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Ligand-Sensitized Lanthanide Nanocrystals: Merging Solid-State Photophysics and Molecular Solution Chemistry

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- 1 Ligand-Sensitized Lanthanide Nanocrystals: Merging Solid-State Photophysics and Molecular
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Abstract

10 To date, the breadth of scientific research that has been devoted to investigating the

11 photochemical and photophysical behavior of the lanthanide elements has generally fallen into

- 12 one of two camps: solution studies of luminescent lanthanide metal-ligand complexes or
- 13 investigations of solid-state nanoparticles, composed primarily of or doped with, lanthanide
- 14 lumiphores. In the latter case, most research of lanthanide nanocolloids has precluded any

15 investigations regarding the use of organic ligands to overcome the difficulties associated with f-

16 f excitation of lanthanides. Instead, most work on condensed-phase lanthanide luminescence has

17 centered on strategies such as d-f charge separation in divalent lanthanides and the sensitization

18 of lanthanide excited states using quantum dots. Current work now aims at bridging the camps of

19 condensed-phase lanthanide photophysics and the solution chemistry of ligand-lanthanide

20 molecular complexes. Recent efforts have partly focused on the fundamental characterization of

21 NaGd_{1-x}Ln_xF₄ nanoparticles featuring surface display of the sensitizer ligand 3,4,3-LI(1,2-

22 HOPO), showing these structures to be capable of converting absorbed ultraviolet light into

23 luminescence from Eu³⁺ and Tb³⁺ ions. These results suggest such use of ligand sensitization as a

24 tool of choice to overcome the constraints of UV solar spectrum/semiconductor band-gap

25 mismatch and low absorption cross-sections in solid-state lanthanide systems.

26

27 Introduction

Our current understanding of the periodic table's f-block marks the result of decades of research 28 29 pursued by scientists of various stripes; from chemists to materials scientists, physicists and biologists. As a consequence, the modern scientific catalog of understood phenomena arising 30 31 either directly or indirectly from lanthanide chemistry and electronic structure is expansive, ranging from multi-photon up- and down-conversion processes 1-16, to ligand-lanthanide 32 chelation thermodynamics in solution 17-20, to lanthanide binding and efflux mechanisms in 33 biological systems^{21–28}. In particular, it has become both widely accepted and abundantly clear 34 that the diversity in the f-orbital electronic structures of the lanthanide elements makes them 35 prime candidates for various possible applications in optoelectronic devices, deep-tissue imaging 36 and luminescence sensing^{10,29–32,23,8,33}. Interest in these applications provides much of the 37 motivation driving both applied and fundamental research in lanthanide photochemistry. 38

39 Molecular Solution Chemistry of Lanthanides

40 With respect to their possible photonic applications, the bulk of photochemical literature of the lanthanides has involved work on either molecular complexes, nanocrystalline forms of 41 42 lanthanides, or glasses where emissive lanthanides occur as either doped or primary constituents^{1–3,9,10,34–39}. In the case of molecular complexes, lanthanide luminescence is 43 44 typically interrogated through the photo-sensitization of lanthanide excited states using organic ligands^{9,23,40}. The employ of this "antenna" effect has the benefit of yielding lanthanide excited 45 46 states with far greater efficiency than generally possible through direct, intra-band pumping of f-47 levels, a partial result of the generally high molar absorptivities of organic/aromatic ligands (~

10³-10⁵ M⁻¹ cm⁻¹). This process stands in stark contrast to the relatively dismal absorption 48 49 coefficients observed for the direct excitation of most lanthanide transitions (~ 1-10 M⁻¹ cm⁻¹), a 50 consequence of the symmetry forbidden and, in some cases, spin-forbidden nature of 51 intraconfigurational f-transitions present in these elements. Sensitization of lanthanides in 52 molecular complexes has also been achieved through the construction of heteronuclear, bi-53 metallic ligands, intricate ligand constructs capable of discriminating between a d-block transition metal and a lanthanide at distinct binding sites^{33,41,42}. There, excitation of a 54 55 transition metal featuring orbitals overlapping with either bridging ligand states or lanthanide-56 centered states allows for energy transfer to the f-element with subsequent emission by the 57 lanthanide. However, rational development and successful solution assembly of these molecular 58 complexes requires an understanding of lanthanide chelation chemistry. The shielding of f-59 electron density by filled 6s and 5p levels in these atoms result in metals that behave 60 predominantly as hard acids, with binding that is primarily ionic, rather than covalent, in 61 character. As hard-soft acid-base (HSAB) theory would predict, these hard, cationic cores are 62 most susceptible to binding by hard donor (base) anionic ligands, explaining the prevalence of oxygen donor atoms in most lanthanide chelators and macrocyclic ligands^{17–19}. Solution 63 64 thermodynamics have played a major role in the quest to develop efficient lanthanide-sensitizing 65 ligands suitable for a variety of applications that utilize f-element properties in the presence of multitudes of additional components such as other metal ions, small organic compounds, or 66 biological macromolecules. Attempts to form stable luminescent complexes range from the 67 68 insertion of simple carboxylate groups onto chromophore-bearing multidentate scaffolds to the 69 use of dye-functionalized dipicolinic acid derivatives or cyclen-based ligands (DOTA,

70 DO3A)^{43,44}, to the synthesis of multidentate ligands incorporating aromatic chelating units such

as the sensitizing 2-hydroxyisophthalamide or hydroxypyridinone⁴⁵. Macrocyclic and

72 octadentate ligands have displayed the highest affinities towards lanthanide ions in solutions with

73 complex stability constant ($\log \beta_{110}$) values nearing or exceeding 25, making them most relevant

⁷⁴ for solution-based applications such as fluorescence-based biological assays⁴⁶.

75 Solid-State Photophysics of Lanthanides

90

76 The direction of research on solid-state glasses and nanocolloidal lanthanide systems has largely been governed by a desire to understand and utilize the photophysical traits of these 77 78 elements. Studies of luminescence behavior in nanoparticulate systems have been especially 79 progressive, with much of the work being geared towards the development of multi-photon conversion for applications in solar energy utilization^{3,4,10}. Extracting spectroscopic 80 81 information from such systems has often demanded the use of laser excitation sources for 82 generating lanthanide luminescence, as their much higher power densities relative to lowerintensity, conventional halogen or arc lamp sources allow for circumventing the issues associated 83 84 with the low luminescence yields of f-f absorption. Higher efficiency lanthanide emission in 85 solid media has also been achieved through lanthanide sensitization by d-block metal ions as co-86 dopants and the recent achievement of lanthanide substitution in ZnS and CdS quantum dots^{35,36,47}. This latter scheme leverages the very high absorption coefficients that have been 87 88 found for metal-chalcogen quantum dots, with the host lattice sensitizing the lanthanide dopant. 89 Lanthanide sensitization by organic ligands typically proceeds through either Forster

energy transfer, Dexter energy transfer, or a combination of the two³³. In the case of Forster

91 transfer, energy is relayed through the dipolar coupling between the triplet excited state of some 92 donor ligand and the 4f orbitals of a lanthanide acceptor. In the limit of pure Dexter transfer, 93 energy flow between a donor and acceptor proceeds through the initial population of the donor excited state, followed by transfer of the excited electron onto the acceptor atom⁴⁸. A second (or 94 95 concerted) exchange then occurs from the acceptor's ground state to the donor's ground state. 96 Relaxation of the electron occupying the acceptor excited state back to ground restores the 97 system to its initial configuration. Forster transfer is primarily a through-space interaction, marked by a strong dependence on spectral overlap between the donor luminescence profile and 98 99 the acceptor absorption spectrum. As a result, this mode of energy transfer generally occurs over 100 several tens of angstroms, with a transfer efficiency governed by:

$$\eta = \frac{1}{1 + \left(\frac{r}{r_0}\right)^6}$$
101 (equation 1),

102 where η is the energy transfer efficiency, r is the donor-acceptor distance, and r_0 , the Forster 103 distance, is defined as the distance yielding a 50% transfer efficiency⁴⁸.

104 The relatively generous distance dependence observed for Forster energy transfer runs 105 contrary to the Dexter mechanism, where energy transfer is restricted to donor/acceptor distances 106 of ~ 5 Å or less. This limitation of the Dexter formalism is a natural consequence of its 107 dependence on electron exchange. Rapid charge transfer between a donor and acceptor is best 108 facilitated via through-bond interactions – the physical overlap of donor and acceptor orbital 109 wavefunctions – which have persistence lengths of only a few angstroms. The result is an energy

5

transfer mechanism displaying a rate-distance dependence that corresponds with thesemiclassical Marcus model of electron transfer:

112
$$k = Aexp(-\beta[r-r_o])$$
 (equation 2).

Here, *k* is the electron transfer rate. β , with units of inverse length, is a medium-dependent property describing the tunneling distance decay of an electron between donor and acceptor. Terms *r* and r_0 retain their earlier definitions⁴⁹. As a consequence of eq. (2), Dexter transfer will have minimal contributions between energy transfers in dilute solutions but may become the dominant channel for energy transfer between a donor and lanthanide acceptor in condensed phases, particularly when donor/acceptor spectral overlap is poor.

119 Nanoparticle research has also focused on modulating lanthanide emission through 120 judicious choice of host lattice; transition probabilities are significantly improved when 121 embedding lanthanides in hosts featuring small phonon energies and low-symmetry (in 122 particular, non-centrosymmetric) metal coordination sites, explaining the prevalence of rare-earth 123 fluoride hosts, particularly the NaGdF₄ and NaYF₄ crystal systems, in luminescent nanoparticle research. ^{50–57}. Such sites are capable of breaking f-orbital symmetry and promoting inter-124 125 orbital mixing, relaxing the conditions for f-f transitioning. In principle, such changes should 126 result in both higher absorption coefficients and more intense luminescence lines, from the 127 metals.

128 Bridging Fields

129 Translating the advances made by researchers in these two camps into projects combining130 the best aspects of these disparate approaches seems an intuitive direction in which to take

research in lanthanide photophysics. However, the amount of literature on the subject of utilizingligand sensitization in solid-state lanthanide structures remains relatively

small^{7,13,15,44,46,52,56,58–68}. The most prominent work produced in this vein include the 133 first ever report of lanthanide nanocrystal sensitization in the IR using tropolonate ligands⁶⁸ and 134 135 a 2012 finding by Zou et. al demonstrating the first example of two-photon up-conversion 136 luminescence in β-NaYF₄:Er:Yb nanocrystals, with initial infrared light absorption through an organic ligand⁵². Despite the promise of this strategy for achieving higher total quantum vields 137 138 for lanthanide lumiphores, the body of scholarship in this area has remained remarkably small since Zhang et. al's first report in 2007⁶⁸. The persistence of this interdisciplinary gap between 139 140 solution chemists and materials scientists/physicists studying lanthanides becomes even more 141 remarkable considering the standing precedent for the role ligand sensitization has played in other areas of solid-state chemistry and physics – with the maturing field of dye-sensitized 142 semiconductor solar cells serving as perhaps the best example⁶⁹. 143

The current research efforts in our group have been adopted while keeping this disparity in mind. Recent work has involved the investigation of luminescence in NaGd_(1-x)Eu_xF₄ and related nanoparticles surface-functionalized with the ligand 3,4,3-LI(1,2-HOPO), an octadentate chelator consisting of a linear spermine backbone derivatized with hydroxypyridine-2-one moieties⁷⁰. The project's conception was partly inspired by projects out of the Abergel group in which the photophysical behavior of f-elements complexed by 3,4,3-LI(1,2-HOPO) in solution were studied^{18,19,71}. This ligand, which displays very high thermodynamic stabilities for 151 lanthanide complexation in aqueous solutions ($\log\beta_{110} = 20.2$ for the [Eu^{III}(3,4,3-LI(1,2-HOPO))]⁻ complex)¹⁹, has been shown to be especially adept at sensitizing Eu³⁺ emission, a feat facilitated 152 by spectral overlap of the ligand's triplet state (centered around 19,000 cm⁻¹) with europium's ⁵D_J 153 states (Figure 1)^{71,72}. The appeal of sensitization in a nanoparticle host featuring low phonon 154 155 energies also promises to help solve the issues associated with non-radiative solvent quenching of luminescent metal ions commonly found in molecular complexes, as the majority of Eu³⁺ ions 156 in nanocrystalline structures reside in the nanoparticle bulk and are thus shielded from 157 interactions with solvent. Following an initial demonstration of europium luminescence 158 159 enhancement in these structures by factors of roughly 5000 fold, we have now focused our 160 efforts on optimizing this system and further characterizing its photophysics.

161 Characterization & Material Optimization

162 Optimizing the Ligand:Nanoparticle Ratio

163 As a follow up to our initial characterizations of rare-earth fluoride nanoparticles modified with 164 3,4,3-LI(1,2-HOPO), we studied the dependence of europium emission on ligand:nanoparticle ratios. Briefly, samples of nanoparticles at concentrations of 0.9 mg mL⁻¹ were incubated with 165 166 varying concentrations of the ligand, spanning from 75 nM to 75 mM. Following overnight 167 incubation at 75 °C and the removal of excess ligand through repeated wash cycles in ethanol, 168 nanoparticle luminescence was measured (Figure 2). These data suggest an incubation ratio of 169 ca. 8 mmol ligand per gram of nanoparticles provides optimal degrees of ligand surface addition. 170 A gradual decay in nanocrystal luminescence can be observed upon increasing incubating 171 concentrations of 3,4,3-LI(1,2-HOPO), despite a steady increase in the amount of surface-bound 172 ligand, as evidenced by nanoparticle absorption spectra (Figure 2, inset). Concentration

173quenching mechanisms between surface-bound ligands may therefore become operative at higher174ligand:nanoparticle ratios. Comparisons of the integrated luminescence spectra at this optimal175incubation and at the previously used value of 75 mM indicate a luminescence enhancement of176roughly 140%. Plots of luminescence intensity as a function of the relative amount of surface-177bound ligand (determined using the 320 nm absorption signal of ligand-modified nanoparticle178suspensions) reveal a moderate, linear correlation ($R^2 = 0.77$) between the two quantities in the179concentration regimes explored (Figure 3).

180 Internal Quantum Efficiency Calculation

As a complement to our previous report of the external quantum yield in these nanocrystals, we have recently determined the approximate internal quantum yield of europium emission in this system. Measurement of this quantity was achieved using the ${}^{5}D_{0} \rightarrow {}^{7}F_{1}$ transition as basis for calculating total luminescence output. This transition in europium has been generally accepted as purely magnetic-dipole in nature, making it insensitive to the ligand field effects that would otherwise modulate intraband f electronic transitions as a function of chemical environment⁷³. In general, the internal quantum yield, Φ_{lm} , is defined as

$$\Phi_{\text{Int}} = \frac{\tau_{obs}}{\tau_{rad}} = \frac{k_{rad}}{k_{obs}}$$
 (equation 3),

189 where τ_{obs} is the observed luminescence lifetime of some transition and τ_{rad} is the luminescence 190 lifetime in the absence of non-radiative quenching processes.

191 A value for τ_{rad} , is found through the relation

188

$$\frac{1}{\tau_{rad}} = An^3 \left(\frac{I_{total}}{I_{MD}} \right)$$
 (equation 4),

where *A* is the spontaneous emission rate of 49.44 s⁻¹ for the ${}^{5}D_{0} \rightarrow {}^{7}F_{1}$ emission line, *n* is the 193 medium's refractive index, I_{total} is the entire integrated luminescence spectrum for Eu³⁺ emission, 194 and I_{MD} is the integrated intensity of europium's magnetic-dipole transition (580 – 600 nm)^{73,74}. 195 Applying the refractive index reported for $NaYF_4/NaGdF_4$ host media (1.5)^{75,76} to eq. (4) yields 196 197 a radiative lifetime of 2.9 ms for this transition. Determination of the observed lifetime, τ_{obs} , is 198 commonly found exclusively through the time-resolved measurement of the intense hyperfine emission (${}^{5}D_{0} \rightarrow {}^{7}F_{2}$) that often dominates europium spectra. However, such an assumption 199 200 remains valid only for cases where the vast majority of photon emission occurs through this 201 radiative decay channel. This approximation breaks down in cases such as the system 202 investigated here, where spectral integration of nanoparticle luminescence shows that relaxation to the ⁷F₂ state represents only about 60% of the total emission from these samples. This point is 203 highlighted by the marked differences between the luminescence profiles of the [Eu^{III}(3,4,3-204 205 LI(1,2-HOPO))]⁻ molecular complex and the modified nanocrystals, where the intensity ratio of 206 the main hypersensitive transition to the other emission lines is far greater in the spectra of the 207 nanocrystals than those of the solution chelate (Figure 4). As a result, it becomes necessary to 208 report k_{obs} , which actually represents the *total* rate of deactivation of the ⁵D₀ excited state (rather 209 than merely its rate of relaxation to any particular ⁷F_J multiplet) as the sum of individual rate constants for ⁵D₀ depopulation to the first, second and fourth stark levels of the ⁷F ground state 210 211 $(k_{J=1,2,4})$. We focus on these particular levels as their summed emissions comprise over 90% of the 212 total observed luminescence in these nanoparticles, and their appreciable signal intensities make 213 acquisition of time-resolved luminescence data comparatively easy relative to the low-intensity ${}^{5}D_{0} \rightarrow {}^{7}F_{0.3.5.6}$ transitions. 214

215 Observed lifetimes for the ${}^{5}D_{0} \rightarrow {}^{7}F_{1,2,4}$ transitions were determined from time-resolved luminescence data acquired under ligand excitation at 317 nm. Emission wavelengths of 590 nm, 216 612 nm and 695 nm were employed for monitoring decay to the ${}^{7}F_{1}$, ${}^{7}F_{2}$ and ${}^{7}F_{4}$ states, 217 respectively. These data reveal all three transitions display bi-exponential decay character. 218 219 Determination of the physical basis for the presence of multiphase decays in the observed 220 europium emission remains a work in progress, but is possibly the result of different chemical 221 environments seen by bulk Eu³⁺ emitters versus solvent-exposed surface ions. As a result of the multi-phase nature of these decays, rates of deactivation to each ⁷F sub-level are expressed as 222 223 weighted averages of the two distinct rates, k_1 and k_2 :

224
$$\langle k_{J} \rangle = c_{1}k_{1}+c_{2}k_{2}$$
 (equation 5),

.1 .

where $\langle k_J \rangle$ is the averaged rate for decay to the J^{th} level of the ⁷D ground-state manifold, and the 225 weighting coefficients, c_n , are derived from the fractional contribution of each decay phase to the 226 227 total signal intensity, as found through fitting of the time-resolved luminescence data. The expression $\langle k_1 \rangle + \langle k_2 \rangle + \langle k_4 \rangle$ yields a value for k_{obs} , which we report as 3105 s⁻¹ ($\tau_{obs} = 0.32$ 228 229 ms) for this system; applying this observed lifetime and the natural lifetime to eq. (1) yields a subsequent approximation for the internal quantum efficiency of 11%. The ratio Φ_{Ext}/Φ_{Int} yields a 230 231 value for the sensitization efficiency. Using the previously reported value of Φ_{Ext} = 0.033 for 232 europium-centered luminescence, we report an estimate for the sensitization efficiency of 0.30. Intersystem Crossing (ISC) 233-

234 The efficiency of triplet formation represents a component of the sensitization process, which we 235 have estimated through measurements of the steady-state luminescence spectrum of 3,4,3-LI(1,2-236 HOPO) bound to undoped, NaGdF₄ nanoparticles at liquid-nitrogen temperatures (Figure 5).

Absence of Eu³⁺ doping guarantees that dissipation of the ligand triplet state will occur only 237 238 through radiative triplet decay to the ligand's ground state or through non-radiative channels, 239 rather than through triplet-donor/metal-acceptor energy transfer, as the first excited state in Gd³⁺ (32,100 cm⁻¹) is too high to be accessed via triplet sensitization by this ligand. Comparison with 240 241 the room temperature spectrum provides a basis for assigning singlet and triplet states. The room 242 temperature data show a broad luminescence peak with skewed Gaussian character. The ligand 243 emission is sharply resolved into two transitions upon cooling to 77 K. The minor, high-energy 244 peak centered around 415 nm, is assigned as the ligand singlet state (Figure 4). A more intense 245 emission occurs around 525 nm and is attributed to ligand phosphorescence, arising from a triplet state serving as a donor of sufficient energy to overlap with the ${}^{5}D_{1,0}$ states (~19,000, 246 17,250 cm⁻¹) localized on the europium ion. Taking the integrated peak ratios of these two ligand 247 248 states seen in the cryogenic spectrum provides an approximate value for the efficiency of 3,4,3-249 LI(1,2-HOPO) triplet state production, for which we find a value of ca. 0.87. This value agrees well with the qualitative observation that ISC in paramagnetic systems containing larger atoms, 250 such as Gd³⁺, are susceptible to the "heavy atom effect," which is known to greatly enhance the 251 efficiency of singlet/triplet intersystem crossing⁷³. 252

253 Future Directions

While our most extensive characterizations of ligand-sensitized nanoparticles have focused on europium-doped constructs, attention has now been turned towards the ultimate goal of moving away from merely downshifted luminescence to develop a two-photon nanocrystal downconverter using the platform established in our europium studies. Materials capable of such non-linear photon production have applications in a number of fields, and are especially well 259 suited for impacting the field of photovoltaics. In the case of downconversion, efficient light absorption at wavelengths where traditional silicon photovoltaics show low photocurrent 260 response, followed by re-emission of multiple lower-energy photons in the infrared regime, holds 261 great potential for designing more efficient solar cells. A great deal of work has been devoted to 262 designing two-photon down-conversion materials using the Tb^{3+}/Yb^{3+} couple^{4,5,77–79}. The 263 general principle here relies on crystal lattices or glasses doped with both Tb³⁺ and Yb³⁺. 264 265 Following the initial excitation to the terbium ion's ${}^{5}D_{4}$ excited state (~ 20,500 cm⁻¹), its depopulation proceeds through cooperative energy transfer to the ${}^{2}F_{5/2}$ level (~ 10,200 cm⁻¹) of 266 267 neighboring ytterbium species. This two-fold energy mismatch between the terbium and ytterbium excited states allows for the possibility of two-photon emission from Yb³⁺ ions, 268 resulting in near-infrared (NIR) luminescence centered around 980 nm. We have been able to 269 270 demonstrate that extending our study to the corresponding terbium-doped, rare-earth lattices does 271 result in terbium luminescence (Figure 5), suggesting that it may indeed be possible to construct 272 a bis-doped material where ligand-sensitized Tb³⁺ ion emission is quenched by energy transfer to 273 adjacent ytterbium co-dopants, with subsequent two-photon, NIR emission. This would represent a vast improvement over the red Eu³⁺ emission we've investigated thus far, as silicon 274 275 photovoltaics (1.1 eV band gap) exhibit peak performance under infrared illumination.

276 Conclusions

We believe the work being conducted in our group represents a good case study of the value in drawing on the lessons of disparate fields and merging them into a coherent research effort. Our confidence in this approach derives largely from the broader trend established in science over the latter half of the last century to the present, with the burgeoning prominence of

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281 interdisciplinary fields yielding scientific insights that may not have been gained otherwise. The 282 relatively young field of bioinorganic chemistry, a modern synthesis of molecular biology and inorganic chemistry, provides strong validation of this truth. Such an outlook informs our current 283 284 research strategy, which we view as the convergence of approaches from the generally distinct 285 realms of solid-state photophysics and solution chemistry. The product of this approach has been 286 our realization of colloidal lanthanide materials that represent significant improvements in light 287 absorption and luminescence enhancement for constructs of this nature. Our expectation is that continued adherence to this scientific philosophy will position us to make lasting contributions to 288 289 the field of spectral conversion and, hopefully, pave the way for the development of ligand 290 sensitized, two-photon down-conversion nanocrystals.

291 Methods

292 Synthesis

293 Europium-doped nanoparticle synthesis and ligand functionalization were performed as

294 previously described^{54,70}. Terbium-doped nanoparticles (1% doped) were synthesized using

similar procedures. In brief, synthesis proceeded through the addition of 1.98 mL

296 Gd(CH₃CO₂)₃•*x*H₂O and 20 µL Tb(CH₃CO₂)₃•*x*H₂O (Sigma-Aldrich) to stirring 1-oleic acid (4

297 mL; Alfa Aesar) and 1-octadecene (6 mL; 90% Sigma-Aldrich) at room temperature. From here,

synthesis proceeded as usual according to the procedures of Wang et al^{54} .

299 Nanoparticle Luminescence Optimization

- 300 Following synthesis, 2.0 mL aliquots of the nanoparticles were precipitated through the addition
- of 2.0 mL ethanol followed by centrifugation was 13,000 rpm for 10 minutes. Afterward, the
- 302 supernatant was decanted and the particles were washed through resuspension in 2.0 mL fresh

303 ethanol via sonication and mechanical mixing by micropipette until a homogenous suspension was formed. This mixture was then centrifuged again for 10 minutes at 13000 rpm, the 304 305 supernatant decanted, and the particles redispersed in 2.0 mL ethanol. The suspension was 306 divided among 10 Eppendorf microcentrifuge tubes as 200 µL aliquots per tube. Ligand 307 modification reactions were set up by addition of 5, 10, 25, 50, 75 or 100 μ L of 75 mM 3,4,3-308 LI(1,2-HOPO) dissolved in pH 6.5 Hepes buffer to each tube. For lower ligand concentrations 309 explored, a 7.5 mM stock solution of 3,4,3-LI(1,2-HOPO) was used for addition to the tubes in 310 volumes of 0.1, 0.5, 1 and 2 µL. Compensating volumes of buffer were applied to each tube to 311 adjust reaction volumes to a total of 300 µL, with final nanoparticle concentrations of approximately 0.9 mg mL⁻¹ reaction⁻¹. The reactions were then heated and mixed on a shaker 312 313 overnight (180 rpm, 60 °C) to promote 3,4,3-LI(1,2-HOPO) binding to the nanoparticle surfaces. 314 Following incubation, nanoparticles were washed three times in 600 μ L ethanol using the 315 procedure described above, before a final resuspension in 600 μ L ethanol. Luminescence 316 intensity as a function of ligand surface density was determined using procedures described 317 below.

318 Steady-State Luminescence

319 Nanoparticle samples were prepped as dilute colloidal mixtures (A_{500} scatter intensity ~ 0.3) in

320 ethanol to ensure stability of the dispersions during luminescence measurements.

321 Steady-state luminescence spectra were acquired on a Jobin Yvon Horiba Fluorolog system.

322 Luminescence spectra of nanoparticles were collected using a 317 nm excitation wavelength

323 sourced from a xenon arc lamp, 1 nm excitation / 3 nm emission slit settings and 1.0 s integration

times at 1 nm resolution. Observation windows of 550-750 nm and 450-650 nm were employed

- 325 for the monitoring of europium and terbium luminescence, respectively. Interfering second-
- 326 harmonic generation originating from the excitation source (~634 nm) was filtered from acquired
- 327 spectra through the placement of a 400 nm long-pass filter between the sample and detector.
- 328 Determination of the triplet state of 3,4,3-LI(1,2-HOPO) (bound to NaGdF₄ control
- 329 nanoparticles) proceeded as described elsewhere⁷⁰.
- 330 Time-Resolved Luminescence
- 331 Europium luminescence lifetimes were acquired by monitoring the emission of the ${}^{5}D_{0} \rightarrow {}^{7}F_{J=1,2,4}$
- transitions at their respective wavelengths of 590, 612, and 695 nm using a Jobin Yvon Horiba
- 333 Fluorolog spectrometer in time-resolved (MCS lifetime) mode. Experiments were conducted
- 334 with the following instrument parameters: 317 nm excitation, 14 nm excitation bandpass; 4nm
- emission bandpass; 10 μ s channel⁻¹ and 3000 channels sweep⁻¹ (30.0 ms observation window).
- 336 Decay times were extracted via multi-exponential fitting in MATLAB (Supporting Information).

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emission. B) From ref. (70): Transmission electron micrograph of hexagonal NaGd_{1-x} $_x$ manoparticles capped with the ligand 3,4,3-LI(1,2-HOPO). C) A simplified Jablonski diagram of europium sensitization by 3,4,3-LI(1,2-HOPO). D) A generic cooperative energy transfer downconversion scheme. E) The molecular structure of the photosensitizer 3,4,3-LI(1,2-HOPO).





(a) Dependence of nanoparticle luminescence intensity (317 nm excitation wavelength) as a function of the concentration of 3,4,3-LI(1,2-HOPO) used in incubations. The slow drop in intensity at much higher ligand concentrations is attributed to concentration quenching between bound 343 moieties at nanocrystalline surfaces. (Inset) Increasing the ligand solution concentration results in an increase in its localization at the nanoparticle surface before plateauing around 10 mM, as shown through monitoring the ligand 320 nm absorption signal (normalized relative to scatter at 500 nm).

(b) Correlation between relative measures of bound ligand, found through 320 nm absorbance of 3,4,3-LI(1,2-HOPO), and peak nanoparticle luminescence at 612 nm, is approximately linear.



Top panel: Normalized luminescence of ligand-modified, Eu-doped nanoparticles and the Eu(III)-3,4,3-LI(1,2-HOPO) solution complex in ethanol.

Bottom Panel: Integrated spectra indicate that the main transition at 612 nm (integration bounds 600-640 nm) constitutes a much smaller portion of total photon emission in the nanoparticles (ca. 61%) compared to the metal-ligand complex (ca. 80%).



Luminescence spectra of terbium-doped NaGdF₄ nanoparticles sensitized by 3,4,3-LI(1,2-HOPO) (black) and the terbium solution complex normalized to the peak (${}^{5}D_{4} \rightarrow {}^{7}F_{5}$) transition. Nanoparticle emission marks the superposition of terbium luminescence and a broad Raman signature previously ascribed to residual oleate ligands on the nanoparticle surface.



Adapted from (70). Luminescence spectra of Gd(III)-3,4,3-LI(1,2-HOPO) in ethanol at room temperature (black dots) and in a 77 K glass (violet diamonds). The room temperature spectrum represents the coalescence of singlet and triplet states into a skewed Gaussian dominated by singlet character. The 77 K spectrum allows for assignment of the peak at 525 nm to the ligand triplet excited state, with a minor contribution from the ligand singlet visible at 415 nm. Integration of these peaks in the cooled spectrum allows for estimation of the efficiency for triplet production via intersystem crossing.

