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Authors
Kalbfleisch, George R
Alvarez, Luis W
Barbaro-Galtieri, Angela
et al.
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Lawrence Radiation Laboratory
Berkeley, California
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OBSERVATION OF A NONSTRANGE MESON OF MASS 959 MeV

George R. Kalbfleisch, Luis W. Alvarez, Angela Barbaro-Galtieri, Orin I. Dahl, Philippe Eberhard, William E. Humphrey, James S. Lindsey,

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Department of Physics and-Lawrence Radiation Laboratory, University of California, Berkeley, California

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We present here evidence chowing the existence of a nonstrange meson of mass 959 MeV .

In the current experiment, the 72-in. hydrogen bubble chamber was expased to a separated beam of $2.45 \cdots, 2.63 \cdots$, and $2.70-\mathrm{BeV} / \mathrm{ck}$ mesons from the Bevatron. Approximately 370000 pictures were taken to date; approximately 300000 have been scanned. The reactions of interest in this paper are

$$
\begin{align*}
& \mathbf{K}^{0}+\mathbf{p} \rightarrow \Lambda^{0}+\mathbf{N M}  \tag{}\\
& \Lambda^{0} \cdot \pi^{+} \pi^{-}  \tag{2}\\
& \Lambda^{0^{0}} \pi^{+} \pi^{0^{\circ}} \pi^{-}  \tag{3}\\
& \Lambda^{0} \mathbb{Z}^{+} \mathbb{Z}^{-}+M M  \tag{4}\\
& \Lambda^{0} \pi^{+} \pi^{+} \pi^{-} \pi^{-}  \tag{5}\\
& \Lambda^{0} \pi^{+} \pi^{+} \pi^{0} \pi^{-} \pi^{-}  \tag{6}\\
& \Lambda^{0} \text { r }^{+} \pi^{+} \pi^{-} \pi^{\circ}+M M \tag{7}
\end{align*}
$$

and

$$
\begin{equation*}
\Lambda^{\circ} 3 \pi^{+} 3 \pi^{-} \tag{8}
\end{equation*}
$$

These reactions are found in the topologies of $\mathrm{a} V$ and $0,2,4$, and 6 prongs. At this time we have measured the $V$-four-prong and $V$-six-prong events in 250000 pictures, the $V$-two-prongs in 135000 pictures, and the $V$-zero-prongs in 100000 pictures.

Figure 1a shows clearly the exiatence of the $959-\mathrm{MeV}$ meson as an enhancement in the $\pi^{+} \pi^{+} \pi^{0} \pi^{-} \pi^{*}$ apectrum from reaction (6). The mass is $959 \pm 2 \mathrm{MeV}$ and the full width is $\Gamma \leqq 12 \mathrm{MeV} .{ }^{1}$ We observe 35 events in the interval $0.86 \leqslant \mathrm{M}^{2}(5 \pi) \leqslant 0.98 \mathrm{BeV}^{2}$. The background is estimated to be less than $10 \%$ of the peak. These events axe produced mainly with a momentum transfer $\Delta^{2} p, \Lambda$ less than $0.5 \mathrm{BeV}^{2} 2$ We have used the momentum transfer only as a means to scparate other decay modes of this meson from the large background in reactions 11 through 5). The distribution of the four $\pi^{+} \pi^{\circ} \pi^{-}$combinations for each of the 25 events in the peak (Fig. ib) clearly shows the presence of the $548-\mathrm{MeV} \eta$ meson. $\cdots$ eh of the 35 events has at least one $\pi^{+} \pi^{\circ} \pi^{-}$triplet at the $\eta$ mass. ${ }^{3}$ We conclude Ex: the 959-MeV mesondecaye into $\pi^{+} \pi^{-} \eta^{-}$.

We now turn our attention to other possible decay modes of this meson. The neutral MM and the $\pi^{+} \pi^{-} \mathrm{MM}$ distributions (at $2.45 \mathrm{BeV} / \mathrm{C}$ only) from reactions (1) and (4) (Fig. 2, a and b) show enhancements at 959 MeV . In addition, the celection of $M M \approx \eta$ in the low $-\Delta^{2} \Lambda^{0}{ }^{+} \dagger^{+} \pi^{-} M M$ events gives a practically clean aample for the $\pi^{+} \pi^{-} \eta$ where the $\eta$ decays into all neutrals (Fig. 2 b ). This last eelection gives a sample of 26 events with $0.86 \leqslant \mathrm{M}^{2}\left(\pi^{+} \pi^{-} \mathrm{MM}\right) \leqslant 0.98 \mathrm{BeV}^{2}$ and $0.27 \leqslant \mathrm{MM}^{2} \leqslant 0.33 \mathrm{BeV}^{2}$ for all the 135000 pictures in which these events were measured. ${ }^{4}$ No appreciable decay into $\pi^{+} \pi^{-}, \pi^{+} \pi^{0} \pi^{*}$. and $2 \pi^{+} 2 \pi^{-}$is observed (Fig. 2, c and d). In addition, no decay into $2 \pi^{+} 2 \pi^{\circ} 2 \pi^{-}$or $3 \pi^{+} 3 \pi^{-}$is observed. ${ }^{5}$ We note that the decay rate into all neutrals ( $\leqslant 20 \Lambda^{\circ}+\mathrm{MM}$ events above background in Fig. 2a) is comparable to the decay rate into $\pi^{+} \pi \pi^{-} \eta$ (four of the $35 \Lambda^{\circ} 5 \pi$ events occur in the $2.45-\mathrm{BeV} / \mathrm{c}$ part of the sample in addition to the 14 $\Lambda^{\circ} \pi^{+} \pi^{-} M M$ events in Fig. 2b).

We now look at the properties of the $\pi^{+} \pi^{-} \eta$ decay mode in an attempt to determine the quantum numbers of this meson. When we construct the Dalitz plot of $M^{2}\left(\pi^{+} \eta\right)$ versus $M^{2}\left(\pi^{-} \eta\right)$ and their projections (Fig, 3a) for the $61 \Lambda^{\circ} \pi^{+} \pi^{-} \eta$ events
(35 $\Lambda^{0} 5 \pi$ and $\left.26 \Lambda^{0} \pi^{+} \nabla^{\infty} \mathrm{MM}\right)^{6}$ we observe no particular structure in the $\pi^{ \pm} \eta^{\prime}$ distributions. Tho distribution of points in the Dalitz plot is consistent with uniformity, The $20-\mathrm{MeV}$ epread in the observed full width of the meson maes smears out the points appreciably (envelopes labeled 950 and 970 MeV in Fig. 3a). However, the distributiom in $M^{2}\left(\pi^{+} \pi^{*}\right.$ outside the $\eta$ ) (Fig. 3b) is not appreciably affected by the resolution. The ${ }^{+} \pi^{-"}$ distribution appears to be enhanced around 360 MeV . We now consider the $\pi T \eta$ system an dipion and $\eta$, with no final-state interaction between the $\eta$ and the pions outside it, and denote the angular momentum of the dipion as \& and that between tho dipion and the eta as L. All J states except $0^{+}$are allowed. For $J=0$ or 1, the possible spins and parities are $0^{-}(\ell=L=0) ; 1^{+}(\ell=0, L=1$, or $\ell=1, L=0)$; and $1^{-}(\ell=L=1)$. For $0^{-}$and $1^{+}$, there can be no correlation between the direction of the $\eta$ and that of one of the pions as viewed in the dipion rest frame. ${ }^{7}$ : For $1^{-}$, a $\sin ^{2} \theta_{\pi \eta}$ correlation is $\because ;$ required. ${ }^{8}$ The distribution in $\cos \theta_{\pi \eta}$ as observed in the dipion rest frame is essentially isotropic, disagreeing with the $\sin ^{2} \theta$ prediction of a $1^{-}$state. (Fig. 3c). The ambiguity of the choice of the $\eta$ in the $\Lambda 5 \pi$ events does not alter this conclusion. ${ }^{9}$. Thus, we conclude that the $959-\mathrm{MeV}$ meson is probably not a vector particle.

The isospin of the meson is either $T=0$ or $T=1$ because of its production from an initial $K^{-p} p$ systom in association with a $\Lambda^{0}$. Absence of appreciable decay into $\pi^{+} \pi^{-} \cdot \pi^{+} \pi^{0} \pi^{-}, 2 \pi^{+} 2 \pi^{-}, 2 \pi^{+} 2 \pi^{0} 2 \pi^{-}$and $3 \pi^{+} 3 \pi^{-}$. implies the absence of appreciable decay into $n \pi^{\circ}$. If the decay into $\pi^{+} \pi \stackrel{\pi}{\eta}$ occurs strongly, then the neutral mode $\pi^{0} \pi^{0} \eta$ determines the lsospin to be $T=0$. However, a "zerol" width for this meson is not excluded ( $\Gamma \leqslant 12 \mathrm{MeV}$ ). ${ }^{10}$ The predicted branching ratios for a strongly decaying $T=0$ ाr $\eta$ eystem are in fair agreement with the observed ratios. ${ }^{11}$ In addition, observation of six $e^{+} e^{-}$conversion pairs and four Compton electrons in association with the $\sim 50 \Lambda^{0}+\mathrm{MN}$ events in the $959-\mathrm{MeV}$ region
points to a high multiplicity of gammas in the all-neutral decay, consistent with $\pi^{0} \pi^{0} \eta^{12}$

In summary, we have observed a meson of mase $959 \pm 2 \mathrm{MeV}$, ${ }^{13}$ full width ${ }^{\text {! }}$ $\Gamma \leqslant 12 \mathrm{MeV}$, and isospin $T=0$ or 1, which decays into $\pi^{+} \pi \pi^{-} \eta$. No appreciable decay into two, threc, four, or six pions is observed. The angular correlation of the $\boldsymbol{\eta}$. and the $\pi^{+}$In the dipion rest system is not $\sin ^{2} \theta_{\pi \eta^{\circ}}$ makingsunifkely the $J^{P}=1^{-}$2ssignment. If we assume a strong decay (nonzero-width), the meson has isospin T $=0$ and G parity $a+1$, with $J^{P}$ probably $0^{-}$or $1^{+}$. ${ }^{14}$ However, electromagnetic decay cannot be ruled out. The path length at $2.45 \mathrm{BeV} / \mathrm{c}$ is appr $3 \mathrm{x}-$ imately $0.6 \mu \mathrm{~b}$ per event. The $\sim 44 \Lambda \pi \pi \eta$ events at $2.45 \mathrm{BeV} / \mathrm{c}$ yield a cross section of about $25 \mu \mathrm{~b}$. The average polarization of the $\Lambda^{0}$ produced in association with this meson is $a_{\Lambda}{ }^{\dagger}=0 \pm .2$.

We wish to acknowledge helpful discussions with Professors Murray Gell- ? Mann and Donald F. Miller and Dr. Nicola Cabibbo. We are indebted to the operators of the 72-in. bubble chamber and the Bevatron for their skill and patience. Furthermore, we thank our scanning and measuring stafis for their untiring efforts, without which this experiment would not have been possible.

## FOOTNOTES AND REFERENCES

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1. The observed full width is $\sim 20 \mathrm{MeV}$, and the resolution is $\sim 16 \mathrm{MeV}$ for the fitted $\Lambda^{0} \pi^{+} \pi^{+} \pi^{0} \pi^{-} \pi^{-}$reaction.
2. The square of the momentum transfer from the proton to the lambda is defined by $\Delta_{p, \Lambda}^{2}=-\left(E_{p}-E_{\Lambda}\right)^{2}+\left(p_{p}-p_{\Lambda}\right)^{2}$.
3. Of the 35 events with $0.86 \leqslant \mathrm{M}^{2}(5 \pi) \leqslant 0.98 \mathrm{BeV}^{2}$. twelve have one, eighteen have two, three have three, and two have four $\pi^{+} \pi^{0} \pi^{-}$triplets near the $\eta$ mass. The $\eta$ mass is defined by $0.282 \leqslant \mathrm{M}^{2}\left(\pi^{+} \pi^{0} \pi^{-}\right) \leqslant 0.322 \mathrm{BeV}^{2}$. This distribution is consistent with a randorn coincidence based on the phase apace for the $\pi \pi \eta$ decay (see Fig. 1b) and the resolution.
4. The resolution for the MM and the $\pi^{+} \pi^{-} \mathrm{MM}$ is $\sim 25 \mathrm{MeV}$, which is essentially due to the $3 \%$ momentum spread in the beam. The observed full width in these MM channels of $\sim 30 \mathrm{MeV}$ gives a full width for the 959 MeV meson of $\leqslant 20 \mathrm{MeV}$.
5. Three $\Lambda^{0}$-six-prong and $28 \Lambda^{\circ}$-four-prong events are consistent with $\Lambda^{0} 6 \pi$ or $\Sigma^{0} 5 \pi$ in the total sample ( $\sim 300000$ scanned pictures). All 31 of these eventshave $\mathrm{M}^{2}(6 \pi$ or $\gamma 5 \pi) \geqslant 1.1 \mathrm{BeV}^{2}$.
6. In the $\Lambda^{0} 5 \pi$ events we chose as the $\eta$ that $\pi^{+} \pi^{0} \pi^{-}$triplet whose mass squared is closest to $M_{\eta}^{2}=0.300$. This choice does not always pick the correct triplet because of the finite resolution. About half of the ambiguities should be chosen correctly statistically, giving 48 (of which 38 are unambiguous) correctly and 13 incorrectly chosen. ${ }^{3}$
7. The $\mathrm{J}^{\mathrm{P}}=0^{-}, 1^{+}, 1^{-}$states can also be formed from $(\ell=\mathrm{L}=1$ or $\ell=\mathrm{L}=2$ ),
8. $\min (\ell=2, L=1$ or $\ell=1, \cdot L=2) ;$ and $(\ell=2, L=2)$; combinations, respectively. The $0^{-}(l=L=1)$ state would give $a \cos ^{2} \theta_{\pi^{+}} \eta^{\text {d }}$ distribution, in disagreement ,with the data. The angular correlations would be relatively complicated in the other cases.
9. The odd intrinaic parity of the $\pi \pi \eta$ oystem requires a $1^{+}$matrix element (proportional to $P_{\pi} \times P_{\eta}$ ) for a vector-meson assignment. N. Cabibbo points out that the resulting $\sin ^{2} \theta_{\pi \eta}$ distribution is essentially independent of any $\pi^{+} r^{-}$. interaction that might be present (such as at 360 MeV . Fig. 3b).
10. The $\Lambda^{0}$ : $5 \pi$ that are unambiguous as to the choice of the eta are shown crosshatched in Fig. 3c. Replotting the "ambiguous" $\Lambda^{0} 5 \pi s o$ that all possible combinations contribute equally to the unit area per event does not appreciably alter the distribution presented in Fig. 3c.
11. The $T=0$ component of the vector-meson octet has $\mathrm{rJ}^{P G}=01^{-0}$. If the minn state with $G=+1$ is this $01^{-\infty}$ meson, it must decay electromagnetically. The $1^{-}$vector assignment is unlikely, because of the $\cos \theta_{\pi \eta}$ distribution (Fig. 3c). An additibnal argument can be made against this assignment. Charge conjugation invariance in the electromagnetic decay requires a $T=1 \pi^{+} \pi^{-} n$ system, so that all neutral decay must be due to the modes $\pi^{0} \gamma, 2 \pi^{0} \gamma$, and $3 \gamma$. However, the decay rate into $\pi^{0} \gamma$ relative to $\pi^{+} \pi^{-} \eta$ is expected to be very large. The $\pi^{0} \gamma$ rate is proportional to a times a large two-body phase-space factor. whereas the $\pi^{+} \pi^{-} \eta$ decay is proportional to $a^{2}$ times a smaller three-body phase-space factor. This is in marked digagreement with the observed all-neutral to $\pi^{+} \pi^{-} \eta$ rates.
12. We use a branching ratio of $\sim 2.5$ for ( 1 decay into all neutrale)/(chargèd decay) and 0.5 for $2 \pi^{0} / \pi^{4} \pi^{-\infty}(2=0)$. For the $\sim 44$ events at low $\Delta^{2}$ at $2.45 \mathrm{BeV} / \mathrm{c}$, this predicte an apportionment of $8,21,4$, and 11 events, respectively. 2s $\pi^{+} \pi^{*} \eta_{\text {chg. }} \cdot \pi^{+} \pi^{-} \eta_{\text {neut. }} \cdot \pi^{0} \pi^{0} \eta_{\text {chg. }}$. and all neutral. We observe 4, 14, §6. ऽ20, respectively, in fair agreement with the predicted apportionment. . In addition, we note that we have not observed any $\pi^{ \pm} \pi^{0} \eta^{0}$ enhancement in $\Sigma^{ \pm}(n \pi)^{\dagger}$ states. This constitutes weak:, evidence against $T=1$.
13. None of the 10 detected $y$ events was consistent with $\Delta y y$ production. The $M M$ in $K^{\prime \prime} p \rightarrow \Lambda+\gamma+M M$ for these events is peaked from 500 to 900 MeV . Thie suggests an average of approximately 4 to 5 gammas for each of the $50 \Lambda^{\circ}+M M$ events. The observed detection efficiency for the gammas is then $10 /(4$ to 5$) \times 50 \approx 4 \%$, a reasonable value for the 72 -in. chamber.
14. We note that evidence for an enhancement in the missing mass opposite the $\Lambda$ 'in the reaction $K^{-}+p \rightarrow \Lambda+$ neutrals near this value has been observed by the Brookhaven bubble chamber group. [See M. Goldberg, M. Gundzik, J. Leitner, S. Lichtman, P. L. 'Connolly, E. L. Hart, K. W. Lai, G. London, G. C. Moneti, R. R. Rau, N. P. Samios, I. O. Skillicorn, and S. S. Yamamoto, Study of $K^{-} p \rightarrow \Lambda\left(K^{0}\right)+$ neutrals at $2.3 \mathrm{BeV} / \mathrm{c}$, Bull. Am. Phys. Soc. 9, 23 (1964).]
15. The existence of a singlet $00^{-+}$meson as well as $1^{+}$mesons has long been conjectured by M. Gell-Mann. J. Schwinger has also proposed a $0^{-} \delta$ meson at a mass of $\sim 1500 \mathrm{MeV}$ [Phys. Fev.r.Letters 12, 237 (1964)].

Fig. 1. (a) Distribution of the effective mass aquared of $\pi^{+} \pi^{+} \pi^{\circ} \pi^{\circ} \pi^{\prime \prime}$ in the reaction $K^{-} p \rightarrow A^{0} \pi^{+} \pi^{+} \pi^{0} \pi^{-}$. The shaded area represente/in which the square of the momentum transfer from the proton to the lambda, $\Delta^{2} p \Lambda^{\prime}$ is less than $0.5 \mathrm{ReV}^{2}$. (b) Distribution of the effective mass squared of all four $\pi^{+} \pi^{0} \pi^{-\quad}$ combinations in the $35 \Lambda^{0} \pi^{+} \pi^{+} \pi^{0} \pi^{-} \pi^{0}$ events in (a) where $0.36 \leqslant \mathrm{M}^{2}(5 \pi) \leqslant 0.98 \mathrm{BeV}^{2}$. The data are at incident momenta of 2.45, 2.63, snd $2.70 \mathrm{BeV} / \mathrm{C}$.

Ng. 2. (a) Distribution of the square of the mass of all neutrals (MM) in the reaction $K " p \rightarrow \Lambda^{0}+$ neutrals (MM). All events consistent with a tit to $\pi+p \rightarrow \Lambda^{0}+K^{0}$ have been subtracted from the plot. (b) Diatribution of the effective mass squared of $\pi^{+}+\pi^{-}+$neutrals (MM) in the reaction $\mathrm{E}^{-} \mathrm{P} \rightarrow \Lambda^{0} \pi^{+} \pi^{+}+$meutrals (MN). The croashatched events have been selected to be consiateat with $\eta^{0}\left(0.28 \leqslant \mathrm{MM}^{2} \leqslant 0.32 \mathrm{BeV}^{2}\right.$ ). (c) Distribution of the square of the effective mass of $\pi^{+} \pi^{0} \pi^{-}$in the reaction $K " p \rightarrow \Lambda^{0} \pi^{+} \pi^{0} \pi^{-}$. (d) Distribution of the square of the effective mass of $\pi^{+} \pi^{-}$and $\pi^{+} \pi^{+} \pi^{-} \pi^{-}$ in the reactions $K^{-} p \rightarrow \Lambda^{0} \pi^{+} \pi^{-}$and $\Lambda^{0} \pi^{+} \pi^{+} \pi^{\oplus} \pi^{-}$. respectively. Shaded events are those with low momentum fransfer from the proton to the lambda $\left(\Delta_{p}^{2} \leqslant 0.5 \mathrm{BeV}^{2}\right)$. All data are at $2.45 \mathrm{BeV} / \mathrm{c}$ incident $\mathrm{K}^{-}$momentum.
Fig. 3. (a) Dalitz plot $\left[M^{2}\left(\pi^{\dagger} \eta^{0}\right)\right.$ versus $M^{2}\left(\pi^{*} \eta^{0}\right]$ and projections for the 61 events la the reartions $K^{-} p \rightarrow \Lambda^{0} 5 \pi$ and $\Lambda^{0} \pi^{+}{ }^{-1}+$ neutrais with $0.86 \leqslant M^{2}(5 \pi)$ or $M^{2}\left(r^{+} r^{\prime \prime}+M M=\eta^{0}\right) \leqslant 0.98$. (b) Distribution of the square of che effective mase of the $\pi^{+} \pi^{-}$outside the $\eta^{0}$ for the 61 events shown in (a). (c) Distribution of the cosine of the angle included between the $\eta^{0}$ and the $\pi^{+} a s$ viewed in the dipion $\left(\pi^{+} \pi^{-}\right)$rest frame. Shaded events are those from the reaction $K^{-} P \rightarrow \Lambda^{0} \pi^{+} \pi^{-}+$neutrals $\left(\eta^{0}\right)$.


Fig. 1


MUB. 2534

Fig. 2


MUB. 2530
Fig. 3

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