ITP6
COMMON/TIMHS/ TIME,DT,DTP,NH,ISZH,C0,C1,C2,C3,C4,C5,C6,NCT
COMMON/TAPE/ ITP13,ITP14,ITPD,ITWR
DATA AMORD1,AMORD2/8H F13.4,8H E13.4,/,AMORD3/8H F12.2,/,
DATA VH,AH,AMORD4,AMORD5/6H VEL.,6H ACCEL.,6H(/12X.,8H9E13.4)/
DATA LABL(1)/8H (112, /
DATA CFLAG,FLAGC/,.FALSE,.,.FALSE,./
C..... SUBROUTINE TO PERFORM IMPLICIT INTEGRATION WITH 1 OR 3 INIT. COND.
DO 2 N=1,NEN
DO 2 J=1,NDF
DUL(J,N)=0.
UL(J,N)=0.
UDL(J,N)=0.
UDDL(J,N)=0.
DO 3 I=2,7
NSIG(I)=1
IF(1BLK.EQ.0) GO TO 4
KU = NUMNP*NSICD
REWIND 7
WRITE(7) DS
C..... SET THE OUTPUT LABELS FOR DISPLACEMENT PRINTS
4 HEAD(12) = 6H TIME=
IFLG =-NSICD
MDEG = NSICD
MDEG2 = MDEG + MDEG
I = 2
DO 5 N = 1,NDIM
LABL(I) = AMORD1
I = I + 1
L = HDIM + 1
DO 7 N = L,9
LABL(I) = AMORD2
I = I + 1
LABL(I) = AMORD4
LABL(I+1) = AMORD5

```

TSO 1C
TSO 2C
TSO 3C
TSO 4C
TSO 5C
TSO 6C
TSO 8C
TSO 9C
TSO 10C
TSO 11C
TSO 12C
TSO 13C
TSO 14C
TSO 15C
TSO 16C
TSO 17C
TSO 18C
TSO 20C
TSO 21C
TSO 22C
TSO 23C
TSO 24C
TSO 25C
TSO 26C
TSO 27C
TSO 28C
TSO 29C
TSO 30C
TSO 31C
TSO 32C
TSO 33C
TSO 34C
TSO 35C
TSO 36C
TSO 37C
TSO 38C
TSO 39C
TSO 40C
TSO 41C
TSO 42C

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000152      DTP = 0.
000153      TIME = 0.
000154      MB = MAXBAN + 1
000155      NE1 = NEN + 1
000156
C..... INITIALIZE INCREMENTAL FORCE VECTOR
6      DO 6 N = 1, NSICD
      DF(N) = 0.0
      DO 10 N = 1, NUMEL
10     IX(NEI,N) = 0
C..... PUT INITIAL DATA ON TAPE
      ISZH = IBUF
      ITRD = ITP13
      ITUR = ITP14
      REWIND ITRD
      IF( (IBLK.GT.0) WRITE(ITRD) ((CU(I,J), I=1, NDF), J=1, KU)
11     DO 11 N = 1, ISZH
      H(N) = 0.
      NTB = (NH+ISZH-1)/ISZH
      WRITE(ITRD) H
      IF(NTB.LE.1) GO TO 13
      NT = NTB*2
      DO 12 N = 1, NT
12     WRITE(ITRD) H
13     REWIND ITRD
      IF( (IBLK.GT.0) READ(ITRD) DM
      NEP = NUMNP + NUMNP
      C6 = 0.0
      NSTEP = 0
      NST = 0
      NUMPLT = 0
      PROP = 0.
      CFLAG = .FALSE.
      DO 900 M = 1, NSEQ
      READ(ITP5,1000) DT, NTS, INT, NNI, NNF, NEI, NEF, NPROP, NFORC, CO, DM, NT, NL
      IF( (NICD, GT.1) C6 = DM
      IF( (NL, NE.0) NST = NL
      NTT = IABS(NST)
      DFLAG = .FALSE.
      IF( (NTT, NE.0) DFLAG = .TRUE.
      NSTEP = NSTEPAINTS
      IF( (M, EQ.1) NUMPLT = NT
      IF( (INT, LE.0) INT = 1
      WRITE(ITP6,2000) O, HEAD, TIME, IPG, DT, NTS, INT, NNI, NNF, NEI, NEF
      IF( (DFLAG) WRITE(ITP6,2002) NST
      IPG = IPG + 1
      IF( (NICD, EQ.1) GO TO 1
      IF( (DT, EQ.0.0) GO TO 901
      CALL UPDATE(3)
      IF( (NPROP, GT.0) PROP = PROPLD(TIME, NPROP)
      IF( (NFORC, GT.0, AND, NPROP, GT.0) WRITE(ITP6,4001)
C..... READ STRESS PLOT INFORMATION
      IF( (NUMPLT, LE.0, OR, M, GT.1) GO TO 5AC
      REWIND 12
      WRITE(ITP6,2005)

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000555 DO 501 N=1,NUMPLT
000562 READ(ITP5,1006) (NEDATA(N ,I),I=1,2)
000600 WRITE(ITP6,2006) N,(NEDATA(N ,I),I=1,3)
000632 CALL PLZERO
000633 VFLAG = .FALSE.
000634 DO 500 NT = 1,NTS
000641 PPROP = PROP
000643 PROP = 1.0
000644 IF(DFLAG)CALL NLNORM(4,MDEG,NT,U,DF,DU,F,PPROP,NCT,NTT,IDEST,IBLK)
000657 IF(NPROP.GT.0) PROP = PROPLD(TIME+DT,0)
000701 IF(NFORC.GT.0) CALL RESET(-NFORC,NUMMP,NDF,F)
000713 NCT = 0
14 REWIND ITUR
000714 IF(IBLK.EQ.0) GO TO 15
000716 IF(.NOT.VFLAG) GO TO 20
000720 DO 21 I = 1,NDEG
000721 DF(I) = 0.
000723 GO TO 20
15 DO 19 N = 1,NUMMP
000730 DO 18 K = 1,NDF
000732 J = IDEST(K,N)
000733 IF(J.EQ.0) GO TO 18
000740 IF(VFLAG) GO TO 17
000741 DO 16 I = 1,MAXBAN
000743 A(J,I) = 0.0
000744 DF(J) = F(K,N)*PROP
17 CONTINUE
18 CONTINUE
19 NH = 1
20 IF(NTB.GT.1) READ(ITRD) H
001017 IF(NCT.GT.0) GO TO 44
C..... OUTPUT THE SOLUTION VECTOR FOR THE CURRENT TIME
001032 IF(NNI.EQ.0.OR.(NT-1)/INT)*INT.NE.NT-1) GO TO 40
001033 MCT = 0
001034 DO 30 N = NNI,NNF
001036 MCT = MCT + 1
001040 IF(MCT.GT.0) GO TO 31
X IF(NICD.EQ.1) WRITE(ITP6,2001) O,HEAD,TIME,IPG,PPROP,M,NT,
X (XHED(I),XH,I=1,NDIM),(UHED(I),UH,I=1,NDF)
X IF(NICD.NE.1) WRITE(ITP6,2001) O,HEAD,TIME,IPG,PPROP,M,NT,
X (XHED(I),XH,I=1,NDIM),(UHED(I),UH,I=1,NDF),(UHED(I),UH,I=1,NDF)
X (UHED(I),UH,I=1,NDF)
IPG = IPG + 1
MCT = 50
31 IF(NICD.EQ.1)WRITE(ITP6,LABELN,(XYZ(I,N),I=1,NDIM),(U(I,N),
X ,I=1,NDF)
X IF(NICD.NE.1)WRITE(ITP6,LABELN,(XYZ(I,N),I=1,NDIM),(U(I,N),I=1,NDF
X ),(U(I,NUMMP-N),I=1,NDF),(U(I,NEP+N),I=1,NDF)
30 CONTINUE
C..... UPDATE (I) * * FOR DYNAMIC SOLUTIONS ONLY
001420 IF(NICD.GT.1) CALL UPDATE(I,MDEG,NICD,U(I,NUMMP+1),U(I,NEP+1),
X DU,F,DF,IDEST,PROP)
001423 IF(IBLK.GT.0) WRITE(ITUR) ((U(I,J),I=1,NDF), J=1,KU )
001457 IP = 0
001512

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TSO 97C
TSO 98C
TSO 99C
TSO100C
TSO101C
TSO102C
TSO103C
TSO104C
TSO105C
TSO106C
TSO107C
TSO108C
TSO109C
TSO110C
TSO111C
TSO112C
TSO113C
TSO114C
TSO115C
TSO116C
TSO117C
TSO118C
TSO119C
TSO120C
TSO121C
TSO122C
TSO123C
TSO124C
TSO125C
TSO126C
TSO127C
TSO128C
TSO129C
TSO130C
TSO131C
TSO132C

TSO136C
TSO.37C

TSO139C
TSO140C
TSO141C
TSO142C
TSO143C

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001513 TEMP = 0.
001514 DO 400 N = 1, NUMEL
001515 NPR = .TRUE.
001516 DO 43 I = 1, 9
001520 NPLT(I,1) = 0
001521 NPLT(I,2) = 0
001524 IF(NCT.GT.0) GO TO 46
001526 IF(N.GE.NEI.AND.N.LE.NEF) NPR = .FALSE.
001540 IF(NUMPLT.LE.0) GO TO 46
001542 DO 45 I = 1, NUMPLT
001543 IF(NEDATA(I,1).NE.N) GO TO 45
001546 J = NEDATA(I,2)
001547 NPLT(J,1) = I
001551 NPLT(J,2) = NEDATA(I,3)
001553 CONTINUE
001556 MA = MOD(IX(NEL1,N),100)
001567 IF(MR.LE.0) MRR = IX(NEL1,N)/1000
001600 IF(MR.LE.0) MR = MRR
001603 DO 60 I = 1, NSTF
001605 FORCE(I,1) = 0.
001612 FORCE(I,2) = 0.
001616 LD(I) = 0
001620 IF(MR.NE.MRR.OR.VFLAG) GO TO 60
001625 DO 50 J = 1, NSTF
001627 ESTIF(I,J) = 0.0
001640 CONTINUE
001643 L = 0
001644 DO 110 I = 1, NEN
001645 K = IX(I,N)
001653 DO 90 J = 1, NDIM
001654 X(J,I) = 0.
001661 IF(K.EQ.0) GO TO 120
001663 NEL = I
001665 DO 100 J = 1, NDIM
001666 X(J,I) = XYZ(J,K)
001701 DO 110 J = 1, NDF
001702 L = L + 1
001704 DUL(J,I) = DU(J,K)
001713 UL(J,I) = U(J,K)
001722 IF(NCT.GT. 0) UL(J,I) = UL(J,I) + DUL(J,I)
001731 IF(NICD.EQ.1) GO TO 110
001734 UDL(J,I) = U(J,K+NUMNP)
001744 UDDL(J,I) = U(J,K+NEP)
001754 IF(NCT.GT.0) UDDL(J,I) = UDDL(J,I) + C6*DUL(J,I)
001765 LD(L) = IDEST(J,K)
002091 DM = TYPE(MA)
002004 IF(DM.NE.TEMP) MCT = 0
002007 TEMP = DM
002010 CALL TICTOC(TIME,6)
C..... COMPUTE ELEMENT STRESSES AND UPDATE FORCES
002013 CALL ELM1B(N,NA,NDIM,NDF,NEL,NEL1,NSTF,NSIZV,NVEC,MCT,DM,D,XYZ,
002047 X,IX,H,FORCE,ESTIF,U,VECT,6)
CALL TICTOC(TIME,5)
C..... FORM STIFFNESSIF NEEDED FOR THE NEXT TIME STEP

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T50144C
T50145C
T50146C
T50147C
T50148C
T50149C
T50150C
T50151C
T50152C
T50153C
T50154C
T50155C
T50156C
T50157C
T50158C
T50159C
T50160C
T50161C
T50162C
T50163C
T50164C
T50165C
T50166C
T50167C
T50168C
T50169C
T50170C
T50171C
T50172C
T50173C
T50174C
T50175C
T50176C
T50177C
T50178C
T50179C
T50180C
T50181C
T50182C
T50183C
T50184C
T50185C
T50186C
T50187C
T50188C
T50189C
T50190C
T50191C
T50192C
T50193C
T50194C
T50195C
T50196C
T50197C

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002052      IF(.NOT.VFLAG.AND.NR.EQ.MPR).OR.(VFLAG.AND.IX(NE1,N).EQ.1))
XCALL ELMLIB(N,MA,NDI,LD,NDF,NEL,NEL1,NSTF,NSIZY,NVEC,NCT,DM,D,XYZ,
X IX,H,FORCE,ESTIF,U,VECT,3)
002131      IF(MOD(IX(NEL1,N),1000)/100.GT.0.AND..NOT.VFLAG)
X CALL PRMAT(HEAD,IPG,N,NSTF,ESTIF,FORCE,LD,NSTF,0)
C.....  MODIFY FOR THE DISPLACEMENT B.C.
002165      CALL MODIFY(NDF,NEL,NEL1,NEL,IBLK,NSTF,PROP,IX,ICOD,F,FORCE,ESTIF,
X N)
002205      IF(VFLAG) GO TO 300
002212      IF(IBLK.EQ.0) CALL COREIN(A,DF,NEOR,ESTIF,FORCE,LD,NSTF)
002231      IF(IBLK.GT.0) WRITE(7) ESTIF,FORCE,LD
002253      IF(.NOT.FLAGC.OR.CFLAG) GO TO 400
002255      L = 1
002267      DO 130 I = 1,NEL
002270      K = IX(I,N)
002276      DO 130 J = 1,NDF
002277      DU(J,K) = DU(J,K) + FORCE(L,2)
002311      L = L + 1
002317      GO TO 400
002320      CONTINUE
C.....  ADD THE FORCE TO THE SOLUTION FOR A RESOLVE
002320      DO 310 K = 1,NSTF
002322      J = LD(K)
002325      IF(J.GT.0) DF(J) = DF(J) + FORCE(K,1)
002335      CONTINUE
002340      CALL TICOC(TYME,3)
002351      IF(FLAGC.AND..NOT.CFLAG) CALL CONTAC(2,IX,NEL1,NDF,NUMNP,DU)
002364      CFLAG = .TRUE.
002365      IF(NTB.GT.1) WRITE(ITWR) H
002405      IF(IBLK.GT.0) WRITE(ITWR) (IX(NEL1,N),N=1,NUMEL)
002436      IF(.NOT.VFLAG) CALL SOLVEQ(NUMNP,NUMEL,NDF,IP1,M8,MAXBAN,9,NSTF,
I ISZA,NEOB,IBLK,A,DU,DF,IDEST,FORCE,ESTIF,LD,MAXB,NDEG)
I IF(VFLAG) CALL RESVEQ(NUMNP,NDF,M8,MAXBAN,ISZA,NEOB,IBLK,A,DF,DF,
I IDEST,IDEST,MAXB,IFLG)
CALL TICOC(TYME,4)
C.....  UPDATE THE SOLUTION
002527      IF(NT.EQ.NTS.AND.M.EQ.NSEQ) GO TO 900
002542      IF(NCT.GT.0) GO TO 410
002544      DTP = DT
002545      I = ITRD
002546      ITRD = ITWR
002547      ITWR = I
002550      IF(IBLK.GT.0) BACKSPACE ITRD
002554      IF(IBLK.GT.0) READ(ITRD) (IX(NEL1,N),N=1,NUMEL)
002605      REWIND ITRD
002607      IF(IBLK.GT.0) READ(ITRD) ((U(I,J),I=1,NDF),J=1,KU)
002642      IF(DFLAG)CALL NLNORM(3,NDEG,NT,U,DF,DU,F,PPROP,NCT,HTT,IDEST,IBLK)
002666      VFLAG = .TRUE.
002667      IF(NST.LT.0) VFLAG = .FALSE.
002672      NCT = NCT + 1
002674      IF(.NOT.DFLAG) GO TO 700
002675      WRITE(6,6000) NCT,UNDRN,DNORM,ANORM
002711      IF(NCT.LT.NTT.AND.DNORM.GT.UNDRN) GO TO 14
002730      IF(DFLAG) VFLAG = .FALSE.

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002733 C..... UPDATE (2)
X CALL UPDATE(2,MDEG,NICD,U,U(1,N*TRIP+1),U(1,NER+1),DU,F,DF,IDEST,
X PROP)
X IF (FLAGC) CALL CONTACT(3,IX,NEL1,NDF,NUMIP,DU,U,U(1,NUMIP+1),
X U(1,NER+1))
003016 MR = MR - 1
003020 TIME = TIME + DT
003024 CONTINUE
003027 IF (NUMPLT.GT.0) CALL PLOTGO(NDIM,HEL,DM,U)
003040 CALL TICTOC (TIME,6)
003043 WRITE(ITP6,2033) TIME
003052 RETURN
003053 WRITE(ITP6,2030)
003057 IPG = 0
003060 RETURN
C..... FORMATS
003061 FORMAT(F10.0,8I5,2F10.0,2I5)
003061 FORMAT(3I5)
003061 FORMAT(A1,12A6,E13.5,17X,4HPAGE,I4//5X,23HTIME DEPENDENT SOLUTION
X//10X,17HTIME INCREMENT = 1PE15.4/
X 10X,17HNUMBER OF STEPS = 15/
X 10X,17HPRINT INTERVAL = 15/
X 10X,14HPRINT NODES ,15,3H TO,15/
X 10X,14HPRINT ELEMENTS,15,3H TO,15)
003061 2001 FORMAT(A1,12A6,E13.5,17X,4HPAGE,I4//5X,12HMODAL VALUES,5X,
X 27HPROPORTIONAL LOAD FACTOR = ,F8.3,5X, 8HSEQUENCE,14,5X,
X 9HTIME STEP,14//13H NODAL POINT,9(1X,2A6),(12X,9(1X,2A6)) )
003061 2002 FORMAT(10X,17HITERATION LIMIT =,15/1X)
003061 2005 FORMAT(4X,50H DESCRIPTION OF STRESS EVOLUTION PLOTS TO BE MADE, /
X 5X,10H PLOT NO. 5X,8H ELEMENT ,7X, 9H XYZ-CODE,6X,9H SIG-CODE/)
003061 2006 FORMAT(4(7X,15.3X))
003061 2030 FORMAT(5X,53H**FATAL ERROR 33** TIME STEP ZERO FOR DYNAMIC PROBLEM
X/1X)
003061 2033 FORMAT ( , 22X,12HELAPSED TIME//10X,25HINPUT PROPERTIES AND MESH
,F10.3/10X,25HCHECK AND PLOT INPUT DATA,F10.3/10X,14HFORM STIFFNESS
X,F21.3/10X,21HSOLUTION OF EQUATIONS,F14.3/10X,15HOUTPUT STRESSES,
XF20.3/10X,27HIMPLICIT ALGORITHM + DISPL.,F0.3/10X,19HTOTAL TIME/IM
XP,ICIT,F16.3/1X)
003061 4001 FORMAT(5X,69H**WARNING 3** BOTH THE PROPORTIONAL LOADING AND FOR
XCE ARE BEING RESET ON EACH TIME STEP )
003061 6000 FORMAT(5X*INCT*15,* UNORM*E15.5* DHORM*E15.5* ANORM*E15.5)
003061 END
TS0245C
TS0246C
TS0247C
TS0248C
TS0249C
TS0250C
TS0251C
TS0252L
TS0253C
TS0254C
TS0255C
TS0256C
TS0257C
TS0258C
TS0259C
TS0260C
TS0261C
TS0262C
TS0263C
TS0264C
TS0265C
TS0266C
TS0267C
TS0270C
TS0271C
TS0272C
TS0273C
TS0274C
TS0275C
TS0276C
TS0277C
TS0278C
TS0279C
TS0280C
TS0281C
TS0282C
TS0242C

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000016 SUBROUTINE UPDATE(ISW,NDEG,NICD,U,UD,UDD,DU,F,DF,IDEST,PROP)
000016 C***** UPDATE ***** 12/14/73 *****
000016 DIMENSION U(1),DU(1),F(1),DF(1),IDEST(1),UD(1),UDD(1)
000016 COMMON/TAPES/ ITP5,ITP6
000016 COMMON/TIMHIS/ TIME,DT,DTP,NH,ISZH,C0,C1,C2,C3,C4,C5,C6,NCT
000016 GO TO (100,300,500), ISW
000024 UPDATE (1) * * PREUPDATE OF ACCELERATIONS FOR DYNAMIC SOLUTIONS.
000026 DO 200 N = 1,NDEG
000035 UDD(N) = C4*UDD(N) - C5*UD(N)
000035 RETURN
000035 C..... UPDATE (2) * * UPDATE THE SOLUTION FOR GENL. DISPL, VEL, ACCL.
000037 DO 400 N = 1,NDEG
000043 DU(N) = DF(N)
000043 IF (IDEST(N).EQ.0) DU(N) = F(N)*PROP - U(N)
000053 TEMP = DU(N)
000056 IF (NICD.EQ.1) GO TO 400
000056 P = JDD(N)
000061 UD(N) = C1*UD(N) + C2*P + C3*TEMP
000065 UDD(N) = P + C6*TEMP
000073 U(N) = U(N) + TEMP
000100 RETURN
000100 C..... UPDATE(3) * * SET INTEGRATION CONSTANTS
000102 C6 = C6 + 0.5
000104 IF (C0.EQ.0.0) C0 = 0.25
000104 WRITE (ITP6,2002) C6,C6
000114 C5 = 1./C0/DT
000116 C4 = 1. - 0.5/C0.
000121 C3 = C6*C5
000123 C2 = (1. - C6/C0/2.)*DT/C4
000127 C1 = 1. - C6/C0 + C2*C5
000134 C6 = C5/DT
000137 RETURN
000137 FORMAT(10X,30HNEWMARK INTEGRATION PARAMETERS/
000137 C10X17HBETA VALUE = F6.3/
000137 C10X17HGAMMA VALUE = F6.3)
000137 END
000137
UPD 1C
UPD 2C
UPD 3C
UPD 4C
UPD 5C
UPD 6C
UPD 7C
UPD 8C
UPD 9C
UPD 10C
UPD 11C
UPD 12C
UPD 13C
UPD 14C
UPD 15C
UPD 16C
UPD 17C
UPD 18C
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UPD 34C
UPD 35C

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