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ON THE ANGULAR DISTRIBUTION AND THE POLARIZATION OF THE DEUTERON IN THE REACTION  $P + P \rightarrow D + t$

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### Publication Date

1951-07-12

UCRL/393

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Contract No. W-7405-eng-48

ON THE ANGULAR DISTRIBUTION AND THE POLARIZATION  
OF THE DEUTERON IN THE REACTION  $P + P \rightarrow D + \pi^+$

K. M. Watson and C. Richman

July 12, 1951

Berkeley, California

ON THE ANGULAR DISTRIBUTION AND THE POLARIZATION  
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Radiation Laboratory, University of California  
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Cross sections at several angles for the reaction  $P + P \rightarrow \pi^+ + D$  have been measured<sup>1,2,3</sup>. For energies not too far above threshold (i.e., for those at which the cross sections have been measured) and in the center-of-mass frame of reference there is reason to believe<sup>4,5</sup> that only orbital S- and P-states of the meson with respect to the deuteron contribute appreciably to the cross section.

There will be no S- and P-wave interference in the differential cross section, so for the moment we can restrict our considerations to just the contribution from meson P-waves. Combining the intrinsic angular momentum of the deuteron with the P-wave orbital motion, we obtain a total angular momentum  $j = 0, 2$  ( $j = 1$  is forbidden by parity considerations<sup>5</sup> if the meson is assumed to be pseudoscalar, in accord with present evidence). These values of  $j$  arise from the  $^1S_0$  and  $^1D_2$  states, respectively, of the initial p-p system.

The wave functions for the orbital motion of the meson will then be described by the spherical harmonics  $Y_1^m(\theta, \phi)$

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(with the polar axis defined by the proton beam direction) and those of the internal coordinates of the deuteron by  $\sum_1^\mu (\mu=1, 0, -1)$ . Then the linear combinations of these,  $Q_0$  and  $Q_2^0$ , which transform as a scalar and as the spherical harmonic  $Y_2^0$  will arise from the  $^1S_0$  and  $^1D_2$  states, respectively, of the initial p-p system.

The asymptotic form of the outgoing wave will be proportional to

$$\frac{e^{iqr}}{r} \left[ Q_0 + \eta Q_2^0 \right], \quad (1)$$

where  $q$  is the meson momentum and  $\eta$  is a parameter. Writing  $\eta = \epsilon e^{i\sigma}$ , the resulting angular distribution of the mesons will be

$$F(\theta) = 1 + \frac{\epsilon^2}{2} + \sqrt{2} \epsilon \cos \sigma + 3 \cos^2 \theta \left( \frac{\epsilon^2}{2} - \sqrt{2} \epsilon \cos \sigma \right) \quad (2)$$

Including an arbitrary amount of contribution to the cross section from mesons in S-states, we arrive at the following generalization of the cross section given previously<sup>4</sup>:

$$d\sigma = \sigma_0 \frac{d\Omega}{4\pi} \left[ 1 + a \frac{q^2}{\mu^2 c^2} F(\theta) \right], \quad (3)$$

where  $F(\theta)$  is given by Eq. (2) and  $\mu$  is the mass of the meson.

Near threshold it is assumed that  $a$  is a constant and that  $\sigma_0$  is just  $q$  (from the phase space factor) times a constant.

Integrating, we obtain the total cross section

$$\sigma = \sigma_0 \left[ 1 + a(1 + \epsilon^2) \frac{q^2}{\mu^2 c^2} \right], \quad (4)$$

which is of the specific form given previously<sup>4</sup>, with  $a(1 + \epsilon^2) \equiv b$ . We see that  $\epsilon^2$  represents the relative amounts of  $^1S_0$  and  $^1D_2$  contribution of the initial p-p system to the P-wave partial cross section.

The approximate value,  $b = 8$ , was deduced earlier<sup>4</sup>. This value is consistent with the observed angular distribution<sup>3</sup>, which however implies that  $b$  cannot be much less than 8. From the experimental cross sections<sup>3</sup>, it would then follow that at 340 Mev  $\sigma_0 \simeq 5(10)^{-29} \text{ cm}^2$ . The observed angular distributions are consistent with a pure  $\cos^2 \theta$  dependence from the meson P-states, which would imply  $\cos \sigma = -1$ ,  $\epsilon^2 = 2$  (the quoted experimental errors imply that these relations are also quite approximate). Of more positive significance is the conclusion from the observed angular asymmetry that the large contribution from the  $^1D_2$  state of the initial protons implies, independently of a specific model, that non-central interactions are involved<sup>6</sup> (since the deuteron is mostly in an S-state).

Eq. (1) suggests that the deuterons observed at an angle  $(\Theta, \Phi)$  may be polarized<sup>7</sup>. Indeed, the direction of the polarization is given by the average values of the components of

intrinsic angular momentum of the deuteron, which are

$J_x = -R \sin \Phi$ ,  $J_y = R \cos \Phi$ ,  $J_z = 0$  (i.e., at right angles to the plane defined by the direction of the protons and the deuterons).

The degree of polarization (considering just the P-wave component of the cross section),  $R'$  ( $0 < R' < 1$ ), is given by the numerical magnitude of

$$R = 3\sqrt{2} \mathcal{E} \sin \sigma \cos \Theta \sin \Theta \left[ F(\Theta) \right]^{-1}, \quad (5)$$

where  $F$  is given by Eq. (2).  $R$  vanishes for a pure  $\cos^2 \theta$  angular distribution.

The detection of this polarization, if non-vanishing, would appear feasible. The beam of deuterons produced is more intense than the accompanying meson beam<sup>8</sup>, which has been used to study meson scattering<sup>9</sup>. A possible means of detection would be by means of (d, p) scattering, the polarization effects of which have been used by Wouters<sup>10</sup> to measure polarization in (n, p) scattering.

This work was performed under the auspices of the Atomic Energy Commission.



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