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# Biomimetic Hierarchical Assembly of Helical Supraparticles from Chiral Nanoparticles

<sup>3</sup> Yunlong Zhou,<sup>†,‡</sup> Ryan L. Marson,<sup>§,||</sup> Greg van Anders,<sup>†,§</sup> Jian Zhu,<sup>†</sup> Guanxiang Ma,<sup>†</sup> Peter Ercius,<sup>#</sup>
<sup>4</sup> Kai Sun,<sup>||</sup> Bongjun Yeom,<sup>†,¶</sup> Sharon C. Glotzer,<sup>\*,†,§,||</sup> and Nicholas A. Kotov<sup>\*,†,§,||,⊥</sup>

- <sup>5</sup> Department of Chemical Engineering, <sup>§</sup>Biointerfaces Institute, <sup>II</sup>Department of Materials Science and Engineering, <sup>L</sup>Department of Biomedical Engineering, University of Michigan, Ann Arbor, Michigan 48109, United States
- 7 \*Wenzhou Institute of Biomaterials and Engineering, CNITECH.CAS-Wenzhou Medical University, Wenzhou, Zhejiang 325011,
- 8 People's Republic of China 9 <sup>#</sup>National Center for Electron Microscopy, the Molecular Foundry, Lawrence Berkeley National Laboratory, Berkeley, California
- <sup>11</sup> Department of Chemical Engineering, Myongji University, Yongin, Gyeonggido 17058, South Korea
- 12 Supporting Information

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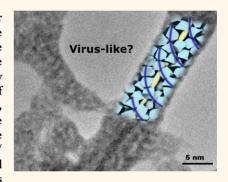
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ABSTRACT: Chiroptical materials found in butterflies, beetles, and other creatures are attributed to biocomposites with helical motifs and multiscale hierarchical organization. These structurally sophisticated materials self-assemble from primitive nanoscale building blocks, a process that is simpler and more energy efficient than many top-down methods currently used to produce similarly sized three-dimensional materials. Here, we report that molecular-scale chirality of a CdTe nanoparticle surface can be translated to nanoscale helical assemblies, leading to chiroptical activity in the visible electromagnetic range. Chiral CdTe nanoparticles coated with cysteine self-organize around Te cores to produce helical supraparticles. D-/L-Form of the amino acid determines the dominant left/right helicity of the supraparticles. The small energy difference between the chiral asymmetry of the amino acid interactions at the surface of the nanoparticles is



amplified by their collective behavior. Coarse-grained molecular dynamics simulations with a helical pair-potential confirm the assembly mechanism and the origin of its enantioselectivity, providing a framework for engineering three-dimensional chiral materials by self-assembly. The helical supraparticles further self-organize into lamellar crystals with liquid crystalline order, demonstrating the possibility of hierarchical organization and with multiple structural motifs and length scales determined by molecular-scale asymmetry of nanoparticle interactions.

**KEYWORDS:** biomimetic nanoparticles, self-assembly, chirality, supraparticles, helices, virus-like nanostructures

xamples of astounding structural complexity of nature's materials raise the challenges of translating biological mechanisms of their synthesis to the industrial scale. 

Inorganic nanoscale components are particularly attractive for many emerging and established technologies due to their special optical and electronic properties and their greater environmental robustness compared to many organic molesse cules. 

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Nature's method to fabricate and diversify its materials is based on the structural hierarchy utilizing a spectrum of small structural units that hierarchically self-organize across multiple scales into ever more sophisticated systems. Therefore, one can ask a question: Can the assembly of small inorganic building blocks with some biomimetic characteristics self-organize into hierarchical systems with a degree of complexity rivaling that found in nature? Answer to this question will also help us understand better the limits of stunning analogies between nanoparticles and globular biomolecules, such as proteins.

Inorganic nanoparticles modified with an amino acid 49 represent one of the simplest biomimetic inorganic building 50 blocks. 9–12 Notably, such particles can also be chiral; that is, 51 their mirror image would not be superimposable with the 52 original. 13 Chirality is ubiquitous in living systems and 53 important in virtually all biological functions. 14 Its role in 54 structural diversification can also be seen in the rapidly 55 expanding spectrum of chiral metal and semiconductor 56 nanomaterials with multiple functionalities. 15–17 Also impor- 57 tant, mathematical relationships between the symmetries of 58 structural units and assemblies as well as their optical 59 manifestations can complement the experimental discovery 60 tools available for evaluation of assembly mechanisms. 18–22

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#### 62 RESULTS AND DISCUSSION

63 Here, we systematically investigate the self-organization of 64 CdTe nanoparticles (NPs) modified with L- or D-cysteine 65 (CYS) as the basic biomimetic building block. Despite a 66 relatively small difference in the energy of interactions 67 associated with chiral asymmetry, 23 it was found that these 68 NPs assemble into helical supraparticles with enantiomeric 69 preference.

D- and L-CYS-stabilized CdTe NPs (denoted as D-NPs and L-71 NPs, respectively) were synthesized by the arrested precipitation method following a standard protocol (see Methods). Note that notation D and L with respect to NPs denotes only the method of synthesis but not the chiral symmetry of the particles. Tetrahedral CdTe NPs could have symmetry that is similar to chiral D or L centers in amino acids when all the four corners have different truncation. Alternatively, the atoms in the apexes of the NPs could also follow the familiar chiral pattern of tetrahedra with four different substituents. In this work, however, we study the effect of chiral surface ligands on the surfaces of the four sides of tetrahedral NPs.

To observe self-organization, the CdTe NPs were 83 precipitated by addition of 2-propanol and centrifuged for 5 84 min, followed by redispersion in deionized water with pH 9.0 in 85 a nitrogen atmosphere (see Methods). The solution of NPs 86 changed color from orange to dark red when left in the dark at 87 room temperature for 8 h, indicating that self-assembly of NPs 88 has occurred. It should be noted that the assembly conditions 89 here are different from previous studies, in which the self-90 assembly was induced by light or different inter-particle 91 forces. 27-30 The conditions here enable the self-limiting 92 assembly at the NP-NP interfaces capped with CYS surface 93 ligands. The assembly process of CYS-stabilized NPs described 94 here was critically influenced by the efficiency of oxygen 95 removal and control over pH value, which are both key factors 96 for successful helical assembly controlled by small anisotropy in 97 NP interactions. When the assembly process was carried out 98 under a high value of pH, i.e., pH 11.0, the final assemblies 99 displayed a morphology of dense NP chains irregularly attached 100 on the surface of Te nanowires. If assembly occurred in an open 101 atmosphere, NPs transitioned to a deeper oxidation state, a 102 greater number of Te nanowires formed, and only a few NPs 103 remained to assemble around them (Figure S1).

Scanning electron microscopy (SEM) showed the formation 105 of nanorods (NRs) as the primary product of the self-assembly 106 process. The lengths of the NRs were ~250 nm for assembly of 107 D-NPs and ~300 nm for L-NPs (Figure 1a,b). The length of 108 NRs can be controlled by the assembly conditions and can be 109 as long as 5  $\mu$ m using NPs synthesized with a molar ratio of 110 CYS/Cd of 1.7:1 in a pH 10.0 solution, indicated in Figure 111 1c,d. The diameters of NRs assembled from D- and L-NPs are ~20 and ~23 nm, respectively, and reveal high uniformity. 113 Both D- and L-NRs self-organize into even higher order lamellar 114 assemblies. Similarly to liquid crystals, the self-organized structures from chiral NRs may also reveal a chiral, for 116 instance, cholesteric, pattern of three-dimensional organization. 117 Figure 1a and b, however, display only a two-dimensional 118 representation of the higher order structures from NRs and are 119 not sufficient to elaborate on the micrometer-scale chirality of 120 the lamellae or potential twist in the mutual orientation of 121 adjacent NRs.

Scanning transmission electron microscopy (STEM) high angle annuar dark-field (HAADF) images (Figure 2a,b,d, e) and

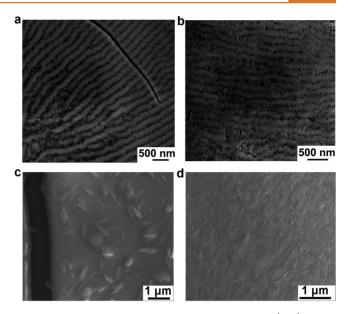


Figure 1. SEM images of NP assemblies from D-NPs (a, c) and L-NPs (b, d).

STEM tomography (Figure 2c,f) showed that NRs have 124 pronounced twist although not as regular as in DNA or some 125 viruses.<sup>6,31</sup> Comparing the organization of the biological and 126 NP systems, we note that the outer shell of viruses also reveals 127 a considerable amount of disorder and may not be perfectly 128 packed especially under room-temperature imaging conditions. 129 On the other hand, the size/shape distribution of NPs does 130 decrease the perfection in the NR packing. Importantly, the 131 self-assembly of NRs is enantioselective, with the twist 132 directions dictated by the starting NPs' chirality. The NRs 133 from D-NPs are left-handed, whereas those from L-NPs are 134 right-handed helices that are easily distinguishable in tomo- 135 graphically reconstructed 3D images from tilt-series STEM-- 136 HAADF images (Figure 2c,f). AFM images also reveal their 137 helical geometry in periodic topographical features typical for 138 helices (Figure 2g,h). We evaluated the handedness of 100 139 samples by AFM images of helical NRs and found consistent 140 dominance of a specific handedness of the helices. The overall 141 yield of left-handed and right-handed helices from D-NPs and L- 142 NPs is ~70% and ~74%, respectively. The average pitch of 143 helices seen in the TEM images is ~10 nm, which is further 144 confirmed by SAXS data, revealing distinct peaks for both D- 145 and L-NRs with  $q = 0.05A^{-1}$ , characteristic of a repetitive 146 structural distance of 12.6 nm (Figure S2).

Nanoscale X-ray energy dispersive spectroscopy (XEDS) 148 (Figure 3a-c) showed a Te/Cd atomic ratio of ~2:1 in the 149 f3 central part of the chiral helical structures; however, the Te to 150 Cd ratio converges to 1:1 for the outside layer. High-resolution 151 TEM (HRTEM) images (Figure 3d,e, Figure S3) indicate the 152 CdTe NPs reside on the outside of the helical NRs; the center 153 part is a single crystal nanowire with a diameter of ~5 nm. The 154 thickness of the twisted CdTe layer is about 10-15 nm. 155 Electron diffraction for the helical NRs reveals the presence of 156 two phases: cubic CdTe and hexagonal Te. The diffraction 157 spots for the CdTe phase correspond to (111) (d = 0.39 nm), 158 (220) (d = 0.22 nm), and (311) (d = 0.19 nm) lattice planes. 159 The diffraction patterns for the central Te-rich part show the 160 maxima from (101) (d = 0.32 nm), (012) (d = 0.23 nm), and 161 (110) (d = 0.20 nm) lattice planes. These data indicate that the 162 helical NRs have unusually complex core-shell morphology, 163

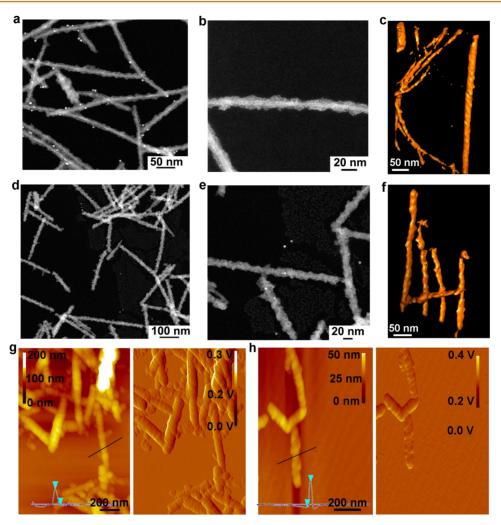


Figure 2. Nanoscale geometry of left- and right-handed helical NP assemblies. STEM—HAADF (a,b) and STEM (c) tomography images of *left*-NRs assembled from D-NPs. STEM—HAADF (a,b) and STEM (f) tomography images of *right*-NRs assembled from L-NPs. (g, h) AFM images of NRs obtained from D-NPs and L-NPs, left: height images, right: amplitude images. Note: In panels a, d, and e, the bright points are gold NPs added as markers for STEM tomography.

164 more sophisticated than what was observed before by Yeom *et* 165 *al.*, <sup>27</sup> with a solid crystalline Te rod in the center and a 166 polycrystalline NP shell on the outside. A gradual growth of the 167 central  $\sim$ 5 nm Te rods/nanowires along the (001) direction is 168 due to the slow oxidation process of CdTe (Figure 3a–f), <sup>32</sup> and 169 thus the control of the oxygen content in the media during the 170 assembly is essential.

From the circular dichroism (CD) spectra (Figure 3h) one can see that the helical NRs have distinct chiroptical activity in the region from 300 to 800 nm, confirming enantioselectivity of helical NRs with mirror-image CD spectra, whereas the precursors to the supraparticles, *i.e.*, NPs, and their early assemblies appearing as NP chains (Figure 4a, Figure S4) show no chiroptical activity in the visible range (Figure 3g). Only CD peaks associated with individual D- and L-CYS moieties in the UV spectral region between 200 and 300 nm appear.<sup>24</sup>

To further understand the mechanism of the assembly process and the driving forces responsible for the enantiose-183 lective transition of D/L-NPs to NRs, we examined the dispersions using different spectroscopy and microscopy 185 techniques at the time points of 8, 24, 48, and 72 h (Figure 186 4, Figures S4–S8). In the first 8 h, TEM bright-field (BF)

images revealed transformation of short NP chains for both 187 enantiomers of CdTe particles (Figure 4, Figure S4). The L-NP 188 assemblies show a broad and weak positive CD peak at ~430 189 nm compared with the NP solution and D-NP assemblies 190 (Figure S6b,d). By 24 h, the chains transformed into short NRs 191 that display chiroptical activity (Figure 4, Figure S6). Neither 192 morphology nor CD spectra for mixed D/L-NP assemblies 193 could show the same changes in 48 h. After 48 and 72 h, both 194 D-NP and L-NP assemblies show self-organization into short 195 helical rods, for which the corresponding CD spectra display 196 peaks and Cotton effects of opposite signs (Figure S6).<sup>24</sup> Note 197 that the intermediate stages en route from single NPs to helical 198 NRs are markedly different from previous observations for 199 assemblies of organic and inorganic helical systems. 18,33-35 The 200 growth of Te cores takes place concurrently with the 201 attachment of CdTe.

The zeta potential decreased with time for both D-NP and L- 203 NP assemblies (Figure S7), indicating the continuously 204 decreasing repulsive interactions between the NPs and the 205 NRs. Such a temporal trend in  $\zeta$  is conducive to the self- 206 limiting aggregation of NPs into supraparticles. The 207 uniformity of NR diameters leading to their hierarchical 208 organization into the lammelar phase (Figure 1a,b) is also 209

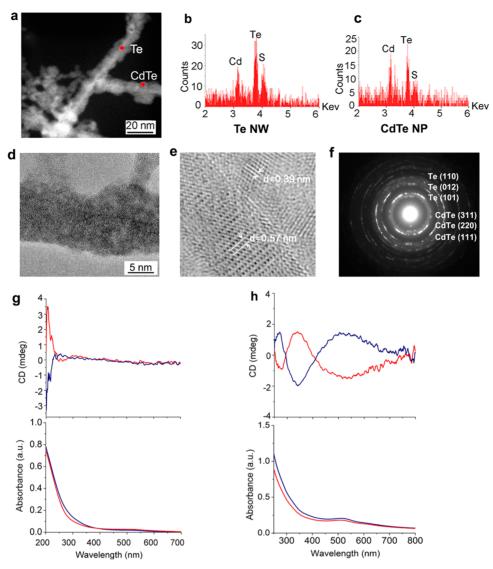


Figure 3. Nanoscale structure and chirality of helical NP assemblies. (a–c) XEDS analysis of peripheral and core parts of *left*- handed helical NRs. (d, e) HRTEM image of peripheral CdTe NPs and Te core. (f) Electron diffraction of the *right*-handed NRs. (g) CD/UV-vis spectra of supraparticle precursors: NP chains obtained after 8 h of assembly of D- (red) and L-NPs (blue). (h) CD spectra of *left*- (red) and *right*-handed (blue) helical NRs obtained after 2 weeks of assembly. Complementary CD/UV-vis spectra of L/D-cysteine are given in Figure S5.

210 indicative of the self-limiting terminal assembly pattern. 211 Photoluminescence (PL) and UV-vis spectra remain essen-212 tially unchanged (Figures S8, Figure S6c,d); the red-shift of the 213 luminescence due to aggregation is (over)compensated by the 214 blue-shift due to oxidation, leading to the formation of Te cores. PL lifetimes increased from 22 ns for D-NPs to 52 ns for the left-handed helix and from 19 ns for L-NPs to 35 ns for the right-handed helix (Figure S9). There could be several reasons 218 for such a noticeable change in PL dynamics. One of them is 219 the decrease of the spatial confinement of the electrons and 220 holes. Another could be the formation of a thin layer of CdS on the surface of CdTe NPs from the decomposition of cysteine surface ligands. The CdS layer can elongate the lifetime of the electron-hole pair by temporary localization of some of the 223 charge carriers. 38,39

Previous reports indicated that the self-organization of helices from organic building blocks arises mainly from the concerted action of multiple types of short-range and long-type range interactions. Hydrogen bonding and hydrophobic interactions are some of the most important forces in the

self-organization of chiral biotic assemblies, *e.g.*, the DNA 230 double-helical structure and chiral peptides—amphiphile nano- 231 fibers. A0,41 Biomolecule-stabilized inorganic NPs can interact *via* 232 a similar set of intermolecular interactions as biological species 233 of the same size, although the van der Waals attraction between 234 them is typically stronger due to higher values of pairwise 235 Hamaker constants. Therefore, we hypothesize that the 236 formation of helical NRs is due to the chirality of interactions at 237 the NP–NP interfaces capped with CYS moieties as opposed 238 to chiral preference of the packing of NP cores, NR 239 twisting, or iented attachment, or screw axis dislocation.

Compared to other types of inter-NP interactions, <sup>47</sup> which <sup>241</sup> include van der Waals forces, dipole—dipole, charge—charge, <sup>242</sup> and charge—dipole interactions, hydrogen bonds, and other <sup>243</sup> interactions between NPs, this chiral contribution may appear <sup>244</sup> to be small, but it is actually appreciable, especially for chiral <sup>245</sup> molecules on the surface of NPs. The energy difference <sup>246</sup> between intermolecular bonds binding two homochiral CYS <sup>247</sup> molecules or two heterochiral CYS molecules is comparable to <sup>248</sup> other forces. According to Kühnle *et al.*, the heterochiral dimer <sup>249</sup>

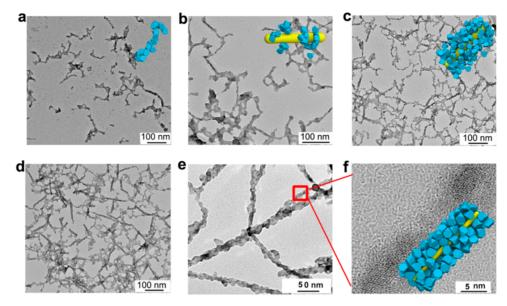


Figure 4. Intermediate stages of self-organization of helical hierarchical assemblies from p-NPs. (a-d) TEM BF images of the formation process of left-handed NRs by assembly of p-NPs in 8 h (a), 24 h (b), 48 h (c), and 72 h (d). TEM BF (e) and HREM (f) images of helical NRs obtained after 72 h of assembly. Inset images are the simulated assembly of helical NRs from a CdTe NP; the central Te core is yellow, and the surrounding CdTe NPs are blue.

of gold-bound CYS was found to be energetically less favorable than the homochiral dimers by *ca.* 0.2 eV, *i.e.*, 20 kJ/mol at 295 kg. The three-point bonding involving carboxyl groups and amine—gold interactions give an LL pair advantage over an LD pair. CYS molecules in our system are bound to the surface of CdTe, and the same three-point intermolecular interactions involving the carboxylic group and amino group coordination to the cadmium atoms on the CdTe surface are likely to take place. The energy cost for breaking the cysteine—Cd²+ tetrahedral coordination bond affected by the rotation of the amino acid segments and/or its chirality can be as high as 117.3 kcal/mol (490.8 kJ/mol); thus its deformation depending on the even seemingly minor alteration of the atomic structures can easily amount to 20 kJ/mol.

For comparison, the energy of dipolar attraction, known to be a significant factor determining the geometry of self-assembled structures for 3.2 nm CdTe NPs, can be estimated to be ca. 9.7 kJ/mol using, for instance, the classical formula for lass aligned dipoles (Supporting Information, Part I). This juxtaposition with dipole—dipole interactions shows the importance of the chiral interactions as a determinant of the assembly geometry, especially considering their cooperative nature for the face-to-face contact between NPs.

To confirm that the molecular-scale chirality of CYS surface 274 ligands can indeed translate into the nanoscale chirality of the 275 NRs, we performed coarse-grained molecular dynamics (MD) 276 simulations of NP self-assembly (Supporting Information Part 277 II). Four primary features of the experimental system, *i.e.*, the 278 excluded volume effects due to the NP shape, the attraction of 279 the NPs to the surface of the Te core, dipolar interactions, and 280 the chiral asymmetry of NP–NP interactions originating from 281 the CYS layer on the NP surface (Figure 5a), included in the 282 simulations. Among others, excluded volume effects are 283 captured *via* a purely repulsive Weeks—Chandler—Andersen 284 potential from the NP surface. To model the axial form of the 285 Te core, we confine particles using a cylindrical harmonic well  $U = \frac{1}{2}kx^2$  (where *x* denotes the radial distance from the *z*-axis)

at the center of the simulation box. To account for the twisting 286 preference induced by the amino acid surface ligand, we include 287 a specially developed chiral potential that acts between 288 interacting faces of the particles. The potential describes the 289 cumulative torque from collective interactions of chiral surface 290 ligands on the NP surfaces; it changes signs for the clockwise 291 and counterclockwise direction of rotation, as can be seen in 292 the description of the potential with respect to face normals 293 given in the SI (Figure S12). For the entire NP the total energy 294 of chiral interactions affecting its motion in the simulations was 295  $\pm 1$ , 10, and 100  $k_{\rm B}T$  (0.41, 4.1, and 41.0 kJ/mol at 295 K), the 296 lower values of which are an underestimate considering the 20 297 kJ/mol at 295 K difference stemming from the chiral 298 asymmetry of CYS. 23

Simulated assemblies without the D/L asymmetry of CYS- 300 CYS interactions resulted in chiral structures, but without 301 preferred handedness. The chirality of these terminal 302 assemblies is the result of the confinement of multiparticle 303 assemblies to a cylindrical surface. The pitch and geometry of 304 such structures can be described analytically and strongly 305 depend on the ratio of diameters between the NP and the 306 structure-directing cylinder. 51 In contrast, when we included 307 the chiral interaction between NPs, the assembly of the 308 supraparticular NRs was enantioselective for interaction 309 energies of 10  $k_BT$  (4.1 kJ/mol) and above (Figure 5b,c). 310 This value is noticeably smaller than the estimated 20 kJ/mol 311 energy difference for surface-bound homo- and heterochiral 312 CYS dimers made above. Multiple and cooperative CYS-CYS 313 interactions can occur for face-to-face contacts between 314 tetrahedral NPs, thereby further increasing the energy 315 preference and enantiomeric bias.

Structures with predictable chirality were observed in 317 simulations of particles between 100 and 300 particles and 318 for rod diameters of 1 and 2 NP widths; the effects are 319 particularly pronounced in simulations where the particle 320 number, rod thickness, and box length were commensurate 321 with a close packing of tetrahedra around the rod.

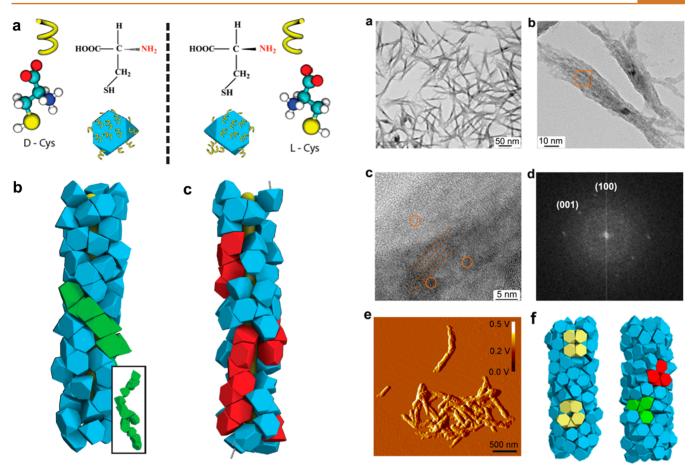


Figure 5. Simulated helical hierarchical assemblies of D- and L-NPs. (a) Schematics of the NP surfaces coated with D- and L-CYS. Chemical structures are related to a "steric coil". Structures produce a predetermined twist based upon the chirality of the stabilizer by biasing local NP motifs. Either left (b) or right (c) handed structures are produced, depending upon the choice of "twist", as determined by a chiral interaction between NP faces. The Te core is shown in yellow, with an outer layer of NPs removed. Insets show the unaltered wire (top) and centers of mass of the NPs connected with bonds along closest neighbors (bottom).

The data described so far show that molecular chirality of 323 324 stabilizers induces the formation of chiral superstructures. To 325 confirm the finding that molecular-scale chirality is reflected in 326 the nanoscale geometry of the assemblies, we also need to test 327 the case of racemic stabilizer-capped NPs. Therefore, we carried out a complementary experimental and computational study of 329 self-assembly processes from rac-NPs taken in the same amount 330 and assembled under the same conditions. The in-silico 331 assemblies driven by an achiral interparticle pair-potential 332 demonstrated (Figure 6f) no distinct handedness, as opposed 333 to the helical assemblies in Figure 5b,c. Concomitant TEM and 334 AFM results of the experimental supraparticle NRs made from 335 rac-NPs show intertwined rod-like assemblies with a diameter 336 of ~40 nm, which is almost double the diameter (Figure 6 and 337 Figure S10a) of helical NRs in Figure 1. The experimental 338 terminal assemblies of rac-NPs are composed of many 339 intertwined thin, ~2 nm nanowires attached with CdTe NPs 340 according to HRTEM image and XEDS analysis (Figure 6c, 341 Figure S11), in which the presence of the hexagonal Te phase 342 can be discerned from the fast Fourier transform pattern 343 (Figure 6d). CD spectra show a weak band around 350-550 344 nm, which may be surprising (Figure S10b). This CD band

Figure 6. Supraparticle NRs resulting from self-assembly of rac-NPs. (a, b) TEM BF images of the NR formed from the rac-NPs. HRTEM (c) and its corresponding FFT pattern (d) characterizations of twinning structures of Te attached with CdTe NPs and corresponding fast Fourier transform pattern of Te nanowires. The fast Fourier transform pattern of the hexagonal Te phase corresponds to (100) (0.38 nm) and (001) (0.6 nm). (e) AFM amplitude image of self-assembled supraparticle NRs from rac-NPs. (f) Simulations of assemblies of tetrahedral rac-NPs (left: front, right: back).

observed for *rac*-NP assemblies is broader than the CD bands 345 obtained for NPs with homochiral surface ligands. At the same 346 time, we see no evidence of helical or other assemblies of NPs 347 with specific chiral preference in TEM and AFM data. The 348 small chiral bias visible in CD spectroscopy for NP mixtures 349 with expected racemic behavior can be attributed to 350 stochastically assembled small NP clusters and should be 351 investigated further.

# CONCLUSIONS

Overall, our findings in this study demonstrate that a small 354 biomimetic unit can drive hierarchical self-assemblies across 355 several scales. The self-assembly of NPs with atomic-scale 356 chirality imparted by the surface ligands results in the formation 357 of helical supraparticles structurally resembling biological 358 prototypes exemplified by the tobacco mosaic virus. The 359 nanoscale twists of helical supraparticles are deterministically 360 controlled by the molecular-scale chirality of the constitutive 361 NPs, which lays the groundwork for understanding the 362 collective behavior of nanoscale biomimetic units. The resulting 363 helical NRs can further self-organize into lamellar liquid 364

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365 crystals, which expand the family of hierarchical assemblies of 366 anisotropic inorganic particles. 27,37 Although possible, chiral 367 order in the lamellar assemblies at the micrometer scale was not 368 observed in this work, and further studies regarding potential 369 chiral hierarchy in these superstructures need to be undertaken. 370 The self-assembly phenomena reported here present multiple 371 venues for further research toward biologically inspired 372 assembly processes and optoelectronic technologies. 52,53

## **373 MATERIALS AND METHODS**

Materials. Cadmium perchlorate hexahydrate was purchased from 374 375 Alfa-Aesar. Aluminum telluride powder was purchased from Materion 376 Advanced Chemicals. L-cysteine (L-CYS) hydrochloride and D-cysteine 377 (D-CYS) hydrochloride were purchased from Sigma-Aldrich. All 378 chemicals are used as received.

Synthesis and Assembly of (D, L, and rac)-cysteine-Stabilized 380 CdTe NPs (D-NPs, L-NPs, and rac-NPs). The molar ratio of CYS to 381 Cd<sup>2+</sup> used in the synthesis of the NPs was 2.2:1, and the synthesis 382 concentration of Cd<sup>2+</sup> was 0.01 mol/L. In order to reduce the 383 possibility of decomposition of cysteine in highly basic solutions, the 384 solution of cadmium perchloride and cysteine was taken at pH 9.6. We 385 briefly dissolved 0.373 g of cadmium perchloride hexahydrate and 386 0.343 g of D-CYS, L-CYS, or rac-CYS in 80 mL of 18 M $\Omega$  pure water, subsequently adjusting the pH value to 9.6. The solution was purged 387 with nitrogen for 1.0 h to remove oxygen. CdTe NPs nucleate after 389 subsequent purging of H<sub>2</sub>Te generated separately by mixing 0.5 M 390 H<sub>2</sub>SO<sub>4</sub> and 0.08 g of Al<sub>2</sub>Te<sub>3</sub>. The CdTe NPs' growth was conducted at 391 90 °C for 4 h and then cooling to room temperature with nitrogen 392 bubbling. The assembly procedure can be summarized as followed: We 393 added a 1:1 volume ratio of 2-propanol to the freshly prepared CdTe 394 solution and centrifuged the mixture solution at 6000 rpm/min for 10 min. After removing the supernatant containing free cadmium complexes and other salts, we immediately dissolved CdTe in pH 396 397 9.0 ultrapure water and bubbled with nitrogen for 5 min. Solutions 398 were kept in 20 mL glass vials sealed with Parafilm at room 399 temperature. It can be found that the solution of nanoparticles changes 400 in color from orange to dark red when left in a dark place at room 401 temperature for 8 h, giving a visual indication that the self-assembly of 402 NPs has occurred.

Characterization. Samples for characterizations were first 403 404 centrifuged at 5000 rpm/min for 5 min, and then the precipitates 405 were redispersed in water in order to remove the random network 406 assemblies in the precursors. In order to get the real-time morphology 407 in the process of self-assembly, the CD spectra and other data about 408 intermediate states were obtained with samples without separation of 409 NP chains and supraparticles NRs. Circular dichroism spectra were 410 collected on a JASCO-815. Fluorescence spectra were carried on a 411 Horiba Fluoro MAX-3. AFM images were taken using the tapping 412 model on a Bruker Multimode-II. TEM images were taken by a JEM-413 3011 transmission electron microscope. Annular dark-field STEM 414 images and XEDS analysis were taken using a JEOL 2010F scanning 415 transmission electron microscope at 200 kV. SEM images were 416 conducted on a FEI Nova NanoLab scanning electron microscope. 417 STEM tomography images were collected on a FEI Titan in annular 418 dark-field mode at 200 kV at the Molecular Foundry, Lawrence 419 Berkeley National Lab. The collection angle was from  $-70^{\circ}$  to  $+70^{\circ}$ . 420 Tomographic alignment of a tilt series and 3D reconstruction were performed using IMOD software. This is achieved by first tracking the 422 small movement of fiducial gold markers with a diameter of 10 nm. 423 Then, the reconstructed volume was filtered using a Gaussian (width, 424 1) for noise reduction in UCSF Chimrea software. SAXS were 425 conducted on a Bruker NanoStar small-angle X-ray scattering (SAXS)

# 427 ASSOCIATED CONTENT

# 428 S Supporting Information

429 The Supporting Information is available free of charge on the 430 ACS Publications website at DOI: 10.1021/acsnano.5b05983.

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### **AUTHOR INFORMATION**

Corresponding Authors	435
*E-mail: sglotzer@umich.edu.	436
*E-mail: kotov@umich.edu.	437
Notes	438
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