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Publication Date 1973-04-01

Submitted to Physical Review Letters

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April 1973

Prepared for the U. S. Atomic Energy Commission under Contract W-7405-ENG-48

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X-RAY PHOTOEMISSION FROM SODIUM AND LITHIUM

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April 1973

Abstract:

XPS spectra of Na and Li metals were obtained in ultrahigh vacuum. Rich plasmon structures were observed. Both photoemission and Auger peaks showed large extra atomic relaxation energies. The Na valence band shows approximately the expected $E^{1/2}$ shape, but it is somewhat anomalous.

We have observed x-ray photoelectron spectra from clean surfaces of sodium and lithium. They show multiple plasmon satellite structure on all peaks, with plasmon energies $\hbar\omega_p(\text{bulk})/\hbar\omega_p(\text{surface})$ in good agreement with theory. Large extra-atomic relaxation¹ energies of ~ 5 eV on core lines and ~ 15 eV on Auger lines were clearly identified. These are in excellent agreement with a simple screening model.^{2,3} The sodium valence-band spectrum was observed. It has approximately the expected free-electron $E^{1/2}$ dependence observed in x-ray emission,⁴ but both the Fermi energy position and the shape of the valence-band edge are anomalous.

The measurements were made with a Hewlett-Packard 5950A ESCA Spectrometer, using monochromatized Al K $\alpha_{1,2}$ x-rays, modified for ultrahigh vacuum work. A sodium film was evaporated from 99.99% pure material onto an aluminum substrate

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and studied at 8×10^{-11} Torr. A Li film was evaporated from 99.95% material onto chromium and studied at 2×10^{-10} Torr. Carbon and oxygen 1s lines from possible surface contamination were undetectable.

Figure 1 shows the Na 1s and Li 1s lines with their characteristic plasmon loss structures. The Na 1s, Na 2s, Na 2p, and Li 1s lines each show a surface plasmon loss peak, plus 5, 4, 3, and 2 bulk plasmon peaks, respectively. Plasmon frequencies ω_p are set out in Table I. The Na ω_p (bulk) values agree quite well with electron-loss results.⁵⁻¹⁰ They also definitely favor the free-electron plasma model prediction of $(4\pi ne^2/m)^{1/2}$ over the corepolarization corrected values⁵,¹¹ in sodium. For lithium our $\hbar\omega_p$ (bulk) values are intermediate between those of Kunz⁵ and of Fellenzer,⁶ and lower than theory. The ratio ω_p (bulk)/ ω_p (surface) is quite close to the theoretical $\sqrt{2}$, although the surface plasmon peaks are poorly resolved.

Core-level binding energies relative to the Fermi level, E_B^F , and to the vacuum level $E_B^V = E_B^F +$ (work function), are set out in Table II. The Fermi energy $E_F(Au)$ was determined from the spectrum of a single crystal of gold. It was checked by evaporating sodium directly onto gold and comparing core-level positions. For both Na and Li the core-level E_B^F values based on $E_F(Au)$ agree very well with x-ray edges.⁴ The E_B^V results are systematically lower than E_B values for free atoms, as estimated either from optical¹² plus x-ray¹³ data or from theoretical free-atom values.¹⁴ The latter two are in excellent agreement (Table II). These results appear to provide particularly clear examples of extra-atomic relaxation. The last column in Table II gives estimates of extraatomic relaxation shifts from a simple "equivalent-cores" model³ that assumes complete screening in the atomic cell, using Mann's¹⁵ Slater integrals. In this model the Li ls shift, for example, would be given by

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$$\Delta E = E_{B}(atom) - E_{B}^{V} \cong \frac{1}{2} \mathscr{F}(ls \ 2s)_{Beryllium}$$

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where $\mathcal{F} = (F^0 - \frac{1}{2}G^0)$ is the two-electron interaction between the K hole and the screening electron. The very good agreement of the last column in Table II with the preceding two strongly supports this model.

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Auger energy shifts from extra-atomic relaxation are even larger, because two holes are screened in the final state. Again a simple "equivalent cores" screening model² predicts extra-atomic relaxation energies in (Table II, Column 8) very good agreement with experiment. The sodium KLL Auger spectrum is shown in Fig. 2. The predicted energy in this case is, for the sodium ¹D line,

$$\Delta E_{ea} = \mathcal{F}(2p \ 3s)_{Al} + \mathcal{F}(2p \ 3p)_{Al} - \mathcal{F}(1s \ 3s)_{Mg}$$

Free-atom Auger energies were estimated, using optical and x-ray data, both from one-electron binding energies² (Columns 4 and 6) and from NaII and NaIII states¹² (Columns 5 and 7). These large Auger extra-atomic relaxation energies should be very useful in surface-physics studies.

The lithium valence-band cross-section was too small to give usable spectra in 20 hours. The valence-band spectrum of sodium (Fig. 3) agrees roughly in shape and width with the x-ray emission spectrum.⁴,¹⁶ An approximate $E^{1/2}$ dependence of intensity on energy is observed, as expected, because sodium is the most free-electron like of the alkali metals.¹⁷ The reduced valenceband width was 3.1 ± 0.1 eV if measured from the high-energy edge, or 2.5 ± 0.1 eV if measured from $E_F(Au)$. The latter value agrees well with the x-ray emission data of Crisp and Williams.¹⁶ Of course x-ray photoemission and x-ray emission and absorption are very different processes, and the spectra need not be identical.

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Our spectrum does not show the excitonic structure observed in x-ray emission^{16,18-20} and attributed to the Nozieres-De Dominicis singularity.²¹ It is unusual, however, both in the position of the valence band relative to $E_{\rm F}({\rm Au})$ and in the gradual slope of the valence-band edge. Both of these features may result from final state relaxation accompanying photoemission from the valence bands. Further work is needed to clarify this question.

FOOTNOTES AND REFERENCES

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Work performed under the auspices of the U. S. Atomic Energy Commission. ⁺Present Address: IBM Thomas J. Watson Research Center, Yorktown Heights, New York 10598.

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FIGURE CAPTIONS

Fig. 1. Sodium and lithium 1s lines, with characteristic plasmon structure. The Li P₁ peak is slightly enhanced in this spectrum from Na 2s contamination.

Fig. 2. Sodium KLL Auger spectrum, with plasmons.

Fig. 3. Sodium valence band. The data has been treated by adding four channels and then performing a sliding linear fit to three adjacent points. Insert shows a wider unsmoothed scan, with plasmons.

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Line	e(p ₁) ^a	E(P ₂)	E(P ₃)	E(P ₄)	e(p ₅)	E(P) ^b	e(p _s)	<u>E(P)</u> /E(P _S)	E(P) _{th}	E(P) ^{*c} th	E(P) ^d e	d E(P _S)e
Na ls	5.8(1)	5.8(1)	5.9(1)	5.9(2)	5.9(2)	5.83(4)	4.0(1)	1.45	5.95 ^{5,e}	5.58 ⁵	5.71(10) ⁵	3.85(10) ⁵
Na 2s	5.9(1)	5.8(1)	5.8(2)	5.8(1)		5.80(2)	4.1(1)	1.41			5.85(5) ⁶	3.82(10) ⁶
Na 2p	5.8(1)	5.9(1)	5.7(2)			5.77(6)	4.0(2)	1.44 ·			5.877	4.01(17) ⁷
Na VB	6.0(2)	5.0(3)				5.5(2)	4.0(3)	1.38			5.4(2) ⁸	
Auger ¹ D (KL ₂₃ L ₂₃)	5.7(2)	5.7(2)	5.7(3)		, 	5.7(1)						
Auger ^l P (KL _l L ₂₃)	5.5(3)	5.9(3)			. 	5.7(2)						
Li ls	7.5(3)	7.3(3)				7.4(2)	5.0(3)	1,48	8.02 ⁵	7.96 ⁵	7.12(10) ⁵ 8 ⁹	4.20(10) ⁵ 4.6 ⁹
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a _{Energy} c ^b Average	of peak N of all th	- Energy ne bulk p	of peak . lasmon en	N - l. ergies.							ан (ал. 1997) Тала	

Table I. Plasmon Energies (eV)

Free electron theory corrected for core polarization.

^dOther measurements.

^eSuperscripts denote references in text.

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Line	E ^F B,K	EV b EB,K	$E^{A}(th)$	E ^A (opt) ^e	$\Delta E^{A}(th)$	$\Delta E^{A}(opt)$	$\Delta \mathtt{E}^{A}(\mathtt{calc})$
Na ls	1071.7(1)	1074.0	1079 ^c	1079.1	5.0	5.1	5.3 ^f
Na 2s	63.4(1)	65.7	71.9 [°]	71.1	6.2	5.4	5.0 ^f
Na 2p	30,4(1)	32.7	36.6°	38.1	3.9	5.4	5.1 ^f
Auger ¹ D (KL23L23)	994.2(1)	991.9	975.8 ^d	977.2	16.1	14.7	16.8 ^g
Auger ¹ P (KL ₁ L ₂₂)	954.7(1)	952.4	937.1 ^d		15.3		16.8 ^g
Li ls	54.8(1)	57.2		64.9		7.7	6.4 ^f

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Table II. Binding, Auger, and Relaxation Energies (eV)

^aBinding energy or kinetic energy relative to $E_{F}^{(Au)}$.

^bWork-function corrections of 2.3 eV (Na) or 2.4 eV (Li) were made.

^CFrom Ref. 14.

d Estimated from binding energies, without relaxation.

^eFrom data in Refs. 12 and 13.

f Estimated as in Ref. 3.

^gEstimated as in Ref. 2.

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Fig. 1



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Fig. 2





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