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1 Photothermal Supercapacitors with Gel 2 Polymer Electrolytes for Wide Temperature 3 Range Operation

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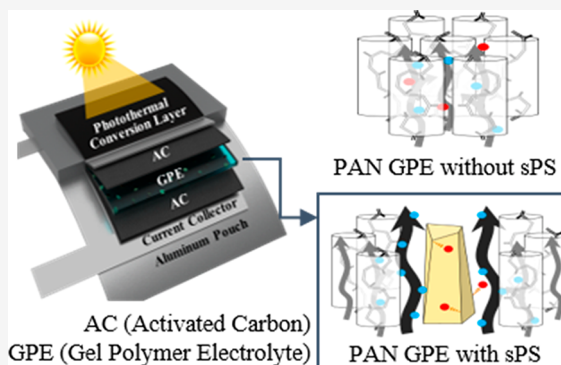
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11 **ABSTRACT:** To deliver electrochemical energy over a wide range of
12 temperatures, this work studies a new gel polymer electrolyte (GPE)
13 for photothermal supercapacitors operating from -60 to 65 °C. The
14 GPE consists of polyacrylonitrile mixed with an active filler sodium
15 polystyrenesulfonate that simultaneously improves ionic conductivity
16 and traps metal cation impurities. The self-discharge rate and impurity
17 diffusion coefficients in the supercapacitors were lowered by the active
18 filler to minimize energy loss in hot environments. For cold settings,
19 the devices were packaged with a photothermal conversion layer that
20 increased the internal cell temperature and raised the energy density to
21 $94 \mu\text{Wh}/\text{cm}^2$ at the output power density of $0.4 \text{ mW}/\text{cm}^2$ in a -60 °C
22 chamber. The combination of our improved GPE and photothermal
23 conversion increased the stored energy and thereby extended the
24 operational time of a motor driven from a cold start by a supercapacitor, demonstrating a high-performance design suitable for
25 harsh environments.



AC (Activated Carbon)
GPE (Gel Polymer Electrolyte)

PAN GPE without sPS
PAN GPE with sPS

20 **A**s electrochemical energy storage devices are widely
21 used in numerous applications from portable elec-
22 tronics to electric transportation, the cell designs are
23 being improved to function in different outdoor settings
24 including cold and hot environments. At low temperature,
25 electrochemical cells show decreased energy and power
26 densities due to sluggish ionic transport, whereas at high
27 temperature, device degradation is triggered by unwanted
28 chemical reactions leading to a short cycle life and accelerated
29 self-discharge. To tackle these issues, the engineering of
30 electrolyte composition^{1–3} can play a major role in expanding
31 the device operational range and stability, especially in electric
32 double layer (EDL) supercapacitors, which are more tolerant
33 of extreme temperatures than batteries.^{4,5} Thus, this work
34 explores a novel gel polymer electrolyte (GPE), consisting of a
35 synergistic combination of a polyacrylonitrile matrix and an
36 ion-exchange resin filler, to advance supercapacitor perform-
37 ance across the temperature range from -60 to 65 °C.

38 Among the polymers used for GPE, polyacrylonitrile (PAN)
39 shows excellent mechanical stability and flame resistance,^{6,7}
40 ideal for merging the electrodes' separator and the electrolyte
41 matrix into one film. However, the ionic conductivity of PAN
42 is limited at low temperature due to crystallization. Prior works
43 have modified PAN with solvent plasticizers or ceramic fillers
44 such as TiO_2 , Al_2O_3 , and graphene oxide to increase the

amorphous phase and lower the glass transition temperature of
the composite for facilitating ion transport.^{8,9} Such a
composite strategy is well-known, but previous studies focused
on <0 °C conditions and did not consider the ramifications at
elevated temperatures >25 °C. It has been observed that
raising ionic conductivity often leads to an increase in self-
discharge,^{10,11} a phenomenon manifested by spontaneous
potential decay in the open-circuit state, resulting in a loss of
stored energy. To decouple this trade-off between ionic
conductivity and self-discharge and in turn enable super-
capacitors that work well at both low and high temperatures,
we choose an active filler based on the cation-exchange resin
sodium polystyrenesulfonate (sPS)¹² to mitigate a dominant
self-discharge mechanism.

The combination of PAN and sPS is introduced for the GPE
to simultaneously increase the ionic conductivity and suppress
self-discharge in the supercapacitors. The sPS filler uniquely
serves the dual functions of creating additional diffusion paths

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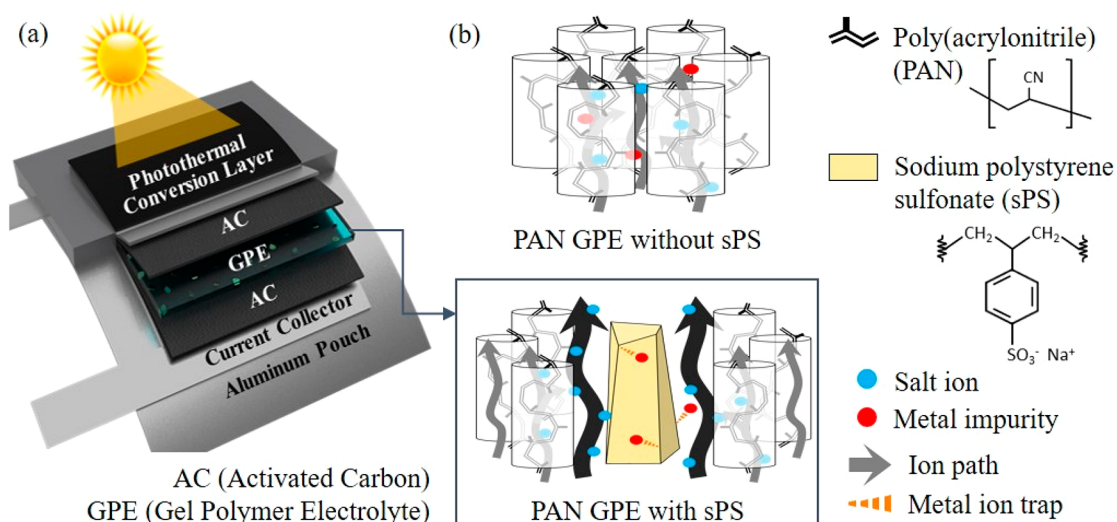


Figure 1. (a) Schematics of the photothermal supercapacitor packaged in an Al pouch. (b) Composition of the gel polymer electrolyte. The addition of ion-exchange resin sPS increases ion transport paths and scavenges metal impurities.

63 for electrolyte ions and scavenging transitional metal
64 impurities, the key contributor to the diffusion-controlled
65 self-discharge process. While the cation-exchange resin sPS is
66 typically used in wastewater treatment to trap heavy metal
67 cations,¹³ here we leverage its preferential binding with high-
68 oxidation-number species ($\geq +2$) to impede the movement of
69 impurities and thereby reduce self-discharge reactions. Mean-
70 while, the transport of electrolyte cations with a low oxidation
71 number of +1 is not affected by sPS.

72 After characterizing the properties of the new GPE, we
73 incorporate it into supercapacitors with activated carbon
74 electrodes and further package the cells with a photothermal
75 conversion layer of carbon black. The photothermal con-
76 version layer absorbs light and converts the energy into heat
77 due to phonon vibrations, effectively raising the cell temper-
78 ature above the environmental temperature. As such, prior
79 photothermal supercapacitors^{14–19} were demonstrated to
80 function down to $-50\text{ }^{\circ}\text{C}$ but suffer from severe self-discharge
81 and very low energy density. In this work with the improved
82 GPE and photothermal conversion effect, the supercapacitors
83 are shown to operate at $-60\text{ }^{\circ}\text{C}$, which notably is below the
84 freezing point of the electrolyte solvent propylene carbonate
85 ($-48\text{ }^{\circ}\text{C}$). The following study compares the device
86 performance and self-discharge characteristics over a wide
87 temperature range, culminating in a demonstration of using
88 our photothermal supercapacitor to drive a motor from a cold
89 start and extend its run-time, toward the goal of advancing
90 high-performance energy storage for harsh environments.

91 *Properties of the gel polymer electrolyte (GPE).* As illustrated in
92 Figure 1a, the supercapacitor structure consisted of activated
93 carbon electrodes separated by a GPE, enclosed in an
94 aluminum pouch with a photothermal conversion layer on
95 top. The GPE was comprised of polyacrylonitrile (PAN) and
96 tetraethylammonium tetrafluoroborate (TEABF₄) salt dis-
97 solved in propylene carbonate (PC), either with or without
98 sPS. The GPE compositions and preparation procedure are
99 provided in the Experimental Methods section in the
100 Supporting Information (Table S1). Figure 1b presents
101 schematics on the sPS filler's roles in enhancing ion diffusion
102 paths and trapping of metal cation impurities by the sulfonate
103 end groups.

Figure 2a,b shows the cross-sectional morphology of GPEs
without and with sPS, respectively, through scanning electron
microscopy (SEM). The SEM images indicate additional free
volume and increased porosity in the PAN matrix around sPS
filler particles. The changes in surface morphology as a
function of sPS weight percentage are presented in
Supplemental Figures S1 and S2. As the outgassing of solvent
during heat treatment created the porous morphology in PAN,
the GPE without sPS tended to have pores perpendicular to
the film surface, while the GPE with sPS showed porous
structures formed around sPS because the evaporated solvent
was vented to the nearby free volume.

To confirm the improvement in porosity in the GPE with
sPS fillers, the specific surface area and pore size distribution
were extracted from nitrogen adsorption-desorption isotherms
in Figure 2c. Specific surface areas of the GPE with and
without sPS were determined to be 458.7 and 268.9 m²/g,
respectively, a 1.7-fold increase in surface area from
incorporating sPS. The GPE with sPS displayed a clear
hysteresis loop in the relative pressure range from 0.07 to 1.
The shape of the hysteresis loop, generated by differences in
adsorption and desorption paths, was attributed to interactions
with aggregated, nonrigid, plate-like particles.²⁰ The pore size
distribution in Figure 2d was calculated from the isotherms
according to density function theory,²⁰ yielding an average
pore diameter of 6.8 nm in the GPE with sPS and 4.5 nm for
the sample without sPS. The higher specific surface area and
average pore width implied that adding sPS had led to a more
porous morphology in the GPE, and these results correlated
well with the SEM images in Figure 2b.

To determine the ionic conductivity in GPEs, electro-
chemical impedance spectroscopy was carried out to find the
equivalent series resistance R of the films (Supplement Figure
S3). The R_{ES} was used for calculating ionic conductivity σ with
the equation $\sigma = \frac{l}{R \cdot A}$, where σ is the ionic conductivity, l is the
thickness of gel polymer electrolyte, and A is the measurement
area. The ionic conductivity as a function of sPS content is
shown in Supplemental Figure S4. In Figure 2e, at $-60\text{ }^{\circ}\text{C}$, the
ionic conductivity of the GPE with sPS (0.99 $\mu\text{S}/\text{cm}$) was 2
orders of magnitude higher than the GPE without sPS (47 nS/

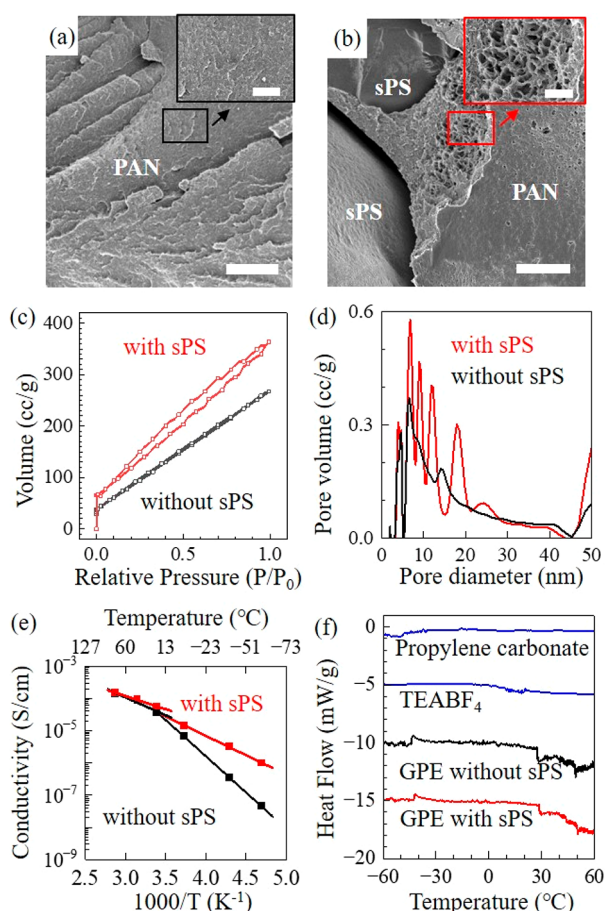


Figure 2. Cross-sectional SEM images of GPE (a) without and (b) with sPS. Scale bars: 20 μm . Insets zoom in on the polymer morphology. Scale bars: 5 μm . (c) N_2 adsorption-desorption isotherms and (d) pore size distribution curves of the GPE with and without sPS. (e) Ionic conductivity as a function of temperature measured on 100 μm thick GPE films. (f) DSC analysis of individual constituents and GPEs. Curves are offset by 5 units each in the y-axis for clarity.

144 cm). The ionic conductivity of the two GPEs converged as the
145 temperature increased to above room temperature.

146 The activation energy E_A of ionic conductivity was extracted
147 by using the Arrhenius equation $\sigma = \sigma_0 \exp\left(\frac{-E_A}{RT}\right)$, where σ_0 is
148 the pre-exponential conductivity, R is the ideal gas constant,
149 and T is the temperature. The fit values of E_A were 104 meV
150 for the GPE with sPS and 194 meV for the GPE without sPS in
151 the temperature range of -60 to 15 $^\circ\text{C}$. Meanwhile, the
152 activation energy values decreased (70–95 meV) after the
153 transition point around 25 $^\circ\text{C}$. The transition point was
154 correlated to the glass transition temperature T_g as measured
155 by differential scanning calorimetry (DSC) in Figure 2f. The
156 solvent PC and salt TEABF_4 did not show any distinct
157 transition in the DSC characteristics, but the GPEs exhibited a
158 glass transition point at 30 $^\circ\text{C}$. While neat PAN is known to
159 show a T_g of ~ 100 $^\circ\text{C}$, the T_g of the composite GPEs was
160 decreased by the salt and solvent serving as plasticizers.⁹ From
161 the analyses of structural and electrical properties, it is evident
162 that sPS enhanced the electrolyte ionic conductivity by
163 increasing the film porosity and free volume, with its effect
164 particularly prominent at below the glass transition temper-
165 ature.

Self-discharge characteristics in supercapacitors with different
GPEs. The self-discharge rate was evaluated for supercapacitors
fabricated using GPE with or without sPS. Self-discharge is
manifested as a loss in cell potential, and it can be caused by
ohmic leakage, charge redistribution, or diffusion-controlled
side reactions.^{12,21–25} The first two contributors can be
minimized, but the diffusion-controlled side reactions are less
avoidable, since carbon-based electrodes with large surface
areas inevitably contain a small amount of metal impurities
even under the most stringent manufacturing process. Below,
we study the effect of using sPS to reduce self-discharge due to
impurity reactions. In Figure 3a,b, the schematics illustrate the
concept that the diffusion of metal cation impurities can be
impeded because of the ion-exchange process in sPS, where the
sulfonate end groups exchange their protons for metal cations
with a high oxidation number. Thus, metal impurities are
trapped and less likely to be transported toward the electrodes
to trigger unwanted reactions leading to self-discharge. To
isolate our study to diffusion effects, the measurement
procedure included a holding period of 60 min at the end
voltage after the charging period, as shown in Figure 3c. The
holding period allowed charge redistribution to equilibrate and
eliminated that mechanism from contributing to self-discharge.
Subsequently, the open-circuit voltage of the supercapacitor
was measured to monitor the voltage decay as seen in Figure
3d,e.

The voltage decay increased with higher temperature as side
reactions were accelerated. The device without sPS showed a
shift in potential decay rate after 20 min at -20 $^\circ\text{C}$, but this
shift might have been due to dimensional changes and contact
problems with thermal contraction in the cold. Meanwhile, the
sPS filler reduced the self-discharge rate across all temperatures
in Figure 3f, where the m values were obtained by fitting the
voltage versus time data in Figure 3d,e to the equation^{10,12}

$$V(t) = V_0 - m\sqrt{t} \quad (1)$$

with

$$m = 2zFAc_i\sqrt{D} / (C\sqrt{\pi})$$

where z is the stoichiometric number of electrons in the
reaction, F is Faraday's constant, A is the electrode area, c_i
is the initial concentration of reacting species, D is the diffusion
coefficient, and C is the device capacitance. The supercapacitor
with sPS showed less self-discharge with m values lower by
>17% compared to its counterpart without sPS. The self-
discharge characteristics of devices using other fillers (ZrO_2
and Al_2O_3) are shown in Supplemental Figure S5, and the
other fillers did not suppress self-discharge. This outcome is in
agreement with our own prior results,¹² and below, we extend
the previous work with additional cyclic voltammetry (CV)
measurements (Figure 3g–j and Supplemental Figure S6) and
analysis method to estimate a key parameter, namely, the
diffusion coefficient of impurities in this system.

The diffusion coefficient can be calculated from diffusion-
controlled CV data according to the Randles-Sevcik equation,
which describes the current controlled by diffusion:

$$i_d = 0.4463\nu^{0.5}zF\left(\frac{nF}{RT}\right)^{0.5}A_c c_i D^{0.5} \quad (2)$$

where ν is the voltage scan rate, z is the stoichiometric number
of electrons in the reaction, F is Faraday's constant, A is the
electrode area, R is the ideal gas constant, c_i is the

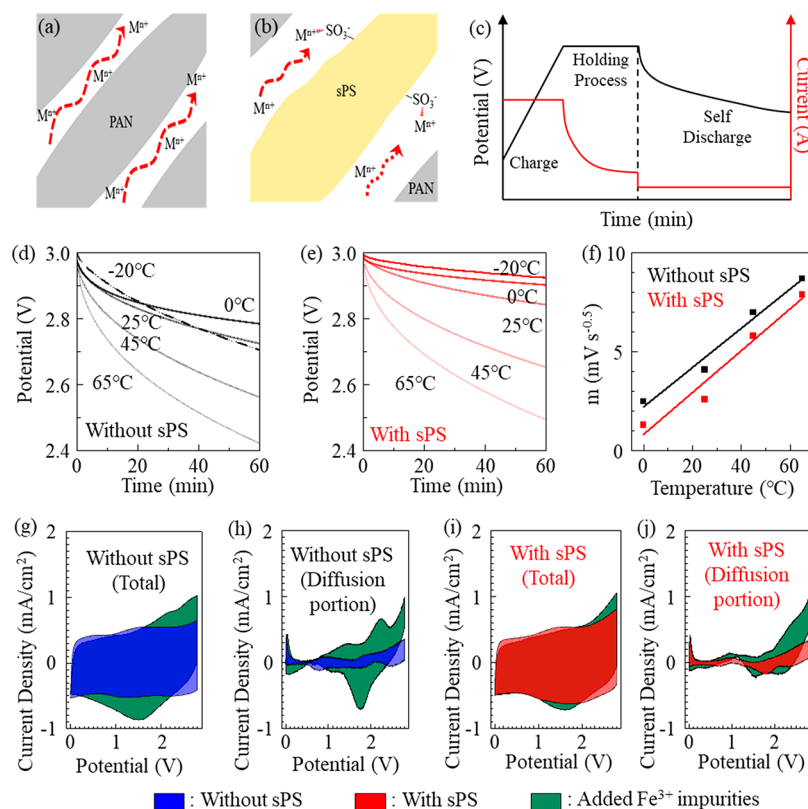


Figure 3. Schematics comparing the movement of metal impurity ions (a) in the GPE without sPS, where metal cations are freely transported across the electrolyte, and (b) in the GPE with sPS, where the sulfonate end groups capture metal cations and suppress their diffusion. (c) The potential (black line) and current (red line) profiles in self-discharge measurements. The potential is held constant for a time period to remove the contribution from charge redistribution. The start of self-discharge is indicated by the dashed line. Upon self-discharge, the voltage drop versus time was recorded at different internal temperatures for supercapacitors (d) without and (e) with sPS in the GPE. (f) The fit value of m in eq 1 as a function of temperature. Comparison of cyclic voltammetry for devices (g,h) without and (i,j) with sPS, measured at 45 °C and a scan rate of 1 mV/s. The contribution of diffusion-limited current (h,j) was extracted from the total current (g,i) by applying eq 3. The different colors represent measurements on a device taken as fabricated (blue and red) versus a modified one, in which Fe^{3+} was intentionally added (green) to check the effect of sPS on impurity mitigation.

223 concentration of the reaction ions, and D is the diffusion
 224 coefficient of the reaction ions.^{17,26} To determine D , the
 225 diffusion current must be extracted first from the total CV
 226 current measured at a set scan rate ν , and then, the unknown
 227 parameters remaining are c_r and D . In our experiment, we
 228 intentionally added a known amount of Fe^{3+} impurities (from
 229 FeCl_3) to the device electrolyte; the added impurities
 230 concentration was 5.7 mM, a small amount relative to the
 231 sPS in the cell to avoid saturating it (sPS has an ion adsorption
 232 capacity¹³ of 12.82 mg/g, and the added impurity concen-
 233 tration was 0.2 mg per 1 g of sPS, much lower than the
 234 maximum capacity). So the lower bound of c_r is known and the
 235 value of 5.7 mM is substituted into eq 2 for estimating D .
 236 For the step of extracting the portion of CV current
 237 contributed by diffusion-controlled ion insertion process (i_d)
 238 as opposed to capacitive EDL effects (i_c), the total CV current
 239 is expressed as^{27,28}

$$240 \quad i(V) = i_c + i_d = k_1\nu + k_2\nu^{0.5} \quad (3)$$

241 or equivalently

$$i(V)/\nu^{0.5} = k_1\nu^{0.5} + k_2$$

242 where k_1 is the capacitive-controlled constant, k_2 is the
 243 diffusion-controlled constant, and ν is the scan rate. A set of
 244 CV data measured at various scan rates were obtained and

245 plotted as $i(V)/\nu^{0.5}$ versus $\nu^{0.5}$ in Supplemental Figure S6, for
 246 which the line slope yielded k_1 and the y -intercept was k_2 . With
 247 these k values, the current contribution from the diffusion-
 248 controlled mechanism was calculated and displayed in Figure
 249 3h,j.

250 Comparing Figure 3h to 3g, or 3i to 3j, the diffusion-
 251 controlled current was less than 30% of the total current in as-
 252 fabricated supercapacitors (red and blue areas). Whereas after
 253 adding Fe^{3+} impurities, the diffusion-controlled current
 254 increased (green areas), although more so in the device
 255 without sPS (Figure 3h) than the one with sPS (Figure 3j).
 256 Notably, redox peaks due to Fe^{3+} were prominent at 1.7 and
 257 2.3 V in the cell without sPS but less obvious with sPS present.
 258 The redox current at 2.3 V was 0.31 and 0.50 mA/cm² for the
 259 supercapacitor with and without sPS. Using these current
 260 values in eq 2, the diffusion coefficient D was calculated to be
 261 1.48×10^{-10} m²/s with sPS and 3.86×10^{-10} m²/s without
 262 sPS. With this quantification, the incorporation of sPS in the
 263 GPE is shown to reduce the diffusion coefficient of impurity
 264 ions by 2.6-fold, trapping the impurities to effectively lower the
 265 self-discharge rate across the entire temperature range.

266 *Effect of photothermal packaging on the supercapacitor*
 267 *performance.* The supercapacitors were packaged in an
 268 aluminum pouch cell, with the surface modified with a
 269 photothermal conversion layer to take advantage of light

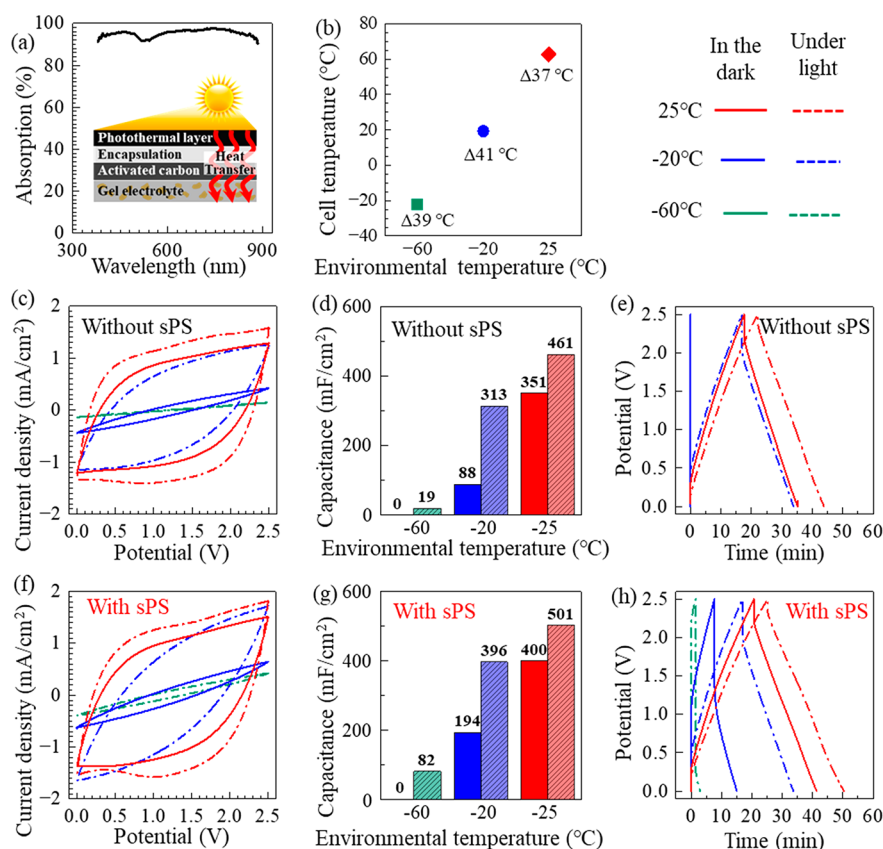


Figure 4. (a) Absorbance of the photothermal conversion layer versus wavelength. The inset illustrates photothermal conversion and heat transfer. (b) Temperature difference due to heat absorption by the photothermal conversion layer under 1 kW/m^2 of light intensity. The legends for colors and solid/dashed lines apply to all the plots in this figure. The plots in the middle row are taken on a cell without sPS and the bottom row on one with sPS. (c,f) Cyclic voltammograms at a scan rate of 5 mV/s . The CV data could not be measured at -60°C in the dark. (d,g) Extracted capacitance from CV data. (e,h) Galvanostatic charge/discharge characteristics at a current density of 0.5 mA/cm^2 .

Table 1. Structures and Metrics of Photothermal Supercapacitors^a

	Electrode	GPE	Potential (V)	Illumination Intensity (kW/m^2)	Temperature ($^\circ\text{C}$)	Photothermal Efficiency (%)	Energy Density ($\mu\text{Wh/cm}^2$)	ref.
a	Activated Carbon	PAN/PC	2.5	1	$-60/-20/25$	~ 52.4	93.5/183/261	This work
b	CNF/Mxene@ SnS_2	PVA/ H_2SO_4	0.6	1	25	~ 86	14.7	16
c	Graphene	PVA/ H_3PO_4	0.7	1	25	~ 18.6	0.38	17
d	LIG + Ppy	PVA/ H_2SO_4	0.7	1	-30	~ 23.5	69.6	19
e	N-MCN@GH	PVA/ H_2SO_4	1.0	3.6	-5	~ 32	41.3	15
f	TiN + MWCNT	PVA/LiCl	0.7	1	-40	~ 62.5	0.68	18
g	Activated Carbon	PVA/LiCl + glycerol	1.0	1	-50	~ 91.5	4.17	14
h	Activated Carbon	PVA/LiCl + glycerol	1.0	1	-20	~ 91.5	6.94	14

^aCNF (cellulose nanofiber), Mxene (two-dimensional metal carbides and nitrides), LIG + Ppy (laser-induced graphene and polypyrrole), N-MCN@GH (n-doped mesoporous carbon-nanosphere-intercalated 3D graphene hydrogel), MWCNT (multiwalled carbon nanotube), PVA (poly(vinyl alcohol)).

energy that could be harvested from the surroundings. The photothermal conversion layer of carbon black absorbed over 90% of the incident light as seen in Figure 4a and converted the light energy into heat due to lattice vibrations. Carbon black had a small heat capacity²⁹ ($0.71 \text{ J g}^{-1} \text{ K}^{-1}$) to readily pass heat to other components within the pouch. The cell temperature was monitored through a thermocouple placed inside the pouch, and the temperature differences between the environmental temperature and the cell temperature were

shown in Figure 4b. With the incident light at 1 kW/m^2 (spectral characteristics of the light source provided in Supplemental Figure S7), the temperature inside the pouch was raised by 37–41 $^\circ\text{C}$ compared to the outside temperature. The temperature differences were similar because the conversion mechanism was the same regardless of the starting temperature. Without the photothermal conversion layer, the pouch temperature was increased by at most 6 $^\circ\text{C}$ compared to the environment. The photothermal conversion efficiency,

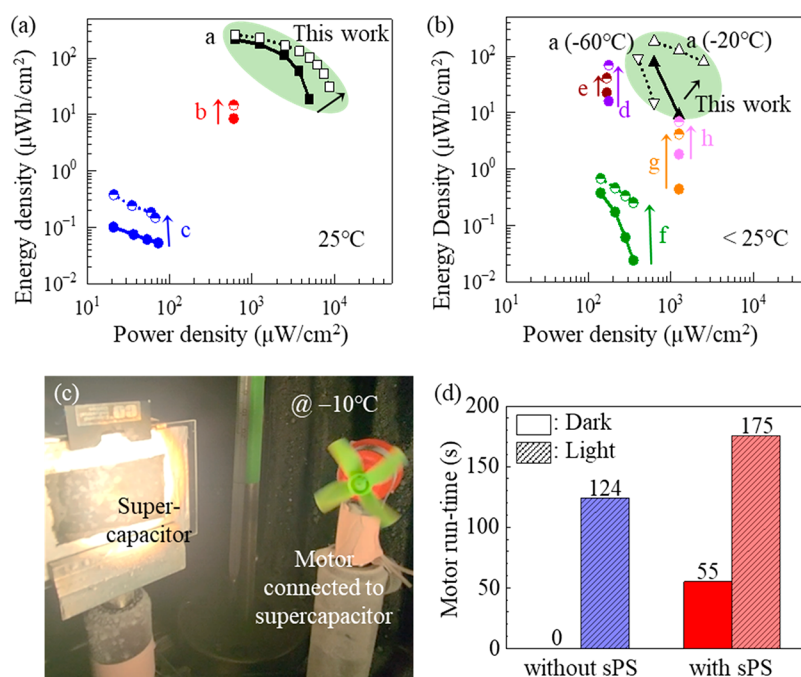


Figure 5. Comparing the performance of photothermal supercapacitors operating at (a) room temperature and (b) below room temperature. The letter labels correspond to the prior works listed in Table 1. The solid and open symbols indicate measurements done in the dark and under light with intensity shown in Table 1, respectively. At $-60\text{ }^{\circ}\text{C}$ in the dark, the devices were not working, and thus, there is no corresponding data. Meanwhile, the solid triangle symbols showed the device performance in the dark at $-20\text{ }^{\circ}\text{C}$. (c) Photograph of a photothermal supercapacitor driving a motor at $-10\text{ }^{\circ}\text{C}$. (d) The operating time of the motor as a function of the energy supply from supercapacitors working under different conditions.

288 calculated according to eq S1 with the parameter values listed
 289 in the Experimental Methods section (Table S2), was
 290 determined to be $\eta_{\text{PT}} \approx 54\%$. This efficiency was in the
 291 middle level among the photothermal supercapacitors being
 292 compared in Table 1, since our photothermal layer was made
 293 with only carbon black, unlike other references which use
 294 plasmonic materials^{14,30} that increased light absorption and
 295 phonon coupling for higher heat generation. In other works,
 296 the photothermal conversion materials were incorporated in
 297 the electrodes, but this design would require the encapsulation
 298 to be transparent while there are limited choices with clear
 299 encapsulation materials. Here, we chose to separate the
 300 photothermal layer from electrodes so the devices could be
 301 fabricated with conventional aluminum pouches typically used
 302 for pouch cell encapsulation. The low specific heat of
 303 aluminum pouch encapsulation (0.89 J/K) has increased the
 304 heat transfer to the supercapacitor components and recovered
 305 the conversion efficiency.

306 The characteristics of photothermal supercapacitors in the
 307 dark and under 1 kW/m^2 light intensity at environmental
 308 temperatures of -60 , -20 , and $25\text{ }^{\circ}\text{C}$ were compared in Figure
 309 4c–e for the device without sPS and in Figure 4f–h for the
 310 device with sPS in its GPE. At each temperature setting, the
 311 CV curves showed current increases under illumination,
 312 because the photothermal conversion effect raised the internal
 313 temperature that increased the ionic conductivity of the GPEs
 314 and in turn improved the device capacitance. The CV data
 315 were used to find the capacitance by the equation of $C = i\Delta t/\Delta V$,
 316 where i is the current, Δt is the measurement time, and
 317 ΔV is the potential window. The extracted capacitance in
 318 Figure 4d,g showed that generally supercapacitors with sPS
 319 achieved higher capacitance due to higher ionic conductivity
 320 over all temperatures and lighting conditions.

The photothermal conversion effect was important for
 321 enabling operation at low temperature. For example at $-60\text{ }^{\circ}\text{C}$,
 322 the devices in the dark did not work, because the electrolyte
 323 solvent PC has reached the freezing point of $-48\text{ }^{\circ}\text{C}$.
 324 Meanwhile, under 1 kW/m^2 light, the cell capacitance was
 325 increased to 19 F/cm^2 without sPS and 82 mF/cm^2 with sPS,
 326 respectively, due to photothermal conversion effectively raising
 327 the internal cell temperature to $-23\text{ }^{\circ}\text{C}$ (Figure 4b). The
 328 photothermal effect induced more change in the device
 329 without sPS than the counterpart with sPS because of the
 330 higher activation energy for the ionic conductivity of the GPE
 331 (Figure 2e). The galvanostatic charge-discharge (GCD)
 332 characteristics in Figure 4e,h provided an additional view on
 333 the equivalent series resistance (ESR) at the start of charge/
 334 discharge step. The voltage drop due to ESR was decreased
 335 under illumination; for example, at an environmental temper-
 336 ature of $-20\text{ }^{\circ}\text{C}$, the voltage drop was reduced from 1.05 V in
 337 the dark (solid blue line) to 0.45 V with a photothermal effect
 338 (dashed blue line), indicating improvement in ESR correlated
 339 to enhanced ionic conductivity.

As the supercapacitor capacitance was increased by sPS and
 341 photothermal effect, the cycling stability was maintained after
 342 10 000 charge–discharge cycles at $25\text{ }^{\circ}\text{C}$, as shown in
 343 Supplemental Figure S8, where the capacitance retention was
 344 100%. At an internal cell temperature of $65\text{ }^{\circ}\text{C}$, the cycling
 345 stability was $>95\%$ after 2000 cycles in the device with sPS and
 346 only 90% in the one without sPS.

Figure 5a,b compares the energy versus power densities of
 348 this work to prior photothermal supercapacitors listed in Table
 349 1. The solid symbols are data taken in the dark, while the open
 350 symbols represent the performance under illumination. To the
 351 best of our ability, in our literature search on state-of-the-art
 352 photothermal supercapacitors, the lowest environmental 353

354 temperature where the device stayed operational was $-50\text{ }^{\circ}\text{C}$
355 at 1 kW/m^2 illumination.¹⁴ Here with our improved GPE with
356 sPS, our photothermal supercapacitor was demonstrated to
357 work at an even lower temperature of $-60\text{ }^{\circ}\text{C}$ with the same
358 intensity of 1 kW/m^2 illumination as previous work while
359 achieving an areal energy density of $94\text{ }\mu\text{Wh/cm}^2$ at the output
360 power density of 0.4 mW/cm^2 (inverted open triangle symbols
361 in Figure 5b). At $-20\text{ }^{\circ}\text{C}$, the photothermal conversion effect
362 increased the device energy density to $183\text{ }\mu\text{Wh/cm}^2$ at the
363 power output of 0.6 mW/cm^2 . Here, our results exceeded prior
364 photothermal studies and were enabled by the expanded
365 potential window and increased ionic transport in our GPE
366 electrolyte as well as the higher film thickness^{31,32} of our
367 activated carbon electrodes.

368 As a proof-of-concept demonstration in Figure 5c, our
369 photothermal supercapacitors were used as the power source
370 for driving a fan motor in a cold chamber of $-10\text{ }^{\circ}\text{C}$. The
371 charging/discharging of the supercapacitor was controlled by a
372 single-pole-double-throw (SPDT) slide switch and a potentiostat
373 (circuit schematics and operational voltage characteristics
374 in Supplemental Figure S9). After the potentiostat charged the
375 supercapacitor at a constant current of 10 mA to reach 3 V , the
376 potentiostat was set to the open-circuit mode for measuring
377 voltage, and the supercapacitor was connected to the motor to
378 discharge current for powering the motor (Supplemental
379 Video S1). This procedure was carried out in the dark and
380 under 1 kW/m^2 illumination. Figure 5b shows the operation
381 time of the fan motor, being dependent on the energy stored in
382 the supercapacitor as a function of electrolyte composition and
383 photothermal effect. The supercapacitor without sPS was not
384 able to drive the motor in the dark but was improved by the
385 photothermal effect to power the motor for 124 s . The
386 supercapacitor with sPS could operate the motor both in the
387 dark and under illumination, demonstrating a capacity to
388 extend the motor operational time to 175 s with the assistance
389 of photothermal conversion.

390 In summary, this work has demonstrated the performance of
391 a GPE with sPS as an active filler in supercapacitors to improve
392 capacitance and suppress self-discharge over a wide temper-
393 ature range. The composite film served as a separator and
394 solid-state electrolyte in one, leveraging the increase in porosity
395 and ion-exchange mechanism due to sPS to enhance
396 electrolyte ionic transport and hinder impurities from under-
397 going side reactions. The diffusion coefficient of impurities was
398 lowered by nearly 3-fold through sPS, leading to better energy
399 retention at elevated temperatures. This sPS GPE offers an
400 economical approach to make the devices more robust against
401 impurities inevitably present in manufacturing processes. The
402 addition of a photothermal conversion layer on the super-
403 capacitor pouch cell further increased the cell performance at
404 low temperature, facilitating operation at a minimum environ-
405 mental temperature of $-60\text{ }^{\circ}\text{C}$. With the combination of sPS
406 GPE and photothermal conversion, the supercapacitors
407 reached state-of-the-art energy and power densities exceeding
408 other reports on EDL photothermal supercapacitors. The
409 supercapacitor structure in this work showed a promising
410 design for high-performance electrochemical energy supplies in
411 applications that span from very cold to hot environments.

412 ■ ASSOCIATED CONTENT

413 Data Availability Statement

414 The data that support the findings of this study are available
415 from the corresponding author upon reasonable request.

416 ■ Supporting Information

The Supporting Information is available free of charge at
<https://pubs.acs.org/doi/10.1021/acseenergylett.3c00207>.

(1) Experimental Methods. (2) Scanning electron
microscopy images of electrolyte films as a function of
active filler contents. (3) Electrochemical impedance
spectroscopy data of gel polymer electrolyte with and
without sPS measured at 0.1 Hz to 100 kHz at -60 ,
 -40 , 0 , 22 , 45 , and $75\text{ }^{\circ}\text{C}$ under 0 V . (4) Ionic
conductivity and device self-discharge characteristics as a
function of filler concentration and (5) filler types. (6)
CV data measured at $65\text{ }^{\circ}\text{C}$ under voltages of 2.5 and 3 .
(7) Spectra of the light source, the photothermal
conversion layer, and the photothermal conversion
layer on an aluminum pouch. (8) Cycling stability of
photothermal supercapacitors with and without sPS at
 25 and $65\text{ }^{\circ}\text{C}$. (9) The measured potential across the
supercapacitor driving a motor (PDF)
Comparison of fan motor operating time with and
without sPS (MP4)

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458 Author Contributions

C.S., L.Y., and T.N.N. designed the experiments and analyzed
the data. C.S. conducted the fabrication and measurement of
the devices. H.L. and P.L. contributed to the pouch cell
fabrication process. The principal investigator T.N.N. conceived
the gel polymer electrolyte concept and supervised the project.
All authors contributed to discussions and writing of the manuscript.

466 Notes

The authors declare no competing financial interest.

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