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Journal

Journal of Physics Conference Series, 200(1)

ISSN

1742-6588

Authors

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Publication Date

2010

DOI

10.1088/1742-6596/200/1/012192

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Peer reviewed

Journal of Physics: Conference Series 200 (2010) 012192

Magnetotransport measurements on $Nd_{1-x}Ca_xB_6$

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Abstract. We performed measurements of the Hall effect, electrical resistivity, and magnetization on random alloys of $Nd_{1-x}Ca_xB_6$ (x < 0.4) in the 2 to 300 K temperature range, and in magnetic fields of up to 5 T. We find that a large anomalous contribution to the Hall effect, observed in NdB₆ single crystals in the paramagnetic regime, decreases sharply with Ca content. In addition, the electron concentration decreases as expected for a divalent substitution for trivalent Nd. We tentatively relate these observations to changes in the Fermi surface volume upon doping.

1. Introduction

Hexaboride compounds have a simple chemical and structural configuration. However, these compounds show unusual physical properties, more so if they are synthesized with divalent rare-earth or alkaline-earth metal cations. In this paper we report results of magnetic and electrical transport measurements on $Nd_{1-x}Ca_xB_6$ alloys, where some of the rare-earth atoms are substituted with divalent Ca atoms. Our aim is to see how a decreasing valence, and the variations in the Fermi surface (FS) that follow from it, affect the properties of these alloys.

The rare-earth hexaboride compound NdB₆ crystallizes in a cubic lattice with CsCl arrangement of Nd and B₆ octahedra. Its magnetic and electronic properties have been the subject of quite intensive research, both experimental and theoretical, for the last few decades. NdB₆ orders in an A-type collinear antiferromagnetic (AF) structure below $T_N \approx 8$ K [1]. The ground state of the Nd⁺³ ions (J = 9/2) is split in a cubic crystal field into two Γ_8 quartets and a Γ_6 doublet [2,3]. The first excited energy is approximately 135 K above the ground Γ_8^2 state. A competition between crystalline–electric field (CEF) and ferro–quadrupolar interactions gives rise to the low–field magnetic anisotropy in NdB₆ that is much weaker than isotropic magnetic exchange interaction. Band structure calculations indicate that 4f levels are rather deep in NdB₆ [4]. The experimentally found frequency branches of the de Haas-van Alphen effect [5] in the AF phase can be well reproduced by calculated Fermi surfaces.

Recently, the topology of FS within the paramagnetic phase of NdB₆ has been explored [6]. It resembles the FS observed in LaB₆ [7], but with additional weak correlations. The FS consists of six large ellipsoids, centered at the X points of the Brillouin zone, which slightly overlap in the ΓMX plane. Upon alloying with divalent atoms, we expect the FS ellipsoids to shrink and eventually not to overlap as the number of conduction electrons decreases. Indeed, we find experimentally that the electron concentration becomes smaller as we substitute Ca for Nd in

NdB₆. We also find that the anomalous part of the Hall coefficient drops nearly to zero with alloying for $x \leq 0.1$. In our previous studies we encountered that the anomalous Hall coefficient in paramagnetic NdB₆ is much larger than the ordinary one and is independent of temperature [8].

2. Experiment

Single crystals of $Nd_{1-x}Ca_xB_6$ were grown from stoichiometric amounts of hexaboride components in Al flux. We checked its composition by electron probe microanalysis. All measurements were performed on small single crystals with approximate dimensions $0.3 \times 1 \times 5.0$ mm³. These lengths are along the axes of the cubic structure. We have used the same crystals in electrical transport and magnetic measurements. In this way, we expect to avoid domain and sample-shape related effects when comparing results of different experiments. We measured the electrical resistivity and Hall effect with a six probe method. Contact leads (25 μ m gold wire) were soldered to the sample using pure indium. The Hall resistivity ρ_H was measured as a function of magnetic field, from -2 T up to 2 T, for all experimental points. In addition, the variation of ρ_H with magnetic field, up to 5 T, was checked at several temperatures. The magnetization measurements in the temperature range from 2 to 300 K and in magnetic fields of up to 5 T were made using the same geometry with a commercial SQUID magnetometer.

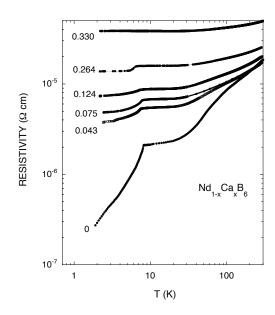


Figure 1. Resistivity as a function of temperature for $Nd_{1-x}Ca_xB_6$ single crystals. An abrupt change in the slope for T < 10 K corresponds to the Néel temperature. The solid lines are guides to the eye.

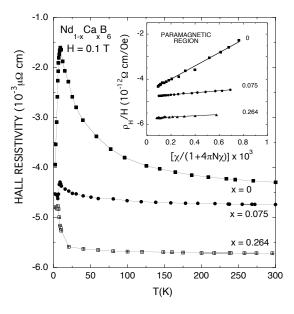


Figure 2. Hall resistivity vs. temperature for various $Nd_{1-x}Ca_xB_6$ single crystals. The inset shows ρ_H/H in the paramagnetic region plotted against the effective susceptibility for the same samples. The solid lines are linear fits to experimental points.

3. Results and Discussion

How the resistivity ρ of Nd_{1-x}Ca_xB₆ single crystals varies with temperature in the 2 to 300 K range is shown in Figure 1. The resistivity drops sharply below T_N as magnetic order sets in. CEF effects give rise to a broad shoulder, observed around 70 K [9]. Temperature-independent residual resistivity becomes larger with increasing content of Ca as a consequence of random

International Conference on Magnetism (ICM 2009)	IOP Publishing
Journal of Physics: Conference Series 200 (2010) 012192	doi:10.1088/1742-6596/200/1/012192

substitution of Nd by Ca atoms. We note that the magnetic part of the resistivity decreases with increasing x.

We now turn to the Hall effect results. Data from some samples are given in Hall resistivity ρ_H versus temperature plots in Fig. 2. At low fields, ρ_H increases as T decreases down to T_N ; ρ_H drops as T decreases below T_N . The Hall resistivity follows quite closely the magnetization of the samples. The low-field magnetization curves, for the same $Nd_{1-x}Ca_xB_6$ single crystals as in Fig. 2, are plotted versus temperature in Fig. 3. To interpret these results, we use for the Hall resistivity the phenomenological expression : $\rho_H = R_o B + R_s 4\pi M_s$, where R_o is the normal Hall coefficient, R_s is the anomalous Hall effect (AHE) coefficient, B is the applied magnetic induction, and M_s is the spontaneous magnetization. In the paramagnetic region $M_s = \chi H$, where χ is the magnetic susceptibility and H is the applied magnetic field. Using this relation we obtain: $\rho_H/H = R_o + 4\pi \chi^* [R_s + R_o(1-N)]$, where $\chi^* = \chi/(1+4\pi N\chi)$, which includes the effects of the demagnetization field, is the effective susceptibility, and N is the demagnetization factor. The total Hall resistivity is linear in H; however it has a normal part and a part which depends on the magnetic susceptibility. Plotting ρ_H/H versus χ^* , we obtain the linear behavior shown in the inset of Fig. 2. This implies that both Hall coefficients are independent of T. The y-axis intercepts yield the values of R_o and the slopes of the straight lines give values which are proportional to R_s . Their variation with Ca content is plotted in Fig. 4.

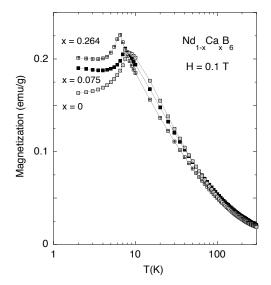


Figure 3. Low-field magnetization as a function of temperature for the same $Nd_{1-x}Ca_xB_6$ single crystals whose Hall resistivity is shown in Fig. 2. The solid lines are guides to the eye.

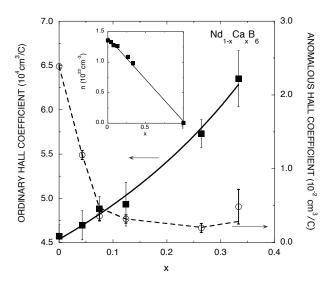


Figure 4. Ordinary (left axis) and anomalous (right axis) Hall coefficients in the paramagnetic region for $Nd_{1-x}Ca_xB_6$ single crystals. The solid and dashed lines are to guide the eye. The inset shows a linear fit to the variation of the electron concentration with Ca content.

As expected for a substitution of trivalent ions by divalent ones, the effective electron concentration n in $Nd_{1-x}Ca_xB_6$ decrease as the content of Ca increases. We find that this decrease is linear in x for x < 0.4. This is shown in the inset of Fig. 4. The linear fit to n(x)reproduces well the value of the electron concentration reported for CaB_6 [10]. Unpublished results of high–magnetic field de Haas van Alphen measurements on $La_{1-x}Sm_xB_6$ single crystals show that the FS volume decreases upon substitution of La by divalent Sm atoms [11]. In addition, this decrease in volume was found to be proportional to the decrease in the number of conduction electrons. For x = 0.10, a multi-connected FS of $La_{1-x}Sm_xB_6$ changes into isolated pockets. One might expect a similar effect in $Nd_{1-x}Ca_xB_6$ alloys. Therefore, for some (small) Ca content in NdB₆, conduction electron ellipsoids would no longer overlap. This should affect transport properties since the Fermi surface changes radically then. Indeed, as seen in Fig. 4, the AHE coefficient decreases with alloying by more than one order of magnitude for $x \leq 0.1$. The AHE values for x > 0.10 are comparable to those reported for other paramagnetic rareearth compounds [12]. However, a large anomalous contribution to the Hall effect for $x \leq 0.1$ in the paramagnetic phase and its strong variation upon alloying are puzzling. The intrinsic AHE in magnetic materials is very sensitive to details of the electronic band structure and gets its significant contributions from band crossing. A calculation of the anomalous velocity for NdB₆, evaluating the Berry curvature tensor over the Fermi surface, would be desirable. In fact, NdB₆ might provide a clean test for some of the current theories of the anomalous Hall effect.

Acknowledgments

This work was supported in part by grant MAT2008/03074 from the Spanish Ministerio de Ciencia e Innovación and by NSF grant DMR-0854781. Additional support from Diputación General de Aragón (DGA-IMANA) is also acknowledged.

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