

Promoter effects of alkali metal cations on the electrochemical reduction of carbon dioxide

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Abstract

The electrochemical reduction of CO₂ is known to be influenced by the identity of the alkali metal cation in the electrolyte; however, a satisfactory explanation for this phenomenon has not been developed. Here we present the results of experimental and theoretical studies aimed at elucidating the effects of electrolyte cation size on the intrinsic activity and selectivity of metal catalysts for the reduction of CO₂. Experiments were conducted under conditions where the influence of electrolyte polarization is minimal in order to show that cation size affects the intrinsic rates of formation of certain reaction products, most notably for HCOO⁻, C₂H₄, and C₂H₅OH over Cu(100)- and Cu(111)-oriented thin films, and for CO and HCOO⁻ over polycrystalline Ag and Sn. Interpretation of the findings for CO₂ reduction was informed by studies of the reduction of glyoxal and CO, key intermediates along the reaction pathway to final products. Density functional theory calculations show that the alkali metal cations influence the distribution of products formed as a consequence of electrostatic interactions between solvated cations present at the outer Helmholtz plane and adsorbed species having large dipole moments. The observed trends in activity with cation size are attributed to an increase in the concentration of cations at the outer Helmholtz plane with increasing cation size.

Keywords: Electrochemical reduction, CO₂, copper, silver

Introduction

The electrochemical reduction of CO₂ offers a means for storing electrical energy produced by renewable but intermittent resources, such as wind and solar radiation.^{1,2} Hydrocarbons and alcohols, rather than carbon monoxide and formic acid, are the preferred products of CO₂ reduction because of their high energy density. Despite extensive efforts aimed at identifying electrocatalysts that can produce these products, copper (Cu) remains the only material capable of doing so with significant yields.³⁻⁵ Nevertheless, Cu requires a high overpotential (~ -1 V versus RHE) and produces a broad spectrum of products.^{6,7} Clearly, the discovery of novel means for reducing CO₂ to desirable products with higher efficiency and selectivity is needed. To do so requires understanding the fundamental processes occurring at the electrode surface. In addition to metal-adsorbate interactions, a full description of electrochemical reactions must include the effects of solvation and interfacial electric fields. A number of studies have shown that one way of influencing the product distribution is through changes in the electrolyte cation.⁸⁻¹³

The role of cations is particularly interesting because a number of studies have shown that both the activity and selectivity of Ag, Hg, and Cu for the CO₂ reduction reaction (CO₂RR) are influenced significantly by the size of the alkali metal cation in the electrolyte.⁸⁻¹³ Similar effects have also been observed for the electrochemical reduction of oxygen, the oxidation of hydrogen, and the oxidation of low molecular weight alcohols.¹⁴⁻¹⁸ For Ag and Hg, which are selective for CO₂ reduction to carbon monoxide (CO) and formate anions (HCOO⁻), respectively, increasing the alkali metal cation size increases the rates of formation of these products.^{8, 10, 11} For Cu, increasing alkali metal size leads to higher selectivities to C₂ products, e.g., ethylene and ethanol.^{9, 12, 13} Explanations for these phenomena have been attributed to differences in local pH, or differences in kinetic overpotentials due to the electrochemical potential in the outer Helmholtz plane being affected by cation size.^{8, 10, 11, 13}

One of the challenges to finding a consistent interpretation for the effects of cation size on the intrinsic kinetics of CO₂ reduction is that many studies are carried out under conditions where mass transport also contributes to the observed activity and product distribution. In recent study, we reported the effects of cation size on the rate of the CO₂RR under strong mass-transport limitations.¹³ We showed that the increase in pH and the decrease in CO₂ concentration near the cathode, caused by inadequate mass transfer through the hydrodynamic boundary layer in front of the cathode, could be offset by the hydrolysis of solvated cations. We showed that the increased buffering effect of solvated cations with cation size could be captured theoretically and suggested that this effect is responsible for the observed changes in the overall current density and the distribution of product current densities with cation size.

In this work, we present a combined experimental and theoretical study of the effects of electrolyte cation size on the intrinsic activity and selectivity of metal catalysts for the reduction of CO₂. In contrast to earlier studies, experiments were conducted at sufficiently low applied voltages to avoid the effects of concentration polarization of the electrolyte, so that the influence of cation size on the intrinsic activity and selectivity of each catalyst could be observed. Most of our work was conducted using epitaxially grown Cu(100) or Cu(111) thin films; a more limited study was carried out using polycrystalline Ag and Sn films. Density functional theory (DFT) calculations were used to investigate how cations and cation size affect the stabilization of reaction intermediates involved in the reduction of CO₂. These calculations demonstrate that adsorbed species with large dipole moments oriented away from the solvent are stabilized by electrostatic interactions with solvated cations, while trends in activity with cation size arise from differences in ion concentration at the Helmholtz plane.

Methods

Electrode Preparation

Single-side polished Si(100) or (110) wafers (Virginia semiconductor, 1-10 Ω cm) were diced into ~ 4 cm² sized pieces that were then used as electrode substrates. Prior to Cu deposition, the native oxide was removed from the Si substrates by submerging them in a 10% HF solution for 5 min. Immediately after HF etching, the Si pieces were transferred to a vacuum chamber for sputter deposition of Cu in an AJA ATC Orion-5 sputtering system. The base pressure of the sputtering system prior to deposition was $\sim 2 \times 10^{-7}$ torr. The flow rate of the sputtering gas (Ar) was 25 sccm and the sputtering pressure was adjusted to 2×10^{-3} torr by controlling the speed at which the chamber was pumped, using a variable butterfly valve. Cu (99.999% Kurt Lesker) was deposited at a rate of 1 $\text{\AA}/\text{s}$, as determined by a calibrated quartz crystal monitor, at ambient temperature. The total film thickness deposited was 100 nm. Polished, polycrystalline foils of Ag and Sn (99.99% Alfa Aesar) were also used as electrodes without further preparation.

Electrode Characterization

The structure of the Cu thin films was characterized by X-ray diffraction. The orientation and epitaxial quality of the films were determined using symmetric θ - 2θ scans, in plane ϕ scans, ω scans or rocking curves, and pole figures. XRD patterns were taken with a PANalytical X'Pert diffractometer, which uses a Cu K α ($\lambda = 1.54056$ \AA) X-ray source. Symmetric θ - 2θ scans were collected on samples fixed onto a flat glass slide in locked-coupled mode with a goniometer resolution of 0.001° . Measured diffraction patterns were compared to known standards taken from the International Center for Diffraction Data (ICDD) PDF4 database (card #71-4610 for Cu).

Electrochemical Measurements

All electrochemical experiments were conducted in a gas-tight electrochemical cell machined from polyether ether ketone (PEEK).¹⁹ The cell was cleaned with 20 wt. % nitric acid and oxidized in UV-generated ozone for 15 min prior to the initiation of an experiment. The working and counter electrodes were parallel and separated by an anion-conducting membrane (Selemion AMV AGC Inc.). A gas dispersion frit was incorporated into the cathode chamber to provide vigorous electrolyte mixing. The exposed geometric surface area of each electrode was 1 cm² and the electrolyte volume of each electrode chamber was 1.8 mL. The counter electrode was a Pt foil (99.9% Sigma Aldrich) that was flame annealed

prior to each experiment. The working electrode potential was referenced against an Ag/AgCl electrode (Innovative Instruments Inc.) that was calibrated against a homemade standard hydrogen electrode (SHE). 0.05 M M_2CO_3 (M referring to an alkali metal cation) solutions were prepared by mixing ultra-pure salts (Sigma Aldrich 99.995%) and 18.2 M Ω DI, and were used as the electrolyte without further purification. The cathode chamber was sparged with CO₂ (99.999% Praxair) at a rate of 5 sccm for 20 min prior to and throughout the duration of all electrocatalytic measurements. Upon saturation with CO₂ the pH of the electrolyte was 6.8, which was maintained throughout the duration of electrolysis.

Electrochemical measurements were performed using a Biologic VSP-300 potentiostat. All electrochemical data were recorded versus the reference electrode and converted to the RHE scale. Potentiostatic electrochemical impedance spectroscopy (PEIS) was used to determine the uncompensated resistance (R_u) of the electrochemical cell by applying voltage waveforms about the open circuit potential with an amplitude of 20 mV and frequencies ranging from 50 Hz to 500 kHz. The potentiostat compensated for 85% of R_u *in situ* and the last 15% was post-corrected to arrive at accurate potentials. The electrocatalytic activity was assessed by conducting chronoamperometry at each fixed applied potential for 70 min.

Product Analysis

The effluent from the electrochemical cell was passed through the sampling loop (250 μ L) of an Agilent 7890B gas chromatograph equipped with a pulsed-discharge helium ionization detector (PDHID). He (99.9999% Praxair) was used as the carrier gas. The effluent of the electrochemical cell was sampled every 14 min. The gaseous products were separated using a Hayesep-Q capillary column (Agilent) connected in series with a packed ShinCarbon ST column (Restek Co.). The column oven was maintained at 50 °C for 1 min followed by a temperature ramp at 30 °C/min to 250 °C, which was maintained for the duration of the analysis. The signal response of the PDHID to each gaseous product was calibrated by analyzing a series of NIST-traceable standard gas mixtures (Air Gas).

The electrolyte from both electrode chambers was collected after electrolysis and analyzed using a Thermo Scientific UltiMate 3000 liquid chromatograph equipped with a refractive index detector (RID). The electrolyte aliquots were stored in a refrigerated autosampler until analyzed in order to minimize the

evaporation of volatile products. The liquid-phase products contained in a 10 μL sample were separated using a series of two Aminex HPX 87-H columns (Bio-Rad) and a 1 mM sulfuric acid eluent (99.999% Sigma Aldrich). The column oven was maintained at 60 $^{\circ}\text{C}$ for the duration of the analysis. The signal response of the RID to each liquid-phase product was calibrated by analyzing standard solutions of each product at a concentration of 1, 10, and 50 mM.

Theoretical Calculations

Plane-wave density functional theory (DFT) employing periodic boundary conditions was used for all calculations in this study. The Quantum ESPRESSO²⁰ code and the Atomic Simulation Environment (ASE)²¹ were used, along with ultrasoft pseudopotentials and the Bayesian error estimation exchange-correlation functional with van der Waals interactions (BEEF-vdW).²² This functional has been optimized for both chemisorption energies and long-range van der Waals interactions. For all calculations, a plane-wave cutoff of 500 eV and a density cutoff of 5000 eV were used, based on convergence tests from previous studies.²³ A Monkhorst-Pack²⁴ k -point grid of $(4\times 4\times 1)$ was used for (3×3) cells while a $(2\times 2\times 1)$ grid was used for (6×6) cells. All systems were modeled using a periodic Cu(111) slab, two ice-like water bilayers, the solvated ion and at least 12 \AA of vacuum perpendicular to the surface. The field effect on various CO_2 adsorbates without ions or solvating waters was also investigated by applying a sawtooth potential in each simulation cell, for moderate field strengths where the vacuum energy remains above the Fermi level.²⁵ As an example, Figure S3 in the Supporting Information (SI) shows the average potential energy along the z -direction for CH_3 on Cu (111), with an applied sawtooth potential.

Constrained minima hopping²⁶ (CMH) was used for the global optimization of water structures surrounding each solvated ion. CMH enforces molecular identity by constraining the atomic distances within each molecular entity through Hookean constraints. The atoms are thermalized in a Maxwell-Boltzmann distribution at an initial temperature of 1000 K, then evolved through molecular dynamics and relaxed to locate new local minima. To save computational cost, CMH was performed on two layers of

Cu. The Cu atoms were fixed, while all others were allowed to relax. The convergence criterion for the CMH was satisfied when three successive lowest energy minima were found within 0.05 eV of each other. Thereafter, the system was relaxed until the forces on all atoms were less than 0.05 eV/Å. Finally, the third layer of Cu was added back into the system to perform a single point calculation and obtain a more accurate value for the energy.

Results

Cation effects on Cu(100) and Cu(111) oriented films

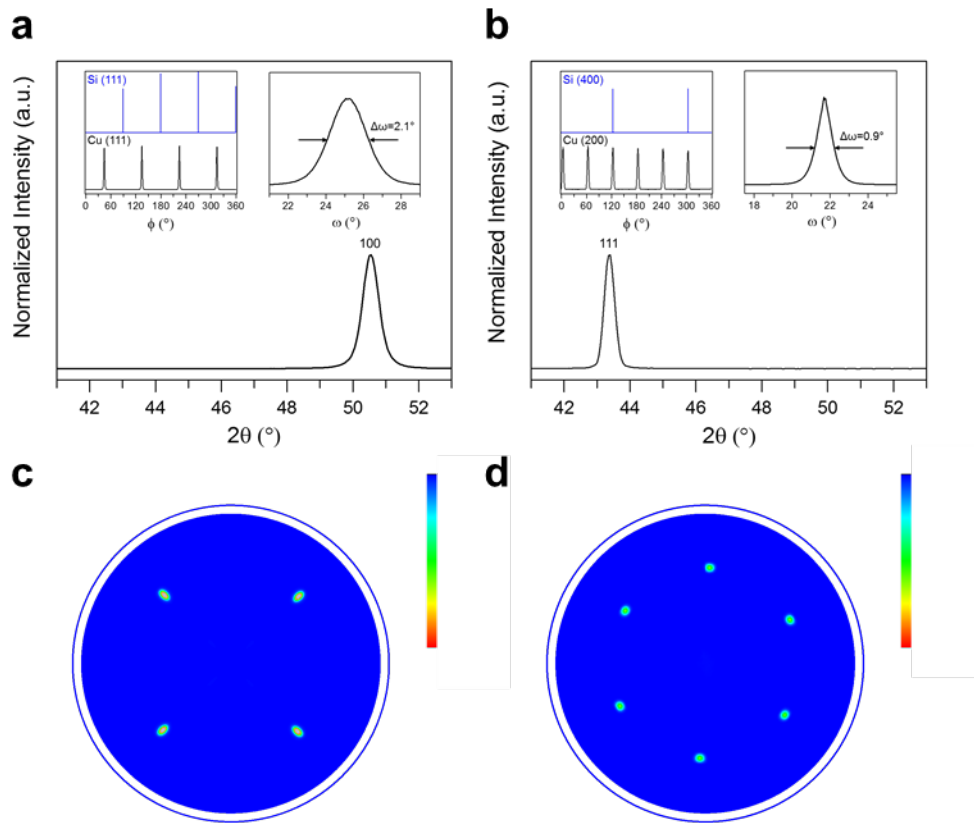


Figure 1: Structural characterization of oriented Cu thin films. a) Symmetric θ - 2θ , in plane ϕ , and ω rocking curve scans for Cu deposited on Si(100). b) Symmetric θ - 2θ , in plane ϕ , and ω rocking curve scans for Cu deposited on Si(110). c) Cu(111) pole figure for Cu deposited on Si(100). d) Cu(200) pole figure for Cu deposited on Si(110).

Epitaxially grown Cu thin films were used as electrocatalysts in order to have well-defined catalytic surfaces.²⁷ These films were prepared by radio-frequency sputtering of Cu onto silicon (Si) single crystal substrates, taking advantage of the epitaxial relationship between Cu and Si substrates of different orientations.²⁷⁻²⁸ Most of our studies were conducted on the Cu(100) surface because it has been shown that polycrystalline Cu consists mostly of (100) oriented crystallites and the Cu(100) surface has a higher selectivity to more desirable products such as C₂₊ hydrocarbons and oxygenates (e.g., ethylene and ethanol).²⁹⁻³² Symmetric θ -2 θ XRD patterns of 100 nm thick Cu films deposited on Si(100) and Si(110) substrates are shown in Figure 1. The observation of a single Cu(200) diffraction peak for Cu deposited on Si(100) indicates that the film is oriented with the (100) crystallographic direction normal to the film plane. Similarly, for Cu deposited on Si(110) only the diffraction peak associated with Cu(111) is observed.²⁷⁻²⁸ While symmetric XRD scans establish the out-of-plane texture relationships, both out-of-plane and in-plane texture analyses are necessary to determine whether the Cu thin films grow epitaxially on Si, as a film with strong fiber texture would show a similar diffraction pattern. To this end, in-plane texture analysis was conducted using X-ray pole figures. Pole figures represent the texture of a material and show the distribution of particular crystallographic directions in a stereographic projection. The Cu(111) X-ray pole figure for Cu(100) on Si(100) exhibits discrete Bragg reflections, indicating cube-on-cube epitaxial growth of Cu on the Si(100) substrate. Four-fold symmetry is observed for the Cu(111) Bragg reflections with an azimuthal angle of 90° apart, which is expected for a Cu (100) single-crystal. For Cu (111) on Si (110), six-fold symmetry is observed for the Cu (200) Bragg reflections, indicating both strong out-of-plane and in-plane texture and thus epitaxial growth on the Si(110) substrate (Figure 1a). Six diffraction spots in the pole figure are seen instead of three because there are two discrete sets of crystallites from twinning with an azimuthal angle of 60° apart. These results demonstrate that Cu thin films can be grown epitaxially with both (100) and (111) orientations.

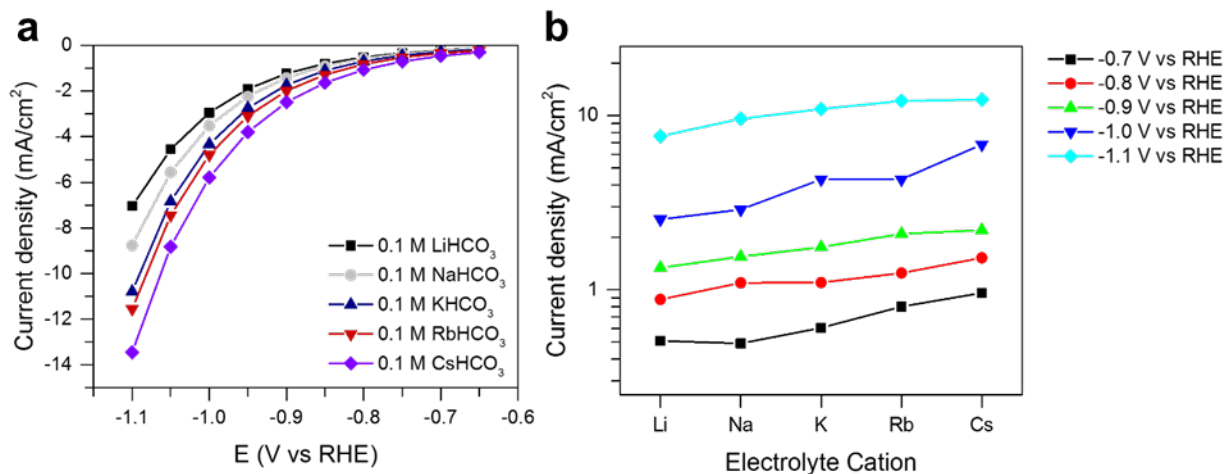


Figure 2: Effect of alkali metal cations on the total activity over Cu(100). a) Linear sweep voltammograms for CO₂ reduction on Cu(100) in CO₂ saturated 0.1 M bicarbonate electrolytes containing different metal cations. b) Average current densities obtained during bulk electrolysis as a function of metal cation at different potentials.

The steady state activity and selectivity of the Cu(100) surface was investigated in CO₂-saturated bicarbonate electrolytes containing different alkali metal cations, by using potentiometric electrolysis at potentials from -0.7 to -1.1 V vs the reversible hydrogen electrode (RHE). Figure 2 shows the total current density as a function of electrolyte cation at various potentials. From both steady state potentiostatic and potential sweep measurements, it is clear that the total current density increases with cation size, and this trend is observed at all applied potentials. To identify the contributions to the increase in the total current density, we examined the effects of cation size on the product distribution. As large differences in total current density exist for electrolytes containing different cations, partial currents provide a better representation of trends in product formation rates than faradaic efficiencies. Complete information regarding faradaic efficiencies can be found in the supplementary information (Figure S1, Table S1).

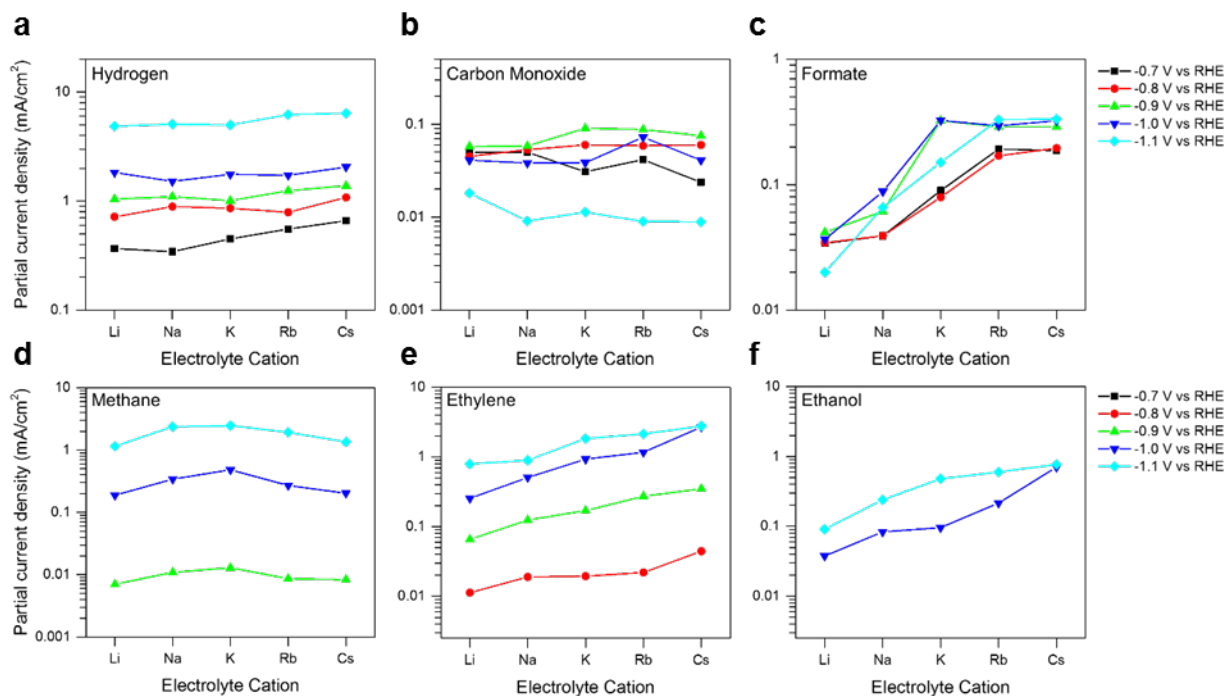


Figure 3: Cation effect on the rates of formation of major products of CO₂ reduction over Cu(100). Partial current densities for each of the major products as a function of the electrolyte metal cation on Cu(100). Data are presented at potential between -0.7 to -1.1 V vs RHE.

Figure 3 shows the partial current density for each of the major products as a function of electrolyte cation. Formation rates of minor liquid products can be found in the supplementary information (Table S1). A key finding is that the production rates of hydrogen and carbon monoxide are relatively unaffected by increasing cation size, whereas the rates of formate, ethylene, and ethanol formation increase monotonically with increasing cation size. The rate of methane production reaches a maximum for Na⁺- and K⁺- based electrolytes and is lower in all other electrolytes. The trends seen in Fig. 3 are consistent with previous amperometric measurements.⁹ It is notable that these trends are independent of the applied overpotential.

To probe the effects of mass transport limitations, the CO₂ flow rate into the electrochemical cell was varied. As the CO₂ flow rate is increased, the hydrodynamic boundary layer thickness will decrease, resulting in a larger maximum diffusive flux of CO₂ to the electrode surface. At low reaction rates (i.e.,

low current densities), for which diffusion of CO₂ is sufficiently rapid to prevent depletion at the cathode surface, increasing the flowrate of CO₂ should have no effect on catalyst activity. Whereas under mass transport limited conditions, increased reactant availability will affect the product distribution measured.^{13, 32} Figure 4 shows that doubling the CO₂ flow rate from 5 to 10 sccm has no impact on the measured selectivity or activity of Cu (100) at -0.9 V vs RHE and -1.0 V vs RHE. However, at -1.1 V vs RHE, doubling the CO₂ flowrate decreases the faradaic efficiency to H₂ and CH₄ and increases the faradaic efficiency for C₂₊ products. As shown previously, these trends are due to the increased rate of mass transfer at higher gas flowrates, which improves the supply of CO₂ to the cathode.³² Therefore, the effects of cation size on the partial current densities observed in Figure 3 at potentials more positive than ~ -1.0 V vs RHE are ascribable to the influence of cation size on the intrinsic kinetics of Cu(100), particularly for the formation of formate, ethylene, and ethanol, and to a lesser degree, methane.

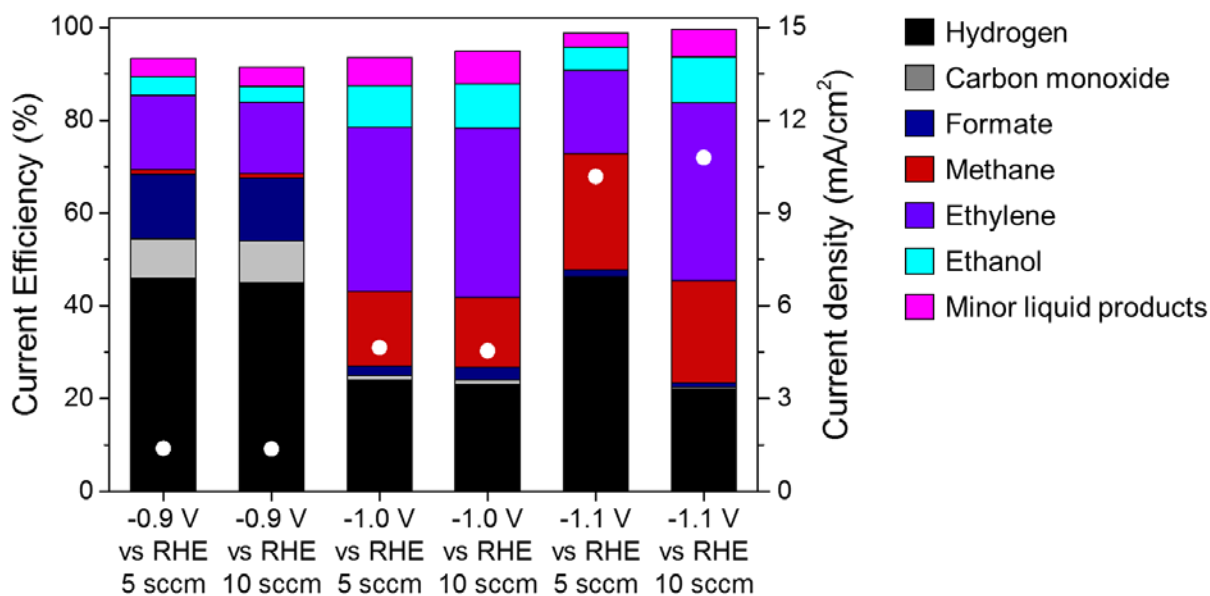


Figure 4: CO₂ flow rate studies for probing external mass transport limitations: Current efficiency (bars) and total current density (dots) for CO₂ reduction over Cu(100) at different flow rates of CO₂.

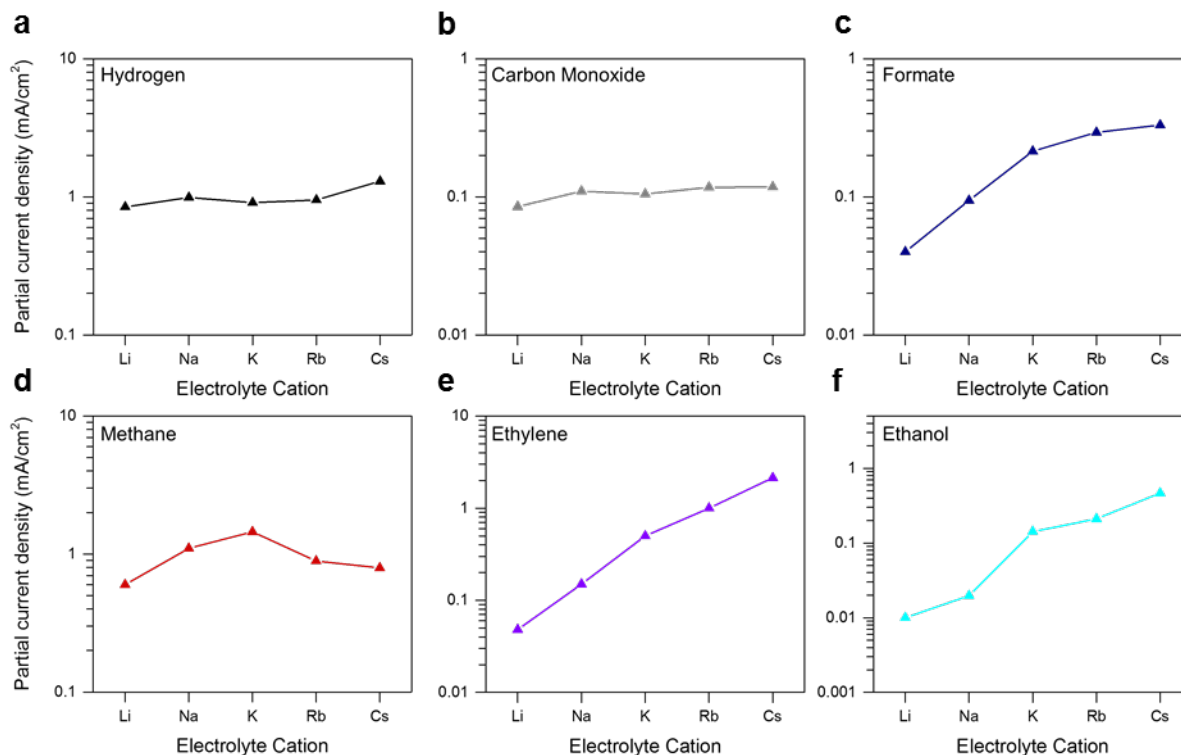


Figure 5: Cation effect on the rates of formation of major products of CO₂ reduction over Cu(111). Partial current densities for each of the major products as a function of the electrolyte metal cation on Cu(111). Data are presented for the potential of -1.0 V vs RHE.

The effects of cation size were also investigated for CO₂ reduction over a Cu(111) surface and the results are presented in Fig. 5. The trends with increasing cation size for the close-packed Cu(111) surface are consistent with those observed for the Cu(100) surface. Partial currents to HCOO⁻, C₂H₄, and C₂H₅OH increase with increasing cation size while rates of formation of H₂, CO, and CH₄ are less affected by cation size. However, the effect of cation size is more pronounced on the close-packed surface, resulting in high selectivity (~40% FE) to C₂H₄ in CsHCO₃ even on a Cu(111) surface, which is known to be more selective for producing CH₄ than C₂H₄ in K⁺-containing electrolytes.^{29, 31, 33-36}

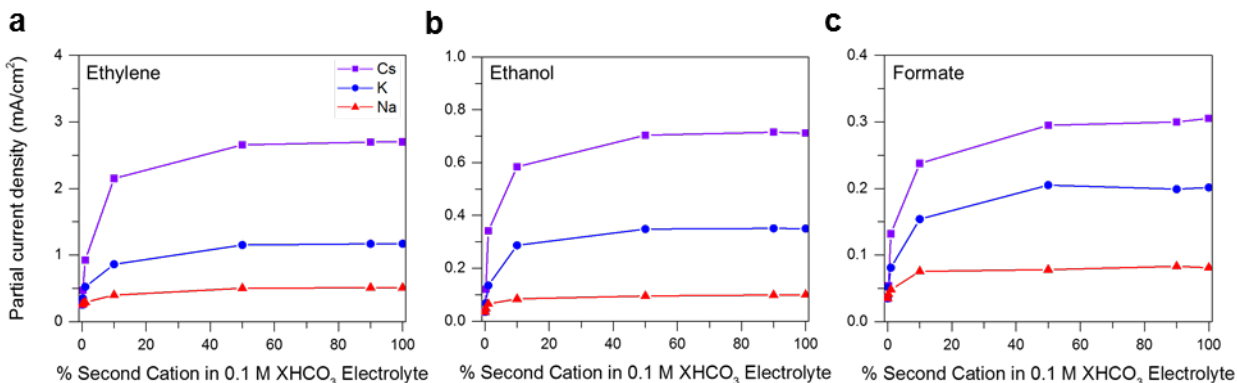


Figure 6: Relative effects of alkali metal cations. Partial currents for ethylene, ethanol, and formate formation as a function of the electrolyte composition on Cu(100). Data are presented for a potential of -1.0 V vs RHE. A mixture of LiHCO_3 and XHCO_3 electrolyte is used with X being Na, K, Cs (second cation). A fixed total salinity of 0.1 M was used in all experiments.

Relative effects of alkali metal cations

To further investigate the effects of the alkali metal cations, electrolysis was performed with mixtures of small (Li^+) and larger (Na^+ , K^+ , Cs^+) cations in the electrolyte, maintaining the total salinity at 0.1 M. The results are shown in Figure 6. It is evident that the composition of the electrolyte has a noticeable effect on the partial currents for HCOO^- , C_2H_4 , and $\text{C}_2\text{H}_5\text{OH}$. No significant effect on the partial currents of H_2 , CO , and CH_4 , was observed. For the latter three products, the effect of the larger cation dominates that of the Li^+ cation, and even at relatively low percentages of the larger cation. Conversely, an electrolyte mixture that is dilute in the smaller cation results in identical selectivity to a solution that does not contain the small cation. These results suggest that the concentration of larger cations near the cathode surface is larger and/or that the larger cations have a much more significant promotional effect on the production of C_2H_4 , $\text{C}_2\text{H}_5\text{OH}$, and HCOO^- .

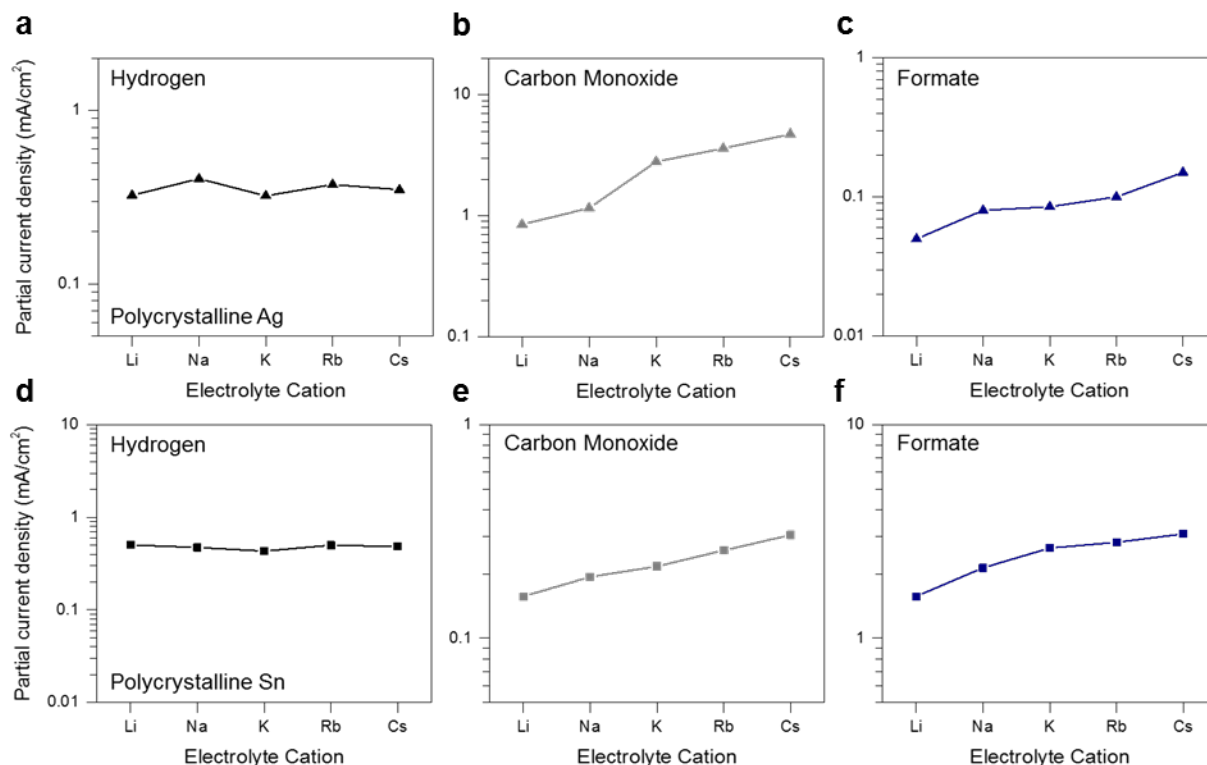


Figure 7: Cation effect on two electron products. Partial currents for each of the major products as a function of the electrolyte metal cation on polycrystalline Ag and Sn. Data are presented for a potential of -1.0 V vs RHE.

Cation effect on polycrystalline Ag and Sn

As carbon monoxide is a reaction intermediate which undergoes further reduction on Cu, the production rate of CO itself is less straightforward to interpret than that of the other major products.³⁷ Therefore, to investigate the effect of cation size on CO production, and to provide further insights into the trends observed for Cu(100) and Cu(111), experiments were conducted on polycrystalline Ag and Sn foils—metals that selectively produce CO and HCOO⁻, respectively.^{3, 38} As seen in Fig. 7, the current densities for CO and HCOO⁻ production on both metals increases with cation size. This suggests that increasing cation size facilitates activation of CO₂. Here again, we see that as with Cu(100) and Cu(111), the rate of H₂ evolution is unaffected by cation size. These trends are also consistent with previous literature reports.^{8-11, 13, 39}

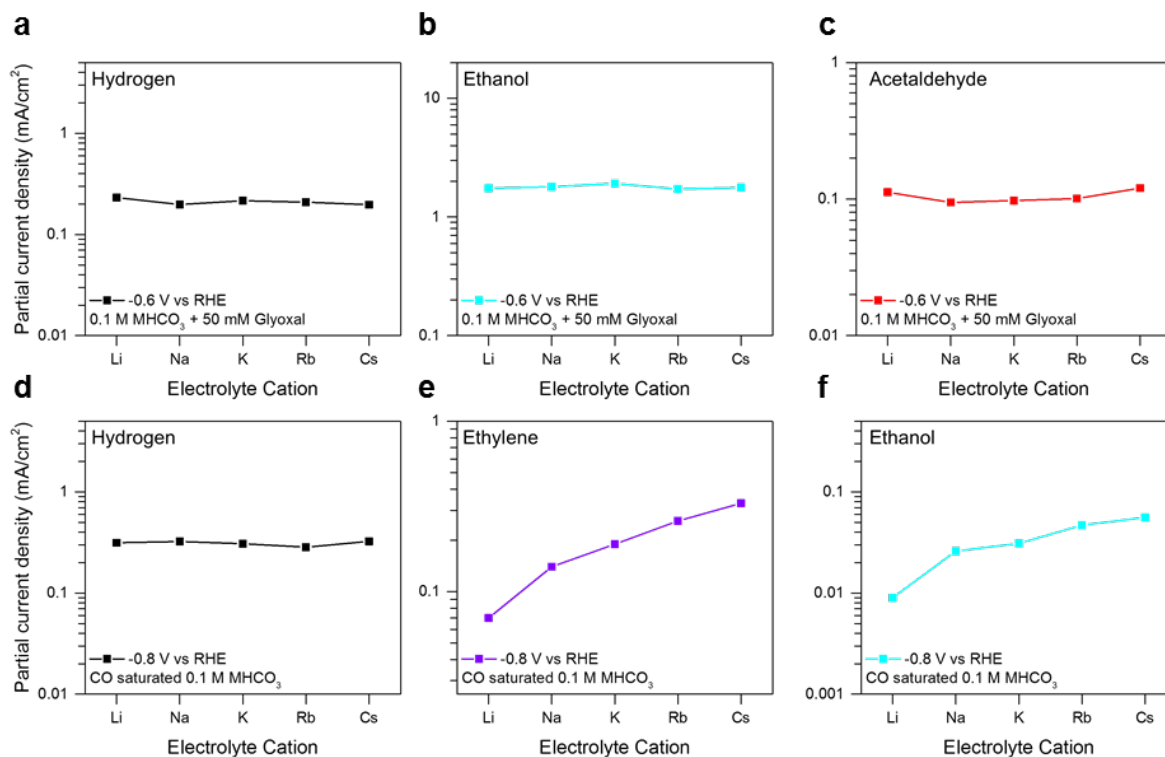


Figure 8: Cation effects of the reduction of reaction intermediates. a-c) Partial currents for hydrogen, ethanol, and acetaldehyde formation from the reduction of glyoxal on Cu(100). Data is presented at a potential of -0.6 V vs RHE. d-f) Partial currents for hydrogen, ethylene, and ethanol formation from the reduction of CO on Cu(100). Data is presented at a potential of -0.8 V vs RHE.

Cation effects on elementary reaction steps: reduction of intermediates

As shown above, the size of the electrolyte cation influences the formation of C₂ products. Recent theoretical studies suggest that C-C bonds are formed via the reaction of two molecules of adsorbed CO or by the reaction of adsorbed CO with a formyl group, HCO.⁴⁰ A recent study has shown that the latter process dominates at applied potentials near -1 V and leads to the formation of adsorbed HC(O)-C(O) species, which then undergoes hydrogenation to form C₂H₄ or CH₃CHO and C₂H₅OH.⁴⁰ Experimental studies suggest that the precursor to CH₃CHO and C₂H₅OH is glyoxal (OHC-CHO).⁴¹ Therefore, to investigate which steps in the formation of C₂ products are affected by the electrolyte cation size, we carried out experiments in electrolyte containing glyoxal. As shown in Figure 8, cation size has no effect on the partial currents of CH₃CHO and C₂H₅OH formation, which in this case require no C-C bond formation as glyoxal is the starting reagent. Consistent with previous observations, no C₂H₄ was observed

during the hydrogenation of glyoxal.⁴¹ These findings strongly suggest that the size of the electrolyte cation affects the rate of C-C bond formation, the rate increasing with increasing cation size, but has a negligible effect on the hydrogenation of glyoxal to CH₃CHO and C₂H₅OH.

As larger cations facilitate CO₂ activation, evidenced by the increase in rates to CO and HCOO⁻ (see Figure 3), experiments were carried out in electrolyte solutions saturated with CO but free of dissolved CO₂, in order to assess whether an increase in the surface coverage of adsorbed CO relative to adsorbed H favors C-C coupling. Consistent with the results for CO₂ reduction over Cu(100), no change in methane activity with cation size was observed from CO reduction at -1 V vs RHE. Figure 8b clearly shows that the rates of C₂H₄ and C₂H₅OH formation increase with increasing cation size. We note that the magnitude of this effect is comparable to that observed during the reduction of CO₂, which shows that the effect of cation size on the formation of multicarbon products cannot be ascribed solely to more facile CO₂ activation.

Discussion

Cation promoter effects

The data presented in this study demonstrate that the size of the alkali metal cations in the electrolyte affects the intrinsic rates of formation of certain reaction products. On Cu (100) and Cu (111), these effects are most notable for HCOO⁻, C₂H₄, and C₂H₅OH; over polycrystalline Ag and Sn, the effects of cation size are observed for the formation of CO and HCOO⁻. It is also shown that larger cations have a stronger effect than smaller cations, and that only a small fraction of the larger cation is needed in the electrolyte to observe its influence. Trends in the partial current densities derived from the studies of the reduction of key reaction intermediates (OHCCHO and CO) suggest that the elementary steps promoted by the presence of metal cations are the activation of CO₂ and the initial C-C coupling, whether through

CO dimerization or CO-CHO coupling, and not the hydrogenation of C₂ intermediates leading to ethylene and ethanol.

We propose that the energetics for the formation of HCOO⁻ and C₂ products are influenced by the presence of hydrated alkali ions located at the edge of the Helmholtz plane, and that the observed effects of cation size are attributable to differences in the cation concentration in the outer Helmholtz plane. This explanation is supported by DFT calculations performed using minimum energy structures of solvated cations at the interface. Since both Cu(100) and Cu(111) surfaces show consistent trends in activity with respect to alkali cation size, the close-packed Cu(111) surface was chosen for the DFT simulations, which has also been frequently used in previous theoretical CO₂R studies⁴²⁻⁴⁵ and where the water structure is more well-defined.⁴⁶

Because the potentials applied during CO₂ reduction are much more negative than the potential of zero charge (PZC) of the low-index facets of Cu $\sim -0.7V_{\text{SHE}}$ ⁴⁷, solvated cations should accumulate near the surface of the electrode during reaction.^{40, 48} Given the very negative reduction potentials of alkali ions, they do not chemisorb under CO₂ reduction conditions. There is also evidence from studies of oxygen reduction that the cations remain partially solvated at the interface.⁴⁹ The presence of the cation at the metal surface gives rise to high electric fields of $\sim -1 \text{ V}/\text{\AA}$ in the vicinity of the ion.^{25, 50} In contrast to classical continuum models of the electrochemical interface, the fields are highly localized when the structure of the cation and solvating waters is modeled explicitly. Refs. ^{25, 50} show that such fields extend roughly 5 Å from the center of the solvated cations and outside of this region decay to zero.

The strength of the adsorbate-field interaction can be probed qualitatively by the application of a uniform field in vacuum. This approach circumvents the challenge of finding a global minimum in the water and solvated cation structure in the presence of various adsorbates on the surface, which introduces many more degrees of freedom than standard atomistic computations in heterogeneous catalysis. It has also been shown previously that the effect of alkali promoters in heterogeneous catalysis can be modeled

approximately by the effect of a uniform field.⁵¹ The interaction energy between an adsorbate and a uniform electric field at the interface is given by⁵²

$$\Delta E = \mu\epsilon - \frac{1}{2}\alpha\epsilon^2 + \dots \quad , \quad (1)$$

where μ and α are the dipole moment and the polarizability (respectively) of the adsorbate, and ϵ represents the electric field. Adsorbate dipoles oriented in the opposite direction of the field will be stabilized, and vice versa.

Figure 9 shows the field stabilization for various adsorbates that have been proposed to be involved in the electrochemical reduction of CO₂ on Cu(111).^{43, 49, 52} The change in adsorption free energy for each species was determined by applying a uniform field oriented perpendicular to the surface in vacuum.⁴³ The corresponding values of μ and α , determined by fitting Eq. (1) the calculated data, are given in Table 1. We note that H* does not have a significant value of μ and is not affected by the electric field, which is consistent with previous experimental studies that have shown that hydrogen adsorption is unaffected by the identity of the cation in solution.¹⁴⁻¹⁶ Since the energy for hydrogen adsorption is a descriptor for the hydrogen evolution activity on metals, the rate for hydrogen production should also be relatively insensitive to the identity of the cation, consistent with the experimental data presented here.⁵³ Recent calculations of the electrochemical barriers for CO₂ reduction using an explicit solvent model have shown that reduction of *CO to form *CHO is the limiting step for CH₄ formation.⁴³ Since the value of μ for *CHO is zero, the cation-induced field would stabilize *CO and therefore increase the energy barrier for hydrogenation to *HCO, the rate-limiting step on the path to CH₄ formation. Therefore, CH₄ formation should occur in regions of the catalyst that are not strongly influenced by the cations located at the outer Helmholtz plane, and correspondingly the partial current density for CH₄ should not be strongly affected by cation identity, which is consistent with the trend in the experimental data shown above.

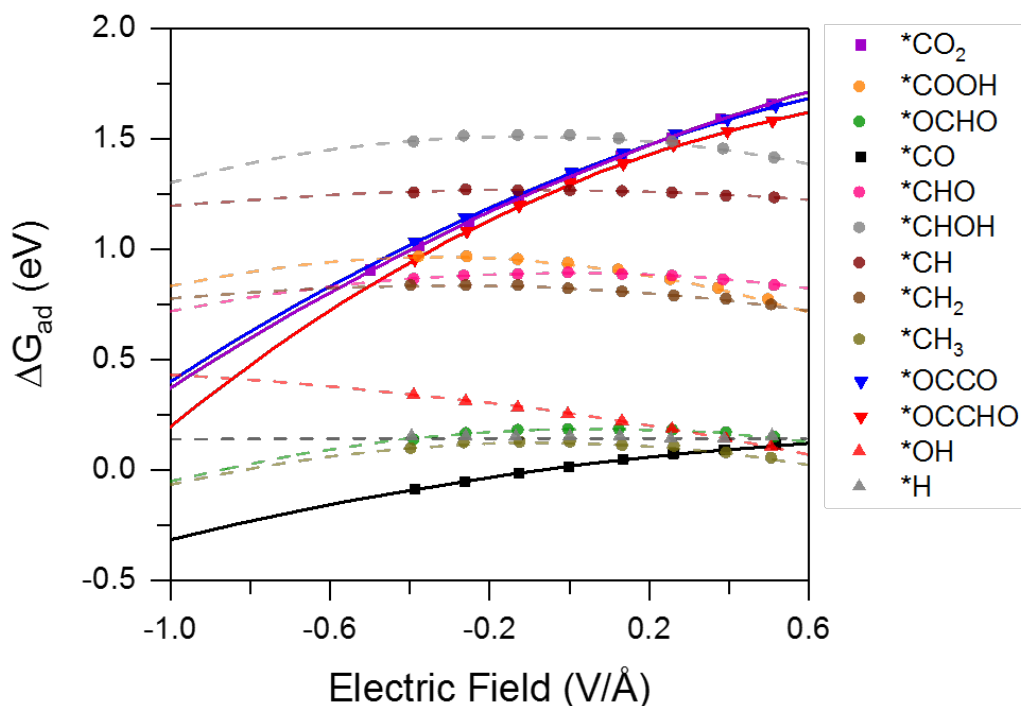


Figure 9: Field effects on various CO₂R intermediates on Cu(111). The energy of each adsorbate is plotted as a function of field strength, which is obtained by applying a uniform electric field oriented perpendicular to the slab. The curves are fit to Equation (1) and the corresponding dipole moment μ and polarizability α of the adsorbates are given in Table 1. Solid curves highlight the adsorbates that are most strongly affected by an electric field. Adsorbates with positive dipole moments μ (oriented in the direction opposite to the electric field) are stabilized, while adsorbates with negative μ (aligned with the electric field) are destabilized.

Table 1 Dipole moments and polarizabilities of adsorbates on Cu(111), determined with Eq. (1)

Adsorbate	μ (eÅ)	α (eÅ ² /V)
*CO ₂	0.76	0.40
*COOH	-0.19	0.60
*OCHO	0.04	0.42
*CO	0.23	0.22
*CHO	0.00	0.36
*OCCO	0.66	0.54
*OCCHO	0.75	0.74
*CHOH	-0.05	0.52
*CH	-0.02	0.18
*CH ₂	-0.08	0.26
*CH ₃	-0.03	0.44
*OH	-0.26	0.14
*H	0.00	0.00

For adsorbates containing C=O bonds oriented perpendicular to the surface—in particular, *CO₂, *OCCO, and *OCCHO—the field stabilization should be large, ~ 1 eV for fields of ~ -1 V/Å in the vicinity of the cations. As shown in Table 1, these species have the largest dipole moments. Previous theoretical calculations have shown that on weakly binding metals such as Ag, the field-induced stabilization of *CO₂ is crucial for the activation of CO₂ to form CO.²⁵ The barrier for the hydrogenation of *CO₂ to *COOH is generally very small,⁵⁴ so stabilization of *CO₂ should increase the coverage of *COOH. The *COOH species can undergo further hydrogenation to form HCOOH or *CO. The chemical hydrogenation barriers for *COOH to form HCOOH are surmountable at 0.69 eV and 0.75 eV on Cu(111) and Cu(211) respectively, and 0.46 eV on Ag(211),²³ consistent with increased HCOOH production observed on all of the surfaces investigated. For weakly binding metals such as Ag and Sn, the field effect would promote CO_(g) production, whereas for moderate to strongly binding metals such as Cu and Pt, *CO produced in the vicinity of a cation would also be field stabilized, thereby decreasing its desorption rate and increasing its *CO coverage there.

On Cu surfaces, formation of the initial intermediates to C₂ products are also field-assisted. The greater stabilization of the *OCCO relative to 2*CO in the presence of a cation is due to the differences in the dipole moments of the product and reactant species,³⁶ which has the effect of lowering the barrier to C-C bond formation. In the absence of the cation-created electrostatic field and effects of solvation by water, the barrier for this process is not surmountable at room temperature.⁵⁵ At higher overpotentials, where *CHO can be produced, *CO-*CHO coupling would also be a feasible route towards C₂ products, and the corresponding barrier should be similarly field-stabilized due to the large dipole on *OCCHO, as shown in Figure 9.⁴⁰ Since the formation of glyoxal via addition of a hydrogen atom to *OCCHO is the first step along the path to forming ethylene and ethanol, the observation (see Figure 8) that the selectivity of products derived from glyoxal is independent of cation size suggests that the rate-determining step in formation of C₂ products occurs at the point of C-C bond formation and prior to formation of glyoxal.⁵⁴

The net effect of increased CO production, CO coverage and C-C coupling on Cu on field-stabilized sites results in a net non-monotonic trend in the partial current density for CO with cation size.

While the influence of the cation-induced field at the catalyst surface can explain the role of cations in promoting elementary steps in the reduction of CO₂ that involve strong dipole moments, it does not explain the observed trend with cation size. We therefore explored the possibilities that this effect is due to disparities in the magnitude of the electrostatic field induced by cations of different size and/or an increase in the concentration of solvated cations at the outer Helmholtz plane as a function of size.

Constrained minima hopping²⁶ (CMH) was used to obtain the optimized structure of water surrounding each solvated ion on Cu(111). The structures of the solvated ions are shown in Figure S5 and the distances of the solvated cation from the Cu(111) surface are given in Table S3. Figure 10 shows that the calculated fields in the vicinity of adsorbed CO (*CO) do not show a systematic variation with cation size. Moreover, any differences in the binding energy of intermediates induced by cations would give rise to exponential differences in current density, which we do not observe.

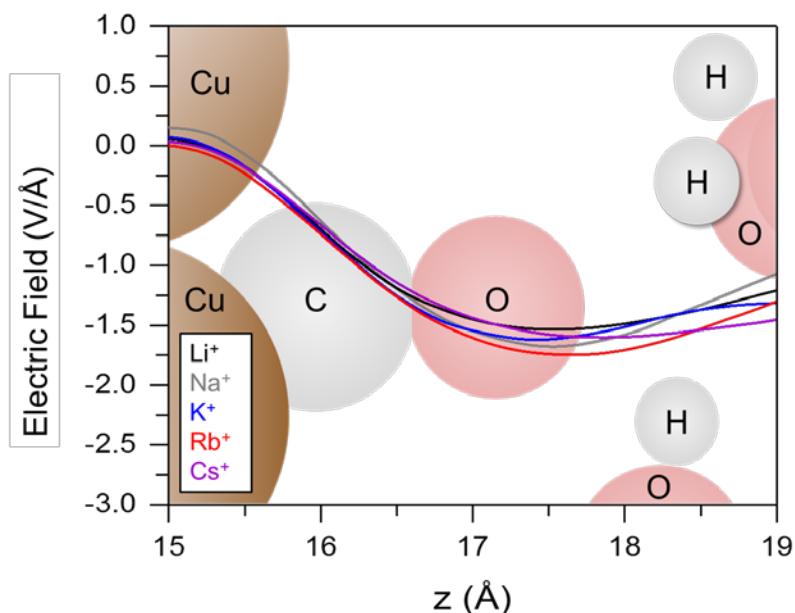


Figure 10. Electric field distribution near the center of the adsorbate plotted as a function of the z -coordinate of the simulation cell for the 2 *CO initial state. The geometry of an example structure is overlaid to clearly illustrate the position of the adsorbate. The cell is viewed from the side and rotated

90°. The minimum z value is taken at the Cu surface and the z -position of the solvated cations is roughly at 20 Å and not included in the overlay.

We hypothesize, therefore, that the ion specificity arises from an increase in the concentration of cations at the outer Helmholtz plane with increasing cation size. To determine the driving force for each cation toward the Helmholtz plane from bulk solution, we used these globally optimized solvation structures for each cation (Table S3) and calculated the energy required to move a solvated cation from the bulk of the electrolyte to a fixed distance of 6.25 Å above the Cu(111) surface the (see the SI for details). Figure 11 shows this energy change with the cathode potential relative to that for Li⁺, $\Delta\Delta E$. The potential U vs SHE is determined through:

$$U = \frac{\Phi - \Phi_{SHE}}{e} \quad (2)$$

where Φ is the work function of the interface and Φ_{SHE} is that corresponding to the standard hydrogen electrode, determined experimentally to be ~4.4eV.⁵⁶ The value of $\Delta\Delta E$ becomes more negative with increasing cation size, its magnitude increasing in the order Li⁺ < Na⁺ < K⁺ < Rb⁺ < Cs⁺. Therefore, the driving force for each cation to be at the Helmholtz plane relative to the bulk electrolyte increases with cation size. This trend suggests that the concentration of cations at the outer Helmholtz plane should increase with increasing cation size, leading to increased product current densities. Consistent with this reasoning, we observed (see Figure 6) that at a constant cation concentration, replacement of even a small fraction of Li⁺ by Cs⁺ results in a shift in the product distribution from one characteristic for Li⁺ to one characteristic for Cs⁺. This supports the hypothesis that there is a stronger driving force for Cs⁺ to be at the Helmholtz plane compared to Li⁺. Without this stronger driving force on Cs⁺ to be in the Helmholtz plane, the trends in Figure 6 would be linear.

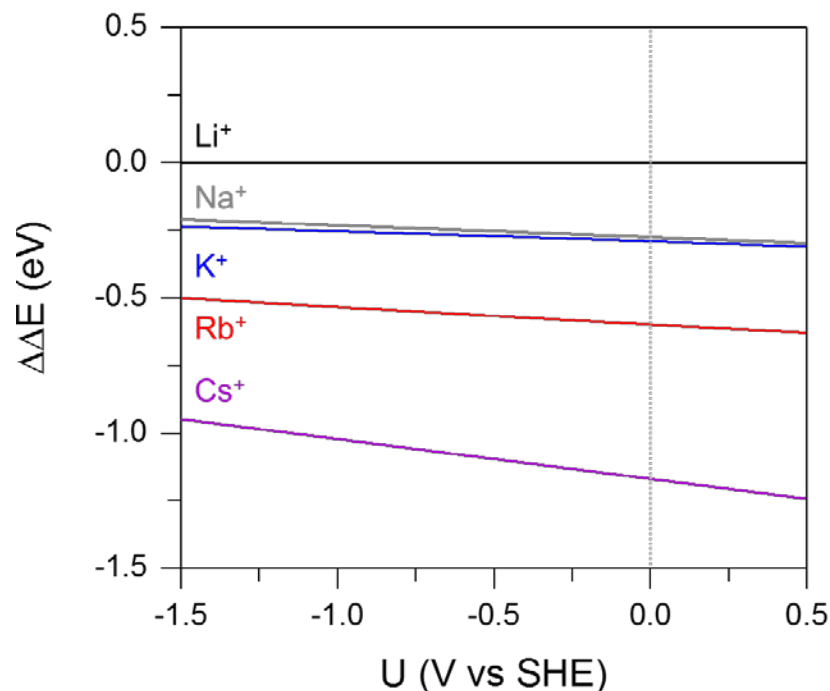


Figure 11: The change in energy for bringing a solvated cation from bulk electrolyte to the outer Helmholtz plane at the Cu (111) facet for different cations shown as a function of the standard hydrogen electrode (SHE) potential. All energies are referenced to that for the Li^+ cation. The dependence of these energies on cathode potential results from the partial electron transfer to the ion upon migration from the bulk to the Helmholtz plane, and is reflected in Eq. 5. This energy becomes more negative with increasing ion size, suggesting a greater concentration of larger cations in the outer Helmholtz plane and giving rise to increased product current densities. The vertical dotted line indicates the potential for the standard hydrogen electrode.

Conclusions

The present study investigated the effects of electrolyte cation size on the activity and selectivity of metal catalysts used for the electrochemical reduction of CO_2 . The experimental work was conducted under conditions where the influence of electrolyte concentration polarization was negligible in order to avoid the additional influence of changes in pH and the consequent effects on the concentration of CO_2 near the cathode surface. For polycrystalline Ag and Sn, increasing the size of the alkali metal cation in the electrolyte increases the partial current densities for HCOO^- and CO, but has no effect on the partial current density for H_2 . Similar experiments conducted using Cu(100) and Cu(111) oriented films

deposited on crystalline Si(100) and(111) surfaces show an increase in the current densities for HCOO^- , C_2H_4 , and $\text{C}_2\text{H}_5\text{OH}$ but little to no effect on the partial current densities for H_2 , CO , and CH_4 as the size of the electrolyte cation is increased.

The observed effects are interpreted on the basis of DFT calculations. Hydrated alkali metal cations in the outer Helmholtz plane create a dipole field of the order $1 \text{ V}/\text{\AA}$, which can stabilize the adsorption of surface intermediates having significant dipole moments (e.g., $^*\text{CO}_2$, $^*\text{CO}$, $^*\text{OCCO}$). This field stabilization decreases the energy for the reduction of $^*\text{CO}_2$ adsorption, the precursor to two-electron products, and C-C coupling to form $^*\text{OCCO}$ or $^*\text{OCCHO}$, the precursor to C_2H_4 and $\text{C}_2\text{H}_5\text{OH}$. Calculations of the stability of the solvated cation in the Helmholtz plane indicate that larger hydrated cations are more energetically favored at the outer Helmholtz plane than smaller ones, which suggest a larger coverage of cations as cation size increases. This result is consistent with the experimental observation that a small percentage of larger cations has a large effect on product distribution.

These findings are consistent with the trends in partial current for C_1 products observed on polycrystalline Ag and Sn and on Cu(100) and Cu(111) surfaces, and for C_2 products observed on Cu(100) and Cu(111) surfaces. Also consistent with the conclusions drawn from the theoretical analysis is that cation size affects the partial currents for C_2H_4 and $\text{C}_2\text{H}_5\text{OH}$ produced during the electrochemical reduction of CO over Cu(100) but has no influence on formation of these products upon reduction of glyoxal. The absence of cation size effect on the partial current for H_2 produced on Ag, Sn, Cu(100), and Cu(111) is attributed to the absence of a dipole for $^*\text{H}$. Similarly, no effect of cation size on the formation of CH_4 on Cu(100) and Cu(111) surfaces is expected since the field of a cation would tend to stabilize $^*\text{CO}$ more than $^*\text{CHO}$, thereby increasing the energy for $^*\text{CO}$ hydrogenation to $^*\text{CHO}$, the rate-limiting step leading to CH_4 .

Finally, we note that differences in the effects of electrolyte cation size on the formation of CO on Ag and Sn vs. Cu(100) and Cu(111) are attributable to the following considerations. For Ag and Sn, CO is a final product, and therefore only influenced by its formation by hydrogenation of $^*\text{COOH}$ which

should have an increased coverage from the field-stabilization of *CO₂. In contrast, for Cu(100) and Cu(111), *CO is an intermediate in the formation of CH₄. Since the dipole electrostatic field created by hydrated cations accelerates both the formation and consumption of *CO, there is no observable effect of cation size on the partial current for CO on Cu(100) and Cu(111).

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References

1. Lewis, N. S.; Nocera, D. G. *Proc. Natl. Acad. Sci.* **2006**, 103, 15729-15735.
2. Chu, S.; Majumdar, A. *Nature* **2012**, 488, 294-303.
3. Hori, Y., Electrochemical CO₂ Reduction on Metal Electrodes. In *Modern Aspects of Electrochemistry*, Vayenas, C. G.; White, R. E.; Gamboa-Aldeco, M. E., Eds. Springer New York: New York, NY, 2008; 89-189.
4. Kuhl, K. P.; Hatsukade, T.; Cave, E. R.; Abram, D. N.; Kibsgaard, J.; Jaramillo, T. F. *J. Am. Chem. Soc.* **2014**, 136, 14107-14113.
5. Hori, Y.; Wakebe, H.; Tsukamoto, T.; Koga, O. *Electrochim. Acta* **1994**, 39, 1833-1839.
6. Hori, Y.; Murata, A.; Takahashi, R. *J. Chem. Soc. Faraday Trans.* **1989**, 85, 2309-2326.
7. Kuhl, K. P.; Cave, E. R.; Abram, D. N.; Jaramillo, T. F. *Energy Environ. Sci.* **2012**, 5, 7050-7059.
8. Frumkin, A. N. *Trans. Faraday Soc.* **1959**, 55, 156-167.
9. Akira, M.; Yoshio, H. *Bull. Chem. Soc. Jpn* **1991**, 64, 123-127.
10. Thorson, M. R.; Siil, K. I.; Kenis, P. J. A. *J. Electrochem. Soc.* **2013**, 160, 69-74.
11. Paik, W.; Andersen, T. N.; Eyring, H. *Electrochim. Acta* **1969**, 14, 1217-1232.
12. Kaneco, S.; Iiba, K.; Katsumata, H.; Suzuki, T.; Ohta, K. *J. Solid State Electrochem.* **2007**, 11, 490-495.

13. Singh, M. R.; Kwon, Y.; Lum, Y.; Ager, J. W.; Bell, A. T. *J. Am. Chem. Soc.* **2016**, 138, 13006-13012.
14. Subbaraman, R.; Tripkovic, D.; Strmcnik, D.; Chang, K.-C.; Uchimura, M.; Paulikas, A. P.; Stamenkovic, V.; Markovic, N. M. *Science* **2011**, 334, 1256-1260.
15. Stoffelsma, C.; Rodriguez, P.; Garcia, G.; Garcia-Araez, N.; Strmcnik, D.; Marković, N. M.; Koper, M. T. M. *J. Am. Chem. Soc.* **2010**, 132, 16127-16133.
16. Strmcnik, D.; Kodama K.; van der Vliet, D.; Greeley, J.; Stamenkovic, V. R.; Marković, N. M. *Nat. Chem.* **2009**, 1, 466-472.
17. Suntivich, J.; Perry, E. E.; Gasteiger, H. A.; Shao-Horn, Y. *Electrocatalysis* **2013**, 4, 49-55.
18. Angelucci, C. A.; Varela, H.; Tremiliosi-Filho, G.; Gomes, J. F. *Electrochem. Comm.* **2013**, 33, 10-13.
19. Lobaccaro, P.; Singh, M. R.; Clark, E. L.; Kwon, Y.; Bell, A. T.; Ager, J. W. *Phys. Chem. Chem. Phys.* **2016**, 18, 26777-26785.
20. Giannozzi, P.; Baroni, S.; Bonini, N.; Calandra, M.; Car, R.; Cavazzoni, C.; Ceresoli, D.; Chiarotti, G.; Cococcioni, M.; Dabo, I.; Dal Corso, A.; de Gironcoli, S.; Fabris, S.; Fratesi, G.; Gebauer, R.; Gerstmann, U.; Gougoussis, C.; Kokalj, A.; Lazzeri, M.; Martin-Samos, L.; Marzari, N.; Mauri, F.; Mazzarello, R.; Paolini, S.; Pasquarello, A.; Paulatto, L.; Sbraccia, C.; Scandolo, S.; Sclauzero, G.; Seitsonen, A.; Smogunov, A.; Umari, P.; Wentzcovitch, R. *J. Phys. Condens. Mat.* **2009**, 21.
21. Bahn, S. R.; Jacobsen, K. W. *Comput. Sci. Eng.* **2002**, 4, 56-66.
22. Wellendorff, J.; Lundgaard, K. T.; Møgelhøj, A.; Petzold, V.; Landis, D. D.; Nørskov, J. K.; Bligaard, T.; Jacobsen, K. W. *Phys. Rev. B* **2012**, 85, 235149.
23. Yoo, J. S.; Abild-Pedersen, F.; Nørskov, J. K.; Studt, F. *ACS Catal.* **2014**, 4, 1226-1233.
24. Monkhorst, H.; Pack, J. *Phys. Rev. B* **1976**, 13, 5188-5192.
25. Chen, L. D.; Urushihara, M.; Chan, K.; Nørskov, J. *ACS Catal.* **2016**, 6, 7133-7139.
26. Peterson, A. A. *Top. Catal.* **2013**, 57, 40-53.
27. Hahn, C.; Hatsukade, T.; Kim, Y.-G.; Vailionis, A.; Baricuatro, J. H.; Higgins, D. C.; Nitopi, S. A.; Soriaga, M. P.; Jaramillo, T. F. *Proc. Natl. Acad. Sci.* **2017**, 114, 5918-5923.
28. Jiang, H.; Klemmer, T. J.; Barnard, J. A.; Payzant, E. A. *J. Vac. Sci. Tech. A* **1998**, 16, 3376-3383.
29. Hori, Y.; Takahashi, I.; Koga, O.; Hoshi, N. *J. Phys. Chem. B* **2002**, 106, 15-17.
30. Schouten, K. J. P.; Qin, Z.; Gallent, E. P.; Koper, M. T. M. *J. Am. Chem. Soc.* **2012**, 134, 9864-9867.
31. Roberts, F. S.; Kuhl, K. P.; Nilsson, A. *ChemCatChem* **2016**, 8, 1119-1124.
32. Kim, Y.-G.; Baricuatro, J. H.; Javier, A.; Gregoire, J. M.; Soriaga, M. P. *Langmuir* **2014**, 30, 15053-15056.
33. Durand, W. J.; Peterson, A. A.; Studt, F.; Abild-Pedersen, F.; Nørskov, J. K. *Surf. Sci.* **2011**, 605, 1354-1359.
34. Schouten, K. J. P.; Pérez Gallent, E.; Koper, M. T. M. *ACS Catal.* **2013**, 3, 1292-1295.
35. Montoya, J. H.; Shi, C.; Chan, K.; Nørskov, J. K. *J. Phys. Chem. Lett.* **2015**, 6, 2032-2037.
36. Sandberg, R. B.; Montoya, J. H.; Chan, K.; Nørskov, J. K. *Surf. Sci.* **2016**, 654, 56-62.
37. Hori, Y.; Murata, A.; Yoshinami, Y. *J. Chem. Soc. Faraday Trans.* **1991**, 87, 125-128.
38. Hatsukade, T.; Kuhl, K. P.; Cave, E. R.; Abram, D. N.; Jaramillo, T. F. *Phys. Chem. Chem. Phys.* **2014**, 16, 13814-13819.
39. Kyriacou, G. Z.; Anagnostopoulos, A. K. *J. Appl. Electrochem.* **1993**, 23, 483-486.
40. Goodpaster, J. D.; Bell, A. T.; Head-Gordon, M. *J. Phys. Chem. Lett.* **2016**, 7, 1471-1477.
41. Schouten, K. J. P.; Kwon, Y.; van der Ham, C. J. M.; Qin, Z.; Koper, M. T. M. *Chem. Sci.* **2011**, 2, 1902-1909.
42. Akhade, S. A.; McCrum, I. T.; Janik, M. J. *J. Electrochem. Soc.* **2016**, 163, 477-484.
43. Liu, X.; Xiao, J.; Peng, H.; Hong, X.; Chan, K.; Nørskov, J. K. *Nat. Comm.* **2017**, 8, 15438-15444.

44. Shi, C.; Hansen, H. A.; Lausche, A. C.; Nørskov, J. K. *Phys. Chem. Chem. Phys.* **2014**, 16, 4720-4727.
45. Xiao, H.; Cheng, T.; Goddard, W. A.; Sundararaman, R. *J. Am. Chem. Soc.* **2016**, 138, 483-486.
46. Ogasawara, H.; Brena, B.; Nordlund, D.; Nyberg, M.; Pelmenchikov, A.; Pettersson, L. G. M.; Nilsson, A. *Phys. Rev. Lett.* **2002**, 89, 276102.
47. Lukomska, A.; Sobkowski, J. *J. Electroanal. Chem.* **2004**, 567, 95-102.
48. Cheng, T.; Xiao, H.; Goddard, W. A. *J. Am. Chem. Soc.* **2016**, 138, 13802-13805.
49. Strmcnik, D.; van der Vliet, D. F.; Chang, K. C.; Komanicky, V.; Kodama, K.; You, H.; Stamenkovic, V. R.; Marković, N. M. *J. Phys. Chem. Lett.* **2011**, 2, 2733-2736.
50. Montoya, J. H.; Shi, C.; Chan, K.; Nørskov, J. K. *J. Phys. Chem. Lett.* **2015**, 6, 2032-2037.
51. Mortensen, J. J.; Hammer, B.; Nørskov, J. K. *Phys. Rev. Lett.* **1998**, 80, 4333-4336.
52. Nørskov, J. K. N.; Studt, F.; Abild-Pedersen, F.; Bligaard, T., *Fundamental Concepts in Heterogeneous Catalysis*. Wiley, Hoboken, New Jersey: 2014.
53. Nørskov, J. K.; Bligaard, T.; Logadottir, A.; Kitchin, J. R.; Chen, J. G.; Pandelov, S.; Stimming, U. *J. Electrochem. Soc.* **2005**, 152, 23-26.
54. Shi, C.; O'Grady, C. P.; Peterson, A. A.; Hansen, H. A.; Nørskov, J. K. *Phys. Chem. Chem. Phys.* **2013**, 15, 7114.
55. Montoya, J. H.; Peterson, A. A.; Nørskov, J. K. *ChemCatChem* **2013**, 5, 737-742.
56. Trasatti, S. *Pure Appl. Chem.* **1986**, 58.

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