Title
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Permalink
https://escholarship.org/uc/item/9q01j3jv

Journal
Physical review letters, 93(8)

ISSN
0031-9007

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Publication Date
2004-08-18

DOI
10.1103/physrevlett.93.081801

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Peer reviewed
Bound on the Ratio of Decay Amplitudes for $B^0 \to J/\psi K^0$ and $B^0 \to J/\psi K^0$
We have measured the time-dependent decay rate for the process $B \rightarrow J/ψ K^{(*)0}$ in a sample of about $8 \times 10^6$ $Y(4S) \rightarrow B \overline{B}$ decays collected with the BABAR detector at the PEP-II asymmetric-energy $B$ factory at SLAC. In this sample we study flavor-tagged events in which one neutral $B$ meson is reconstructed in the $J/ψ K^{(*)0}$ or $J/ψ K^{(*)0}π^0$ final state. We measure the coefficients of the cosine and sine terms in the time-dependent amplitude:

$$A(t) = A_0 + A_1 \cos(\frac{2\pi}{\tau} t) + B_1 \sin(\frac{2\pi}{\tau} t)$$

where $A_0$, $A_1$, and $B_1$ are the coefficients, $\tau$ is the decay time, and $t$ is the time since the $B$ meson was produced.
The standard model of electroweak interactions describes CP violation in weak interactions of quarks by the presence of a complex phase in the three-generation Cabibbo-Kobayashi-Maskawa (CKM) quark-mixing matrix [1]. In this framework, the CP asymmetries in the proper-time distributions of neutral B decays to \( J/\psi K^0 \) and \( J/\psi K^0 \) are directly related to the CP-violation parameter \( \sin 2\beta \) [2]. The time-dependent CP asymmetries for \( J/\psi K^0 \) and \( J/\psi K^0 \) are of opposite sign and, to a very good approximation, equal in magnitude [3]. The decay \( B^0 \rightarrow J/\psi K^0 (B^0 \rightarrow J/\psi K^0) \) proceeds through the CKM-favored, color-suppressed decay \( B^0 \rightarrow J/\psi K^0 \) [4] followed by \( K^0 \rightarrow K^0 \) \((K^0 \rightarrow K^0)\). The so-called wrong-flavor \( B^0 \) decay amplitude to the opposite strangeness final state \( B^0 \rightarrow J/\psi K^0 \) is expected to be negligible in the standard model [3]. Interference between a wrong-flavor amplitude and the favored amplitude can alter the relation between the CP asymmetries, \( A_{CP} \), for the \( J/\psi K^0 \) and \( J/\psi K^0 \) final states. In general, a difference between \( A_{CP}(J/\psi K^0) \) and \( -A_{CP}(J/\psi K^0) \) of more than a few times \( 10^{-3} \) requires a wrong-flavor amplitude [3]. A limit on the CP-odd part of the phase difference between the wrong-flavor amplitude and the favored amplitude can be derived from the measured values of \( \sin 2\beta \) from \( B \) decays to the \( J/\psi K^0 \) and \( J/\psi K^0 \) final states. No test of the modulus of the wrong-flavor amplitude currently exists.

The decay mode \( B^0 \rightarrow J/\psi K^0 \) proceeds via the same quark transition as \( B^0 \rightarrow J/\psi K^0 \). The matrix elements, and therefore the ratio of wrong-flavor to favored amplitudes, are expected to be similar for \( B^0 \rightarrow J/\psi K^0 \) and \( B^0 \rightarrow J/\psi K^0 \) [3]. In this Letter we present a measurement of the ratio of wrong-flavor to favored amplitude for the decay \( B^0 \rightarrow J/\psi K^0 \), from the time-dependent asymmetry, where we use \( K^0 \rightarrow K^- \pi^+ \) to identify the strangeness of the final state. The data sample consists of about 88 \( \times 10^6 \) \( B\bar{B} \) pairs produced in \( e^+e^- \) interactions at the \( \Upsilon(4S) \) resonance, corresponding to an integrated luminosity of 82 \( fb^{-1} \), collected with the BABAR detector [5] at the PEP-II asymmetric-energy collider at SLAC.

Charged particles are detected, and their momenta measured, by a combination of a vertex tracker consisting of five layers of double-sided silicon microstrip detectors, and a 40-layer central drift chamber, both operating in the 1.5-T magnetic field of a superconducting solenoid. We identify photons and electrons using a CsI(Tl) electromagnetic calorimeter. Further charged particle identification is provided by the average energy loss \( (dE/dx) \) in the tracking devices and by an internally reflecting ring imaging Cherenkov detector covering the central region. Muons are identified by their penetration through the iron plates of a magnet flux return.

The analysis method is similar to that of other time-dependent mixing measurements performed at BABAR [6]. We use a sample of events \((B_{J/\psi K^0})\) in which one neutral \( B \) meson is reconstructed in the state \( J/\psi K^0 \) or \( J/\psi K^0 \). The \( J/\psi \) meson is reconstructed through its decay to \( e^+e^- \) or \( \mu^+\mu^- \), and the \( K^0 \) \((\bar{K}^0)\) meson through its decay to \( K^-\pi^+ \) \((K^+\pi^-)\). We examine each event in this sample for evidence that the other \( B \) meson decayed either as a \( B^0 \) or \( B^0 \) (flavor tag).

The pseudoscalar to vector-vector decay \( B^0 \rightarrow J/\psi K^0(892) \) is described by three amplitudes, \( A_0 \), \( A_1 \), and \( A_2 \), for the longitudinal, parallel, and perpendicular transverse polarization [7], respectively, of the vector mesons. In the selection of \( B^0 \rightarrow J/\psi K^0(892) \) there is a small contribution from \( B^0 \rightarrow J/\psi K_s^0(1430) \), whose decay amplitude is denoted with \( A_2 \). The favored decay amplitudes \( A_\lambda(B^0 \rightarrow J/\psi K^+\pi^-) = a_\lambda e^{i\phi_\lambda}e^{-i\phi_\lambda} \) for \( \lambda = 0, \perp, s \) are described by the magnitudes \( a_\lambda \), weak phase \( \phi_\lambda \), and strong phases \( \delta_\lambda \), where \( \lambda = 0, \perp, \perp, s \). The amplitudes for the wrong-flavor decays are given by \( \bar{A}_\lambda(B^0 \rightarrow J/\psi K^+\pi^-) = b_\lambda e^{i\phi_\lambda}e^{-i\phi_\lambda} \). The corresponding decay amplitudes for the charge-conjugate final state \( J/\psi K^-\pi^+ \) are obtained by replacing \( \phi_\lambda \) with \( -\phi_\lambda \), \( b_\lambda \) with \( \bar{b}_\lambda \), \( \delta_\lambda \) with \( \bar{\delta}_\lambda \), and \( \delta_\lambda \) with \( -\bar{\delta}_\lambda \). We assume \( a_\lambda = \bar{a}_\lambda \).

The proper-time distributions of \( B \) meson decays to \( J/\psi K^+\pi^- \) \((J/\psi K^-\pi^+)\), having either a \( B^0 \) or \( B^0 \) tag, can be expressed in terms of the \( B^0 - \bar{B}^0 \) oscillation amplitude and the amplitudes describing \( B^0 \) and \( B^0 \) decays to this final state [8]. The angular-integrated decay rate \( f_+/(f_-) \) to the final state \( J/\psi K^+\pi^- \) \((J/\psi K^-\pi^+)\) when the tagging meson is a \( B^0(B^0) \) is given by

\[
f_{\pm}(\Delta t) = \frac{e^{-|\Delta t|/\tau_{B^0}}}{4\tau_{B^0}^2} \left[ 1 \mp C \cos(\Delta m_s \Delta t) \sin(\Delta m_s \Delta t) \right],
\]

where \( \Delta t = t_{rec} - t_{tag} \) is the difference between the proper decay times of the reconstructed \( B \) meson \( (B_{rec}) \) and the tagging \( B \) meson \( (B_{tag}) \), \( \tau_{B^0} \) is the \( B^0 \) lifetime, and \( \Delta m_s \) is the \( B^0 - \bar{B}^0 \) oscillation frequency. The corresponding decay rates \( \bar{f}_+ \) and \( \bar{f}_- \) for the charge-conjugate final state \( J/\psi K^-\pi^+ \) \((J/\psi K^+\pi^-)\) are obtained by replacing \( C \) with \( -\bar{C} \) and \( S \) with \( -\bar{S} \).
The C and S coefficients are related to the wrong-flavor and favored amplitudes by

\[ C = \frac{a^2 - b^2}{a^2 + b^2}, \quad \text{and} \quad S = \frac{2 \sum \eta a_j b_k \sin(\phi + \delta_k)}{a^2 + b^2}, \]

with \( a^2 \equiv a_0^2 + a_1^2 + a_2^2, b^2 \equiv b_0^2 + b_1^2 + b_2^2, \) and \( \eta = +1(-1) \) for \( \lambda = 0, |1> \). The strong and weak phase differences are given by \( \delta_k = \delta_k^S - \delta_k^A \) and \( \phi = \text{arg}(q/p) + (\phi_b - \phi_a) \), respectively, where \( (q/p) \) contains the weak phase of \( B^0_\bar{B}^0 \) oscillations. The \( \bar{C} \) and \( \bar{S} \) coefficients are given by the same expressions, replacing \( b(\lambda) \) with \( \bar{b}(\lambda) \), \( \delta_k \) with \( \delta_k \), and \( \phi \) with \( \phi \).

In the \( B \to J/\psi K^{*0} \) selection, a \( J/\psi \) candidate must consist of two identified lepton tracks [5] that form a good vertex. The lepton-pair invariant mass must be in the range 3.06–3.14 GeV/c² for muons and 2.95–3.14 GeV/c² for electrons. This corresponds to a ±3σ interval for muons, and, for electrons, accommodates the remaining radiative tail after bremsstrahlung correction [6]. We form \( K^+ \pi^- \) candidate pairs, where the track that is most consistent with being a kaon is assigned to be the kaon candidate. The \( K^+ \pi^- \) pair must have an invariant mass within 100 MeV/c² of the nominal \( K^{*0}(892) \) mass [9]. In the selected mass window the \( K^0_s(1430) \) contributes (7.3 ± 1.6)% of the \( K^+ \pi^- \) events.

The \( B \)-meson candidates are formed from \( J/\psi \) and \( K^+ \pi^- \) candidates with the requirement that the difference

\[ \Delta E = E_B^{\text{cm}} - E_B^{\text{beam}} \]

be between their energy and the beam energy in the center-of-mass frame be less than 30 MeV from zero. The beam-energy-substituted mass \( m_{ES} = \sqrt{(E_B^{\text{cm}})^2 - (p_B^{\text{cm}})^2} \) must be greater than 5.2 GeV/c², where \( p_B^{\text{cm}} \) is the measured \( B \) momentum in the center-of-mass frame. We define a signal region with \( m_{ES} > 5.27 \) GeV/c² to determine event yields and purities, and a sideband region with \( m_{ES} < 5.27 \) GeV/c² to study background properties. If several \( B \) candidates are found in an event, the one with the smallest \( |\Delta E| \) is retained.

A measurement of the asymmetry coefficients \( C, S, \bar{C}, \) and \( \bar{S} \) requires a determination of the experimental \( \Delta t \) resolution and the fraction \( w \) of events in which the flavor tag assignment is incorrect. This mistag fraction reduces the amplitudes of the observed asymmetries by a factor \( 1 - 2w \). Mistag fractions and \( \Delta t \) resolution functions are determined from a sample of neutral \( B \) mesons that decay to final states with one charmed meson (\( B_{Dh} \)) and consists of the channels \( D^{(*)} h^+ (h^+ = \pi^+, \rho^+, \text{ and } a_1^+) \).

The algorithm for \( B \)-flavor tagging is explained in Ref. [10]. The total efficiency for assigning a reconstructed \( B \) candidate to one of four hierarchical, mutually exclusive tagging categories is (65.6 ± 0.5)% for \( J/\psi K^+ \pi^- \) and \( J/\psi K^- \pi^+ \) samples and for the \( B_{Dh} \) sample. Errors are statistical only.

![FIG. 1. Distributions of \( m_{ES} \) for \( J/\psi K^+ \pi^- \) and \( J/\psi K^- \pi^+ \) candidates satisfying the tagging and vertexing requirements. The fit is described in the text.](image)

Effective tagging efficiency \( Q = \sum_i w_i (1 - 2w_i)^2 \), where \( w_i \) are the efficiencies and mistag probabilities, for events tagged in category \( i \), is measured to be \( (28.1 \pm 0.7)\% \).

The time interval \( \Delta t \) between the two \( B \) decays is calculated from the measured separation \( \Delta z \) between the decay vertices of the \( B_{rec} \) and \( B_{tag} \) along the collision \( z \) axis [6]. We determine the \( z \) position of the \( B_{rec} \) vertex from its charged tracks. The \( B_{tag} \) vertex is determined by fitting tracks not belonging to the \( B_{rec} \) candidate to a common vertex, employing constraints from the beam spot location and the \( B_{rec} \) momentum [6]. We accept events with a \( \Delta t \) uncertainty of less than 2.5 ps and \( |\Delta t| < 20 \) ps. The fraction of events satisfying these requirements is 95%.

Figure 1 shows the \( m_{ES} \) distributions of the \( J/\psi K^+ \pi^- \) and \( J/\psi K^- \pi^+ \) candidates that satisfy the tagging and vertexing requirements. The \( m_{ES} \) distributions are fit with the sum of a threshold function [11], which accounts for the background from random combinations of tracks in the event, and a Gaussian distribution describing the signal. In Table I we list the event yields and signal purities for the tagged \( B \to J/\psi K^+ \pi^- \) and \( B \to J/\psi K^- \pi^+ \) candidates. The fraction of events in the Gaussian component of the \( m_{ES} \) fits due to other \( B \) decay modes is estimated to be (1.6 ± 0.4)% based on simulated events.

<table>
<thead>
<tr>
<th>Sample</th>
<th>( N_{tag} )</th>
<th>( P(%) )</th>
</tr>
</thead>
<tbody>
<tr>
<td>( J/\psi K^+ \pi^- ) sample</td>
<td>860</td>
<td>95.5 ± 0.7</td>
</tr>
<tr>
<td>( J/\psi K^- \pi^+ ) sample</td>
<td>856</td>
<td>96.5 ± 0.6</td>
</tr>
<tr>
<td>( B_{Dh} ) sample</td>
<td>25375</td>
<td>84.9 ± 0.2</td>
</tr>
</tbody>
</table>
We determine the $C$, $S$, $\overline{C}$, and $\overline{S}$ coefficients with a simultaneous unbinned maximum likelihood fit to the $\Delta t$ distributions of the tagged $B_{J/\psi K\pi}$ and $B_{Dh}$ samples. In this fit the $\Delta t$ distributions of the $J/\psi K^+\pi^-$ and $J/\psi K^-\pi^+$ samples are described by Eq. (1). The $\Delta t$ distributions of the $B_{Dh}$ sample are described by the same equation with $C = 1$ and $S = 0$. The observed amplitudes for the time-dependent asymmetries in the $B_{J/\psi K\pi}$ sample and for flavor oscillation in the $B_{Dh}$ sample are reduced by the same factor, $1 - 2w$, due to flavor mistags. Events are assigned signal and background probabilities based on the $m_{ES}$ distributions. The $\Delta t$ distributions for the signal are convolved with a common resolution function, modeled by the sum of three Gaussians [6]. Backgrounds are incorporated by means of an empirical description of their $\Delta t$ spectra, obtained from the $m_{ES}$-sideband region, containing prompt and nonprompt components convolved with a resolution function [6] distinct from that of the signal.

There are 48 free parameters in the fit. The fit parameters that describe the signal $\Delta t$ distributions are $C$, $S$, $\overline{C}$, and $\overline{S}$ (4), the average mistag fraction $w$, the difference $\Delta w$ between $B^0$ and $\overline{B}^0$ mistag fractions, and the linear dependence of the mistag fraction on the $\Delta t$ error for each tagging category (12), parameters for the signal $\Delta t$ resolution (8), and parameters to account for differences in reconstruction and tagging efficiencies for $B^0$ and $\overline{B}^0$ mesons (5). The $B_{J/\psi K\pi}$ and $B_{Dh}$ background $\Delta t$ distributions are described by parameters for the background time dependence (8), $\Delta t$ resolution (3), and mistag fractions (8). We fix $\tau_{B^0}$ at 1.542 ps and $\Delta m_d$ at 0.489 ps$^{-1}$ [9]. The determination of the mistag fractions and $\Delta t$ resolution function parameters for the signal is dominated by the large $B_{Dh}$ sample. Background parameters are determined from events with $m_{ES} < 5.27$ GeV/c$^2$.

The fit to the $B_{J/\psi K\pi}$ and $B_{Dh}$ samples yields $C = 1.045 \pm 0.058 \pm 0.035$, $S = -0.024 \pm 0.095 \pm 0.041$, $\overline{C} = 0.966 \pm 0.051 \pm 0.035$, and $\overline{S} = 0.004 \pm 0.090 \pm 0.041$, where the first error is statistical and the second error is systematic. Figure 2 shows the $\Delta t$ distributions and the asymmetries in yields between $B^0$ tags and $\overline{B}^0$ tags as a function of $\Delta t$ for the $J/\psi K^+\pi^-$ and $J/\psi K^-\pi^+$ samples, overlaid with the projection of the likelihood fit result.

We estimate common systematic errors for $C$ ($S$) and $\overline{C}$ ($\overline{S}$). The dominant sources of systematic error are the uncertainties in the level, composition, and time-dependent asymmetry of the background in the selected $B_{J/\psi K\pi}$ sample (0.016 for $C$, 0.017 for $S$), uncertainties in the beam spot location and the internal alignment of the vertex detector (0.016 for $C$, 0.021 for $S$), and the statistics of the simulated event sample (0.016 for $C$, 0.015 for $S$). Another significant contribution to the systematic uncertainty in the cosine coefficients comes from possible differences between the $B_{Dh}$ and $B_{J/\psi K}\pi$ mistag fractions (0.012). The uncertainty in the interference between the suppressed $b \to u c \bar{d}$ amplitude with the favored $b \to c \bar{u} d$ amplitude for the decay modes in the $B_{Dh}$ sample and for certain tagside $B$ decays to hadronic final states [12] contributes to the systematic uncertainty in the sine coefficients (0.019). Finally, there are differences in the angular-integrated efficiency for the $B \to J/\psi K^{*0}(892)$ helicity amplitudes and the $B \to J/\psi K^0(1430)$ amplitude (0.007 for $C$, 0.016 for $S$). The total systematic errors for the cosine coefficients and sine coefficients are 0.035 and 0.041, respectively. Most systematic errors are determined with data and are expected to decrease with larger sample size.

The large $J/\psi K^+\pi^-$ and $J/\psi K^-\pi^+$ samples allow a number of consistency checks, including separation by data-taking period and tagging category. The results of fits to these subsamples are found to be statistically consistent.

The measured values of the cosine and sine coefficients are consistent with $C = \overline{C} = 1$ and $S = \overline{S} = 0$, as expected for no contributions from the wrong-flavor decays $B^0 \to J/\psi K^-\pi^+$ and $B^0 \to J/\psi K^+\pi^-$. We use the measured cosine coefficients $C$ and $\overline{C}$ and assume $|q/p| = 1$ [13] to calculate the wrong-flavor to favored decay rate ratios $\Gamma(B^0 \to J/\psi K^-\pi^+)/\Gamma(B^0 \to J/\psi K^+\pi^-) = |b/a|^2 = 0.022 \pm 0.028 \text{(stat.)} \pm 0.016 \text{(syst.)}$ and $\Gamma(B^0 \to J/\psi K^-\pi^+)/\Gamma(B^0 \to J/\psi K^+\pi^+) = |b/a|^2 = 0.017 \pm 0.026 \text{(stat.)} \pm 0.016 \text{(syst.)}$, where the negative.
central value occurs because $C > 1$. From these measurements the wrong-flavor to favored amplitude ratios for $B \rightarrow J/\psi K^{*0}(892)$ and $B \rightarrow J/\psi K^{*0}(892)$ can be calculated. Using the measured fraction of $B \rightarrow J/\psi K_0^*(1430)$ events contributing in the $B \rightarrow J/\psi K^\mp \pi^\mp$ selection, the upper limits for the decay amplitude ratios at 90% confidence level (C.L.) are found to be $|A(B^0 \rightarrow J/\psi K^0)|/|A(B^0 \rightarrow J/\psi K^{*0})| < 0.26$ and $|A(B^0 \rightarrow J/\psi K^{*0})|/|A(B^0 \rightarrow J/\psi K^0)| < 0.32$. For the single ratio of wrong-flavor to favored amplitude for $B^0$ and $B^0$ combined, we determine an upper limit of 0.25 at 90% C.L.

In conclusion, we observe no evidence for the wrong-flavor decays $\bar{B}^0 \rightarrow J/\psi K^{*0}(892)$ and $B^0 \rightarrow J/\psi K^{*0}(892)$. Together with theoretical information on the relation between the matrix elements for $B^0 \rightarrow J/\psi K^0$ and $B^0 \rightarrow J/\psi K^{*0}$ [3], the results presented here can be used to set a limit on the difference between $A_{CP}(J/\psi K_0^*)$ and $-A_{CP}(J/\psi K^{*0})$.

We are grateful for the excellent luminosity and machine conditions provided by our PEP-II colleagues and for the substantial dedicated effort from the computing organizations that support BABAR. The collaborating institutions wish to thank SLAC for its support and kind hospitality. This work is supported by DOE and NSF (USA), NSERC (Canada), IHEP (China), CEA and CNRS-IN2P3 (France), BMBF and DFG (Germany), INFN (Italy), FOM (The Netherlands), NFR (Norway), MIST (Russia), and PPARC (United Kingdom). Individuals have received support from the A. P. Sloan Foundation, Research Corporation, and Alexander von Humboldt Foundation.

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[4] Charge conjugation is implied throughout this Letter, unless explicitly stated otherwise.