## **Lawrence Berkeley National Laboratory**

**Recent Work**

**Title** NISOL secondary accelerator scoping study

**Permalink** <https://escholarship.org/uc/item/96t37587>

**Author** Gough, R.A.

**Publication Date** 1998-11-01

LBNL-42511

. . ~

 $\ddot{\phantom{a}}$ .,j

·.'

# ERNEST ORLANDO LAWRENCE BERKELEY NATIONAL LABORATORY

<sup>~</sup>\ . . ' .. . "' ; .....

# NISOL Secondary Accelerator Scoping Study

R.A. Gough, J.W. Staples, and M.S. Zisman Accelerator and Fusion Research Division

November 1998



#### **DISCLAIMER**

This document was prepared as an account of work sponsored by the United States Government. While this document is believed to contain correct information, neither the United States Government nor any agency thereof, nor the Regents of the University of California, nor any of their employees, makes any warranty, express or implied, or assumes any legal responsibility for the accuracy, completeness, or usefulness of any information, apparatus, product, or process disclosed, or represents that its use would not infringe privately owned rights. Reference herein to any specific commercial product, process, or service by its trade name, trademark, manufacturer, or otherwise, does not necessarily constitute or imply its endorsement, recommendation, or favoring by the United States Government or any agency thereof, or the Regents of the University of California. The views and opinions of authors expressed herein do not necessarily state or reflect those of the United States Government or any agency thereof or the Regents of the University of California.

#### LBNL-42511

## **NISOL Secondary Accelerator Scoping Study**

R.A. Gough, J.W. Staples, and M.S. Zisman

Accelerator and Fusion Research Division Ernest Orlando Lawrence Berkeley National Laboratory University of California Berkeley, California 94720

November 1998

This work was supported jointly by the Director, Office of Energy Research, Office of High Energy and Nuclear Physics, Division of Nuclear Physics, of the U.S. Department of Energy under Contract No. DE-AC03- 76SF00098, and by the Physics Division of Oak Ridge National Laboratory.

### **Table of Contents**

1. Introduction and Overview

2. Radio Frequency Quadrupoles

3. Heavy Ion Linac Systems

4. Cryogenic Systems

5. Heavy Ion Beam Transport Systems

6. Diagnostics

7. Conclusions

8. Acknowledgments

ii

#### **NISOL Secondary Accelerator Scoping Study<sup>1</sup>**

November, 1998

#### **1. Introduction and Overview**

The secondary accelerator system of a next-generation ISOL facility (NISOL) must provide cw beams of heavy ions with a wide variety of masses at energies up to 15 MeV/n. Because the primary mission of the facility is to provide radioactive, secondary beams, high efficiency and flexibility are also key requirements. As a user facility, the design must be highly reliable and robust. An upgrade path to extend the performance to higher energies is also an important consideration. The secondary accelerator can accept beam from the NISOL target/ion source system or from a stable-beam ion source. These injected beams can be mass analyzed in a high resolution isobaric mass analyzer, or can be delivered directly to the secondary accelerator. In either case, the majority of reaction products in the NISOL target have been removed from the beams prior to injection into the secondary accelerator. To maximize the intensities of the radioactive beams, initial acceleration must be accomplished with charge state  $q = 1$  for masses up through  $A = 140$ . For higher masses, initial acceleration is accomplished with  $q = 2$ . The design ion for this accelerator system has a mass  $= 140$ , and a charge state  $= 1$ . Only one stripping stage is assumed as a reasonable tradeoff between intensity and efficiency. For the design ion, this still permits over 10% of the beam injected into the accelerator system to be delivered the user experiment.

Modem linac technology provides the most cost effective approach to meeting these requirements. We propose to use a linac system comprising two room temperature radio frequency quadrupoles (RFQs) using split coaxial resonators, followed by a heavy ion linac using independently-phased, superconducting resonator cavities. Superconducting cavities are now a well established technology that can be used for beams with energies down to about 20 keV/n. The RFQs in this design efficiently boost the energy from approximately  $2$ to 22 keVIn. The proposed design is based on existing, proven technology; further developments leading to the availability of multiple-charge-state ion sources with high "species efficiency" and superconducting resonators with higher accelerating gradients, could result in more efficient designs.

The results of a survey<sup>2</sup> of available stripping data are shown in Fig. 1. These data are displayed as mean charge state *versus* atomic number for various different stripping energies up to >10 MeV/n. The experimental data are "binned" into in bands corresponding to different stripping energies. The trends are very systematic, with the mean charge state almost independent of atomic number at the lowest stripping energies, and increasing significantly with atomic number for higher stripping energies. For accelerator design optimization, it is convenient to parameterize these experimental data by the solid lines shown at 0.01, 0.02, 0.05, 0.1, 0.2, 0.5 and 1 MeV/n. Using this parameterization, a plot of accelerator length versus stripping energy exhibits a broad minimum around 0.5 MeV/n (see Fig. 2). This energy is sufficient to elevate the charge-to-mass ratio *qlm* to at least

<sup>&</sup>lt;sup>1</sup>This report was prepared by Lawrence Berkeley National Laboratory for ORNL.

<sup>2</sup> Data compiled by M. McMahan, privvate communication.



Figure 1. Summary of experimental stripping data. The mean charge state is shown as a function of atomic number for various stripping energies.



Figure 2. Optimization of stripper location

 $0.126$  (U<sup>30+</sup>). For this study, we have chosen 0.5 MeV/n as the preferred location for a foil stripper.

Provision for a second, non-equilibrium stripper following the second RFQ is recommended for flexibility, but the linac is designed to accommodate all anticipated requirements with a single stripper at  $0.5 \text{ MeV/n}$ , which is sufficient to achieve the highest possible beam intensities at reasonable cost.

A schematic layout of the proposed secondary accelerator and beam transport system is shown in Fig. 3. In this scheme, the beam accelerates to a maximum energy of  $15 \text{ MeV/n}$ and is delivered to a representative suite of experimental areas. By extending the building, a performance upgrade to higher energies and additional space for an expanded experimental program is straightforward. The principal components are briefly described in the following sections.

#### **2. Radio Frequency Quadrupoles**

The two RFQs, operating at 12 MHz, incorporate a split coaxial resonator (SCR) geometry, similar to the 25 MHz RFQ used in the KEK/University of Tokyo RIB complex, and similar to the structure recently investigated at ANL. The first RFQ structure is 3 m in length and accelerates the beam from  $\sim$ 2 to 10 keV/n. The second structure is 6 m in length and continues the acceleration to  $22 \text{ keV/n}$ , sufficient to be further accelerated directly by 4gap, superconducting cavities operating at 24 MHz.

The tank diameter for both RFOs is  $\sim$ 1.5 meters. The internal accelerating structure is divided into modules to ensure mechanical stability and minimize the required rf power needed for a 100 kV peak intervane voltage. The 3-meter-long structure will be divided into two modules, and the 6-meter-long structure into three modules.

3





4

The cw rf power requirements are 20 and 40 kW for each of the two structures, respectively. The thermal power density on the internal cavity surfaces is small, due to the low frequency of operation. Controlled-temperature water cooling will be required. Vacuum requirements for the RFQs are modest, a pressure requirement of  $\sim 10^{-7}$  Torr can be readily met with conventional pumping systems. Most of the gas load will be from outgassing from the internal cavity surfaces.

External bunching is used: the RFQ units include only an accelerator section, accounting for their relatively short length. Other energy trim and bunching cavities together with quadrupole and steering magnets will be needed as part of the RFQ system.

Ions with different mass will exit the mass separator system with different velocities. The cells within the RFQ are of fixed length with spacing tailored to the ion velocity; the velocity profile of the beam must be constant, independent of ion. In particular, there are fixed velocity requirements at the entrance and exit of each RFQ. In order to satisfy this requirement, the RFQs are placed on 300 kV high voltage platforms similar to the one presently in use at the Oak Ridge HRIBF facility. A summary of preliminary RFQ parameters is given in Table 1.



#### 3. Heavy Ion Linac Systems

All of the acceleration beyond the RFQs is provided by independently-phased, quarterwave, superconducting resonators. These systems take the beam from an energy of 22 keV/n to 15 MeV/n. Singly-charged ions are accelerated to 0.5 MeV/n where a stripper is used to elevate the charge state to a charge-to-mass ratio  $q/m \ge 0.126$ . Focusing is provided by superconducting quadrupole triplets. Allowing space for beam diagnostics and a beam extraction system at 2.5 MeV/n, and assuming accelerating gradients achievable with presently available technology, a total length of about 115 m is required.

We describe here a heavy ion accelerating scheme based on three superconducting linac segments distinguished by three rf frequencies, 24, 48 and 96 MHz. A summary of the accelerating schedule with accelerating parameters is provided in Appendix A.

#### 3.1 Superconducting Segment 1

This is a short section consisting of 11 superconducting cavities operating at 24 MHz which raises the energy from 22 to 60 keV/n. The cavities each have 4 accelerating gaps. For the design ion, with a charge to mass ratio  $q/m = 1/140$ , the first four cavities operate with 2 MV/m accelerating gradient and the remaining seven cavities operate with a 3 MV/m accelerating gradient. A separate rf amplifier is used for each cavity. These cavities are 3.3  $\times$   $\beta\lambda/2$  in length and are separated by superconducting quadrupole triplets. Triplet focusing minimizes the beam size in the cavities, reducing the longitudinal-transverse coupling and maintaining the lowest possible emittance of the beam through the linac. The length of each triplet is  $1.5 \times 3.3 \times \beta \lambda/2$ , and can be determined at any point along the linac from the Table in Appendix A. The maximum quadrupole gradient is approximately 300 *Tim.* Some grouping of the quadrupoles is possible to reduce the number of separate designs and the number of power supplies and control points.

The total length required in this segment for cavities and quadrupoles is 5 m. This comprises 11 four-gap cavities and 10 quadrupole triplets, with an additional quadrupole focusing triplet at the entrance to the section. The cavities and quadrupoles are immersed in liquid helium dewars operating at 4.5 deg K. The effective real estate gradient for this segment is 1.1 MV/m. We increase the total length requirement for this section to 6 m to allow for beam diagnostic devices and inter-dewar spaces.

#### 3.2 Superconducting Segment 2

The second segment continues the acceleration of the unstripped beam from 60 to 500 keV/n. It is similar to the first segment, except the four-gap, quarter-wave cavities operate at 48 MHz and have an effective accelerating gradient of 4 MVIm over their effective length. A separate rf amplifier is again used for each cavity. Each cavity is  $3.3 \times \beta \lambda/2$  in length and is separated by superconducting quadrupole triplets. As in the first segment, quadrupole triplets are inserted between each four-gap cavity. They also operate with a maximum gradient of 300 T/m over their effective lengths.

The total length required in this segment for cavities and quadrupoles is 38 m, and comprises 68 four-gap cavities and 68 quadrupole triplets which are immersed in liquid helium dewars operating at 4.5 deg K. The effective real estate gradient for this segment is 1.6 MV/m. We increase the total length requirement for this section to 46 m to allow for beam diagnostic devices and inter-dewar spaces.

#### 3.3 Superconducting Segment 3

The final (poststripper) segment continues the acceleration from 0.5 to 15 MeV/n. Beams in the postripper segment have a charge-to-mass ratio  $q/m \geq 0.126$  (see discussion in section 3.4 below). The configuration of this section is a little different. The quarter-wave, superconducting cavities operate at 96 MHz and contain two gaps. These cavities have an effective accelerating gradient of 4 MV/m. The accelerating length of each cavity is assumed to be 1.2  $\times$   $\beta\lambda$ . Due to the higher velocities and shortness of the cavities, quadrupole triplets are placed between every fifth cavity. The quadrupole gradients are approximately the same as those in the previous sections, so the quadrupole triplet technology can be the same as in all the other quadrupole triplets.

The total length required in this. segment for the two-gap cavities and the quadrupole triplets is 44 m. There are 145 two-gap cavities and 29 quadrupole triplets. The effective real estate gradient for the post-stripper segment is 2.6 MV/m. We increase the total length requirements for the segment to 53 m to allow for beam diagnostics and inter-dewar spaces.

#### 3.4 Stripper section

A foil stripper is used at the 0.5 MeV/n point, located between the second and third superconducting segments. For uranium, the worst case, the most probable charge state is  $U^{30+}$  (q/m = 0.126), with medium-mass ions stripping to q/m up to 0.17. The efficiency for stripping in the medium to high mass range is  $\sim 20\%$ . The lightest ions, such as <sup>11</sup>Li<sup>1+</sup>, will be fully stripped at this energy with very high efficiency  $(\sim 90\%)$ .

A four-magnet chicane is used in the stripper section to analyze the charge state spectrum produced by the stripper and to select a particular charge state for injection into the post stripper segments for further acceleration. This will simplify the tuning of the linac and help to ensure that the maximum possible intensity can be delivered to the user. A spacing of 4 m between linac segments two and three is provided to accommodate both the stripper, the chicane and associated diagnostics.

#### 3.5 Low Energy Beam Pick-off (2.5 MeV/n)

In order to efficiently deliver beams to the low energy experimental areas, another 4-mgap is provided at the 2.5 MeV/n point in the post stripper linac segment. This point is located about 65 m downstream from the start of the sc linac system. As shown in Fig. 2, two 90 deg bends can be used to bend the beam away from the linac axis and direct it into the main beamline used to transport the accelerated radioactive beams to the low energy experimental areas.

#### 4. Cryogenic Systems

The cryogenic systems must provide 4.5 deg K service to the components of the secondary accelerator facility: approximately 225 rf cavities, and 108 quadrupole triplets located in the linac. The rf power supplied to each cavity is only a few watts. The largest consumption of rf power will be by the fast tuners. Most of this power can be dissipated at  $LN<sub>2</sub>$  temperature, not at LHe temperature, to minimize the load on the helium liquifier. These superconducting components are housed in a series of approximately 40 dewars located along the 115 meter length of the linac. It is also economical to plan for cryogenic pumping of the beam transport lines using the cryogenic return line, since this is accomplished without requiring additional cryogenic capacity.

The cryogenic systems required to support the NISOL facility include compressors, a cold box, a distribution system and a dewar system. A cryogenic plant with a room temperature power of  $\sim$  1 MW is required to support the superconducting elements of the facility. Such a system is readily available from the commercial sector. It is customary to separate the compressors in a separate building to isolate the noise and vibration. The compressors will conveniently fit into a 4-m-high building with a footprint of 6 m  $\times$  11 m. Light crane service (5 ton) should be provided for maintenance. The cold box has a diameter of3.4 m. It will be located closer to the accelerator system and housed in a 7.3-mhigh building with a footprint of approximately 6 m  $\times$  7 m. Maintenance of this equipment will also require a 5 ton crane.

The cryogenic components of the linac are housed in a series of approximately 40 dewars. Cold surfaces within and between the dewars provide an excellent vacuum for the accelerator systems. Sufficient space is provided between dewars for diagnostics or other beam access that may be needed. Standard distribution lines for the working fluids are provided between the compressors, cold box, dewars, and beamlines.

#### 5. Heavy Ion Beam Transport Systems

The general layout of the NISOL secondary accelerator systems (see Fig. 2) forms a "Ushaped" configuration with the accelerator systems on one leg and the beam transport systems on the other. The representative suite of experimental systems established at this time by the NISOL planning team is used to define a preliminary set of requirements for these beam transport systems. Radioactive beams must be transported from three points along the accelerator chain: fully accelerated beam (maximum energy, 15 MeV/n), partially accelerated beam (maximum energy, 2.5 MeV/n), and "unaccelerated" beam. The layout allows for future expansion to higher energies by extension of the building, linac and beam transport systems.

The 15 MeV/n system accepts beam from linac and transports it to one of 5 experimental areas. In order to transport a beam of <sup>238</sup>U<sup>30+</sup> ( $q/m = 0.126$ ), the line is designed for a rigidity  $B\rho = 4.44$  T-m. All other beams would have higher  $q/m$  and therefore lower rigidity. A mass  $m = 140$  beam will have a rigidity of about 4 T-m, so designing this line to accommodate uranium beams can be done for essentially the same cost. Standardized

8

magnet designs can be used, and the layout provides adequate space for diagnostics and easy-to-tune beam transport. Dipole magnets with 45 deg bends and straight-through ports could be used to facilitate tuning the beam to the different experimental areas. These straight-through ports also allow placement of Faraday cups in the 180 deg bend section. Together with a cup at the exit of the linac, these can be used to facilitate tuning up the linac and ensuring the beam is properly matched into the beamline.

The 2.5 MeV/n system is used to transport beam to experiments utilizing lower energy beams that don't require operation of the full linac system. The pick-off point for this beamline is approximately 65 m downstream from the start of the superconducting linac system. Four 45 deg magnets are used to turn the beam away from the linac axis and direct it back along the beamline leading toward two of the experimental areas. The maximum magnetic rigidity of these beams is 1.8 T-m.

The unaccelerated beam transport system is used to deliver radioactive beams to an experimental area for mass traps and other apparatus for material science and biomedical studies. These beams have about 100 keV of energy,  $q/m \ge 1/140$ , and a maximum rigidity  $B\rho = 0.5$  T-m.

A preliminary estimate of the number and types of magnets required for these three . beamline systems is given in Table 2. No switching magnets are listed for the 15 MeV/n beam transport line pending further definition of the experimental layout.



The beamlines are cryogenically pumped using the 20 deg K return system from the  $\sim$ main linac cryosystem.

#### 6. Diagnostics

Because the primary purpose of the systems described here is to accelerate *secondary*  beams, the intensities of these beams will vary widely, and for some experiments be very low. Additionally, it is anticipated that there will be a large number of ions and operational modes that will require frequent retuning of the accelerator during its operational lifetime. While a detailed study of beam diagnostics is beyond the scope of the present study, it is clear that efficient operation will require a facility that is well instrumented with beam diagnostics and that these instruments cover a wide dynamic range of intensities from  $\sim 10^2$  $-10^{10}$  ions/sec.

Several features have been incorporated at the conceptual level to facilitate efficient operation of the facility. These include, for example: the use of a separate, stable-beam ion source that can be used for initial commissioning and to tune out "pilot" beams; a charge state analyzer in the stripper area to select a particular charge state for injection in to the post stripper; and the incorporation of a large number of spaces between dewars in the linac which permit frequent access to the beam for placement of diagnostics. Additionally, the general layout of the external beamlines provides ample space for locating diagnostics devices.

#### **7. Conclusions**

The design of this accelerator requires no substantial advances in technology. The design is flexible and robust, using a large number of short, independently-phased cavities, and providing for a wide range of ion species acceptance and energy variability. The triplet focusing between cavities reduces the beam size and renders it round in the accelerating gaps, thereby reducing the potential for emittance blowup. The total length required to house the proposed accelerator is about 140 m. This includes space for the RFQs, the superconducting linac segments, inter dewar spacing, and the various short sections of beamline that are required from the exit of the isobaric separator to and including the 180 deg turn in the 15 MeV/n transport line. The estimated total installed power requirement for the accelerator and beam transport systems is about 2.2 MW (not including the user experimental apparatus). R&D to develop reliable high gradient resonators could lead to a more efficient, compact design, but at this time, it is prudent to base the designs on established performance capabilities.

As discussed in section 1, a second, optional, non-equilibrium stripper following the second RFQ is recommended to add flexibility, but the linac is designed to accommodate all anticipated requirements with a single stripper at 0.5 MeV/n; this approach will achieve the highest possible beam intensities and reliability at reasonable cost.

Some multiplexing of multiple ion species should be possible. For ions of similar rigidities, the quadrupoles can be tuned to a compromise field, and cavity gradients switched to the appropriate levels for the two ions. The cavities have a Q-value on the order of  $10^8$ , indicating that it would take on the order of 1 second to significantly change the cavity gradient. Multiplexed operation may be particularly valuable for setting up a new experiment which requires only occasional or low levels of beam. The bandwidth of the focusing system is quite large, and rigidity ranges of 2:1 should be possible without retuning.

In summary, the development of a secondary accelerator system to meet the requirements of NISOL appears to be straightforward and for the most part can be based on proven technology with established performance parameters.

#### **8. Acknowledgments**

The authors are pleased to acknowledge the support and guidance from Fred Bertrand and Jerry Garrett at Oak Ridge National Laboratory. We also gratefully acknowledge engineering support from Missy Ogan (ORNL) for help with the civil construction aspects of the proposed facility, and from Shlomo Caspi, Ron Scanlan and Ron Y ourd (LBNL) for engineering support. Support from Margaret (Peggy) McMahan with the stripping data and discussions with J. Kalnins on beam transport aspects is also gratefully acknowledged.

#### **Appendix A: Linac Acceleration Table**

The following table summarizes the superconducting linac parameters. Incidental gaps in the accelerator are not included. A single stripper is assumed at 0.5 MeV/n. The design ion is  $q/m =$ 1/140; ions with mass  $m > 140$  are accelerated in the prestripper sections with charge state  $q=2+$ . The geometry of each resonator is specific to its exact energy. A filling factor of 0.5 is assumed in the superconducting structures. Four-gap resonators are assumed in the prestripper segments and two-gap resonators in the post stripper sections.





Strip to  $q/m > 0.126 = 30/238$  at 500 keV/n

## Strip to  $q/m > 0.126 = 30/238$  at 500 keV/n 2-gap cavities; length = 1.2 \*  $\beta \lambda/2$

Focus between every 5th cavity 29 focusing devices Cavity Length =  $\beta \lambda /2 \times 1.2$  Non-focus intercav length  $=$  5.00

 $L_foc/L_cav = 1.50$ 









This section can be continued as a future upgrade to 25 MeV/n

 $\hat{\beta}$ 

 $\bar{\lambda}$ 

ERNEST ORLANDO LAWRENGE EERKELEY NATIONAL LASORATORY ONE GYOLOTION ROAD | BERKELEY, GALIFORNIA 94720

J.

 $\left\langle \cdot \right\rangle_{\alpha}$ 

Repared for the U.S. Department of linery under Contract No. DB-ACOS-ASHOODS