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## 40000 TAU DECAYS

Terry S. Mast, Lawrence K. Gershwin, Margaret Alston-Garnjost, Roger O. Bangerter, Angela Barbaro-Galtieri, Joseph J. Murray, Frank T. Solmitz, Robert D. Tripp, and Bryan R. Webber

August 1968

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## UNIVERSITY OF CALIFORNIA

Lawrence Radiation Laboratory Berkeley, California

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## 40 000 TAU DECAYS<sup>†</sup>

## Terry S. Mast, Lawrence K. Gershwin, Margaret Alston-Garnjost, Roger O. Bangerter, Angela Barbaro-Galtieri, Joseph J. Murray,<sup>††</sup> Frank T. Solmitz, Robert D. Tripp, and Bryan R. Webber

Lawrence Radiation Laboratory University of California Berkeley, California

#### August 1968

An exposure of  $1.3 \times 10^6$  pictures in the Berkeley 25-inch hydrogen bubble chamber has yielded about 50 000 K<sup>-</sup>  $\rightarrow \pi^- \pi^- \pi^+$  decays in flight. The K<sup>-</sup> momenta range from 270 to 470 MeV/c.

A preliminary analysis has been made of the pion spectra. Events were accepted if they satisfied the following conditions:

(a) The decay occurred within a restricted fiducial volume.

- (b) The confidence level for the four-constraint kinematic fit was > 0.1%.
- (c) Each pion had a fitted momentum in the laboratory system greater than 30 MeV/c.

These conditions reduced our sample to 42 640 events. Each event was weighted for Coulomb interactions<sup>1</sup> and for the cut on pion momentum. The mean Coulomb weight was 1.034, and the mean weight for lowmomentum pions was 1.036. The distribution of the decay in the K<sup>-</sup> rest frame shows a small dependence on the beam direction and scanning plane. We think this does not significantly affect the results and that remeasurements of failing events will make the distribution more isotropic.

<sup>&</sup>lt;sup>†</sup>Work done under auspices of the U. S. Atomic Energy Commission.

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Systematic effects from uncertainty in the position on the Dalitz plot have not yet been investigated.

The usual variables x and y were defined from the pion kinetic energies in the K<sup>-</sup> rest frame:

x = 
$$\sqrt{3}$$
 | (T <sub>$\pi_1^-$</sub>  - T <sub>$\pi_2^-$</sub> )|/Q,  
y = (3T <sub>$\pi^+ - Q)/Q.$</sub> 

The y projection of the Dalitz plot (Fig. 1) was weighted for phase space and normalized to 1.0 at y = 0.0. A least-squares fit of the projection to the form  $1 + ay + by^2$  gave  $a = 0.239 \pm 0.010$  and  $b = -0.006 \pm 0.021$ , with an 8% confidence level. A fit to the form 1 + aygave  $a = 0.240 \pm 0.010$ , with a 92% confidence level. The importance of the Coulomb corrections is emphasized by a fit made without them which gave  $a = 0.204 \pm 0.010$ , with a confidence level of 20%.

The events for the x projection (Fig. 2) were weighted by phase space and 1/(1+0.24y). The best least-squares fit for this was to a constant. The confidence level for this fit was low (0.005%) due to the low point at x = 0.6. A search for a systematic effect producing this large fluctuation has so far been unsuccessful. The fit failed to improve when an  $x^2$  dependence was introduced.

The dependence on the variables  $\rho$  and  $\theta$ , where  $x = \rho \sin \theta$  and  $y = \rho \cos \theta$ , has been investigated, and the distribution seems to be consistent with the simple  $1 + a\rho \cos \theta$ .

CP invariance predicts the decay spectra of K<sup>-</sup> and K<sup>+</sup> to be identical. The  $|\Delta I| = 1/2$  rule, with a linear approximation to the matrix element, predicts

 $\alpha_{\tau^1}/\alpha_{\tau} = -2$  and  $\alpha_{K_2^0}/\alpha_{\tau} = -2$ ,

where  $a = \alpha M_K Q/m_{\pi}^2$ . The experimental results (Table I) show good agreement between the slopes of the spectra for  $\tau^+$  and for  $\tau^-$ . The two results for  $K_2^0$  decay are not consistent with each other; only one of them (Nefkens et al.<sup>9</sup>) is consistent with the  $\Delta I = 1/2$  prediction. The slopes for  $\tau^{i+}$  are only in fair agreement with the  $|\Delta I| = 1/2$  rule.

Table I. Experimental results (a = $\alpha M_K Q/m_{\pi}^2$ ).				
	Events	Reference	a	α
$ au^-$	42 640	This expt.	$0.240 \pm 0.010$	
	948	Smith et al. (2)	$0.20 \pm 0.065$	
	1 347	Ferro-Luzzi et al.(3	)0.28 ±0.045	
	44 935	average =	$0.241 \pm 0.010$	$0.127 \pm 0.005$
$ au^+$	6752	Grauman et al. (4)	$0.228 \pm 0.030$	
	6786	Plano(5)	0.28 ±0.03	
	3 587	Huetter et al. (6)	$0.21 \pm 0.04$	
	899	McKenna et al. (7)	0.25 ±0.06	
	18024	average =	$0.243 \pm 0.018$	$0.128 \pm 0.009$
к <sub>2</sub> 0	1 350	Hopkins et al. (8)		$-0.40 \pm 0.02$
	1 198	Nefkens et al. (9)		$-0.28 \pm 0.03$
$\tau$ <sup>+</sup>	1874	Bisi et al. (10)		$-0.40 \pm 0.07$
	1 792	Kalmus et al. (11)		$-0.32 \pm 0.03$

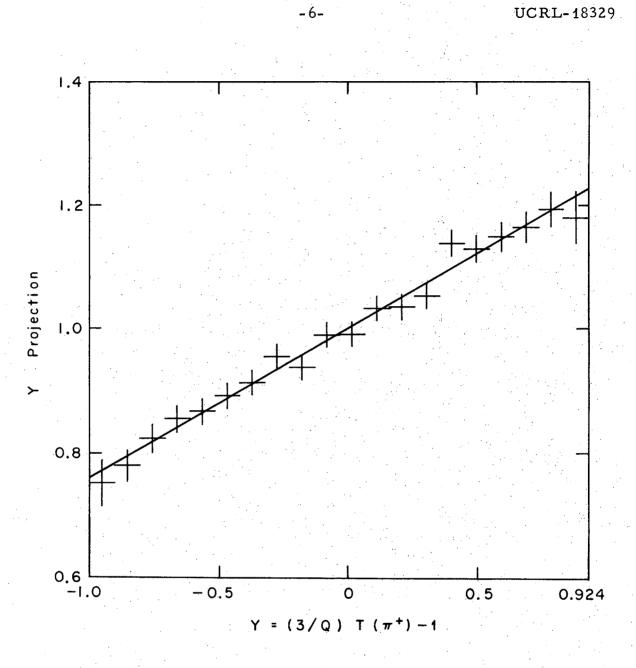
## Footnote and References

- 1. Following Dalitz, Proc. Phys. Soc. <u>69</u>, 527 (1956), it is assumed that the Coulomb interactions between pions i and j multiplies the phase space by a factor  $x/(e^x-1)$ , where  $x = 2\pi e_i e_j/v_{ij}$  and  $v_{ij}$  is the relative velocity between the pions.
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## Figure Captions

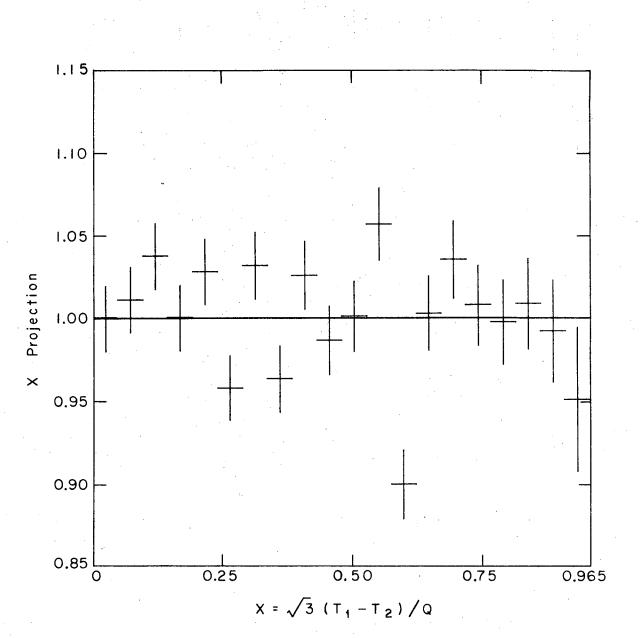
- Fig. 1. Distribution in y weighted for phase space, Coulomb interactions, and a cut on low-momentum pions. The line is a best fit to 1 + ay, where  $a = 0.240 \pm 0.010$ .
- Fig. 2. Distribution in x weighted for phase space, Coulomb interactions, a cut on low-momentum pions, and also weighted by 1/(1+0.24 y).

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Fig. 1



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Fig. 2

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