

1 Rafting-Enabled Recovery Avoids Recrystallization in 3D-

2 Printing-Repaired Single-Crystal Superalloys

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18

## 19 **Abstract**

20       The repair of damaged Ni-based superalloy single-crystal turbine blades has been  
21 a long-standing challenge. Additive manufacturing is a promising route to this end,  
22 but a formidable obstacle is that recrystallization seems inevitable in the dislocation-  
23 riddled heat-affected zone, bringing forth new grains that degrade the high-  
24 temperature creep properties. Here we design a post-3D-printing recovery protocol  
25 that eliminates the driving force for recrystallization, namely the stored energy  
26 associated with the high dislocation content, prior to standard solution treatment and  
27 aging. The post-electron-beam-melting, pre-solution recovery via sub-solvus annealing  
28 is rendered possible by the rafting of  $\gamma'$  particles that facilitates dislocation  
29 rearrangement and annihilation. The rafted microstructure is removed after solution  
30 annealing, leaving behind a damage-free single crystal with uniform  $\gamma'$  precipitates  
31 that is indistinguishable from the the rest of the turbine blade. This discovery offers a  
32 practical means to keep 3D-printed parts from recrystallizing into a polycrystalline  
33 microstructure, paving the way for additive manufacturing to repair, restore and  
34 reshape any superalloy single crystal product.

35

## 36 **Introduction**

37 Modern turbine blades are made of superalloy single crystals (SXs), strengthened  
38 by cuboidal precipitates of ordered  $\gamma'$ -phase in the  $\gamma$ -phase matrix <sup>1</sup>. SX Ni-based  
39 superalloys outperform their polycrystalline counterparts by a large margin, in many  
40 aspects including the resistance to creep, fatigue, and oxidation <sup>2</sup>, by eliminating most  
41 of the defects that form during blade casting. These high-value SX blades are  
42 nonetheless subject to surface damage and cracking upon extensive service in harsh  
43 environments. It is thus critical to find a way to repair the damaged surface while  
44 keeping their single crystalline nature as well as the desired uniform  $\gamma/\gamma'$   
45 microstructure. Such a successful repair will extend the life of turbine blades and  
46 reduce the overall cost significantly.

47 The versatile 3D printing route emerged in recent years appears to be a powerful  
48 option towards this goal <sup>3</sup>. Via “epitaxial” deposition of the alloy, one layer at a time,  
49 additive manufacturing can preserve the crystallographic orientation of the substrate  
50 SX <sup>4-11</sup>. However, as 3D-printing involves fast cooling, the  $\gamma'$  precipitated are  
51 excessively small with uneven sizes and rounded corners, and hence less stable during  
52 high-temperature service. More importantly, a high density of dislocations build up in  
53 the heat-affected zone (HAZ), due to unavoidable (often local) deformation under high  
54 thermal stresses during printing <sup>12,13</sup>. Upon solution treatment at elevated temperatures,  
55 these regions riddled with defects readily undergoes recrystallization (RX) <sup>14</sup> that  
56 renders the microstructure polycrystalline, which would significantly degrade the high  
57 temperature creep performance of blades (Figure S1 in Supporting Information) <sup>2</sup>.

58 These shortcomings are in fact characteristic of all 3D-printed superalloys <sup>15</sup>, such that  
59 the repaired part is no longer as strong as the original SX. There is thus a pressing  
60 need to conceive an innovative treatment that can resolve this problem.

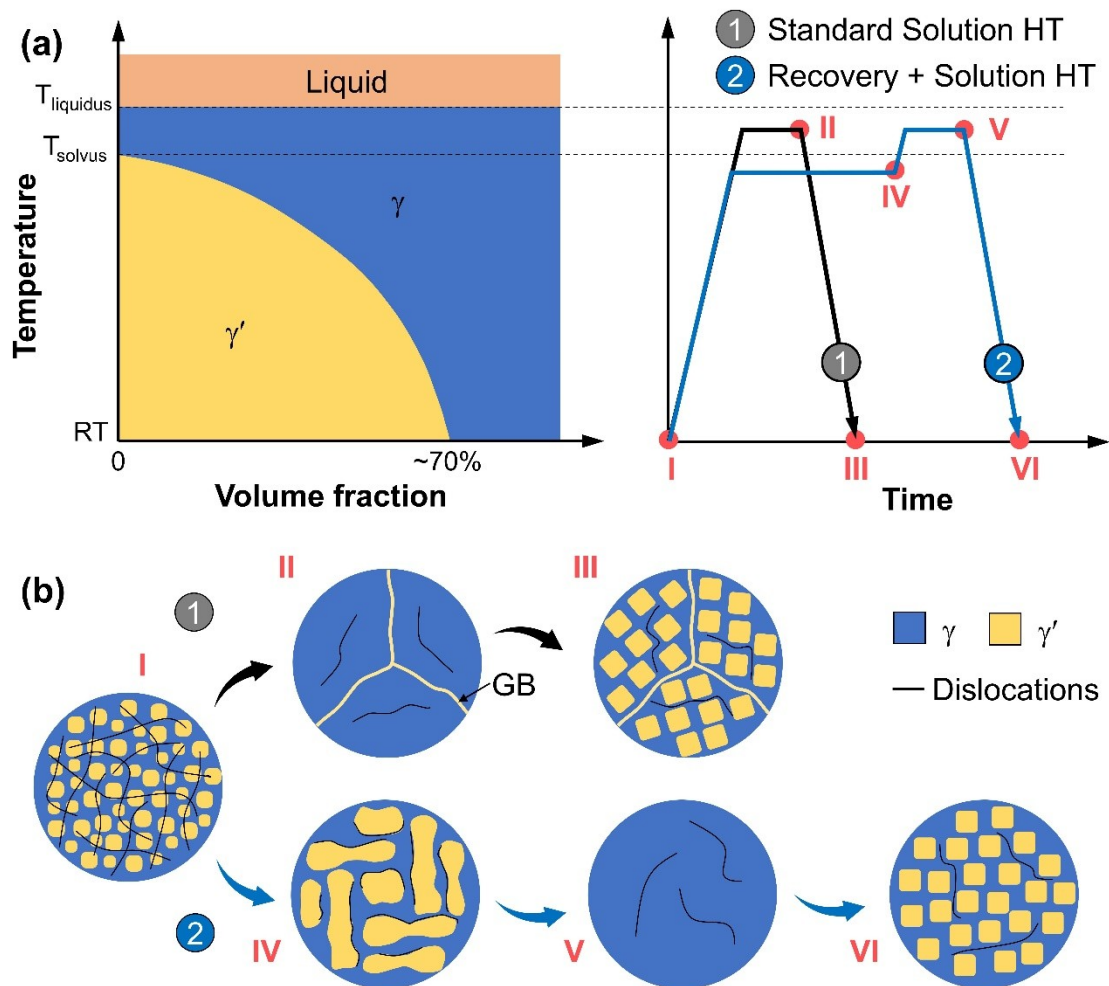
61 To this end, our strategy is to design post-3D-printing annealing to reduce the  
62 driving force for recrystallization that ruins the SX structure. That is, we aim to  
63 remove the accumulated dislocations through a recovery heat treatment (HT) before  
64 the standard solution treatment and aging HT. However, conventional wisdom is that  
65 recovery is difficult to realize in Ni-based superalloys for mainly two reasons. First,  
66 the stacking fault energy of Ni-based superalloys ( $\sim 20$  mJ/m<sup>2</sup>) <sup>16,17</sup> is much lower  
67 than that of pure Ni (125 mJ/m<sup>2</sup>) <sup>18</sup> and Al (166 mJ/m<sup>2</sup>) <sup>19</sup>. A low stacking fault energy  
68 promotes the unit dislocations to dissociate into partial dislocations, hampering their  
69 climb and cross slip, the basic mechanisms for recovery <sup>19,20</sup>. Second, the closely  
70 spaced  $\gamma'$  particles impede the motion of stored dislocations and hence prevent their  
71 annihilation at temperatures below the  $\gamma'$  solvus <sup>19,21,22</sup>. If, instead, the alloy is heated to  
72 temperatures above the  $\gamma'$  solvus, RX sets in quickly well before the stored defects get  
73 effectively removed via recovery <sup>20-23</sup>. Such an RX scenario is demonstrated in Figure  
74 1 (standard solution HT).

75 In this paper, we demonstrate a novel HT process to produce a repaired single  
76 crystal with indistinguishable  $\gamma/\gamma'$  microstructure from the interior. Prior to standard  
77 solution annealing where all the  $\gamma'$  disappear, other defects (with associated excess  
78 energy) accumulated in the HAZ can already be programmed to be relieved by taking  
79 advantage of rafting-facilitated recovery, through an HT step of annealing at an

80 appropriate sub-solvus temperature. During such recovery annealing prior to solution  
 81 HT, rafted microstructure forms in the HAZ, opening pathways to dislocation  
 82 rearrangement and annihilation. This greatly reduces the dislocation density, such that  
 83 in subsequent solution and aging HT, using standard protocol normally used, RX does  
 84 not get triggered to nucleate new grains.

85

86 **Results**

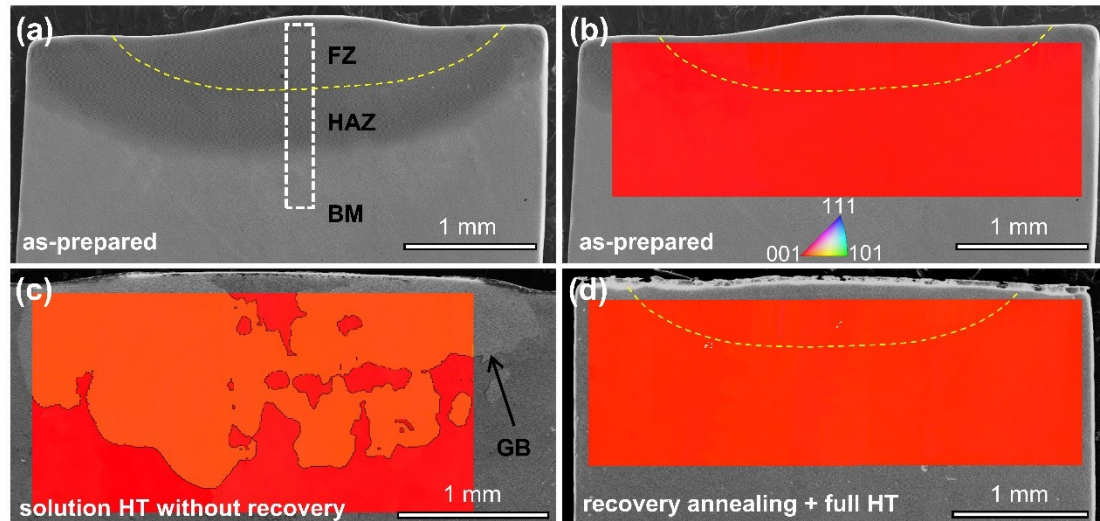


87

88 Figure 1 Microstructural evolution of 3D printed Ni-based single crystals induced by standard solution  
 89 HT, in comparison with our novel HT incorporating recovery annealing. (a) Standard solution HT  
 90 involves one-step annealing between the solvus and liquidus temperatures, while our new HT protocol  
 91 includes a recovery annealing procedure prior to solution HT. (b) RX is triggered by the standard

92 solution HT, but preempted by the new HT.  
93

94 Electron-beam melting (EBM) with no powder feeding was carried out as an  
95 analogue to 3D-printing repairing on AM3, a first-generation SX Ni-based superalloy  
96 used as a model in this study. The electron beam was focused and programmed for line  
97 scans on the substrate (base metal, BM) surface of [001] cast SX boule to generate  
98 fusion zones (FZs). Figure 2 shows the morphology of the EBM sample and the  
99 crystal orientation distribution before and after HT using different protocols. In the  
100 low-magnification image, Figure 2a, three regions can be readily distinguished: un-  
101 affected BM, HAZ, and FZ. The crystal orientation is the same throughout the as-  
102 prepared EBM sample, from the BM to the HAZ to the FZ (Figure 2b). After solution  
103 HT at 1300 °C for 30 min without prior recovery annealing, RX grains and high-angle  
104 grain boundaries are clearly observed in both HAZ and FZ (Figure 2c). In contrast,  
105 after our new 1100 °C recovery annealing for 6 h prior to standard HT, the electron  
106 backscatter diffraction (EBSD) map shows no detectable RX grains (Figure 2d). In  
107 other words, the newly developed HT protocol fulfills the “keeping the single crystal”  
108 requirement. The other requirement that needs to be accomplished simultaneously is  
109 the “uniform  $\gamma'$  microstructure”.



110

111 Figure 2 EBM samples before and after different HT protocols. (a) Three regions are visible in the  
 112 cross-sectional scanning electron microscope (SEM) image of the as-prepared sample: fusion zone  
 113 (FZ), heat-affected zone (HAZ), and unaffected base metal (BM). EBSD orientation maps indicate that  
 114 (b) the as-prepared EBM sample is a single crystal, but (c) RX occurs and high-angle grain boundaries  
 115 are generated upon solution HT at 1300 °C for 0.5 h, while (d) the RX is successfully prevented to keep  
 116 the single crystalline nature by applying a recovery HT at 1100 °C for 6 h before standard HT.  
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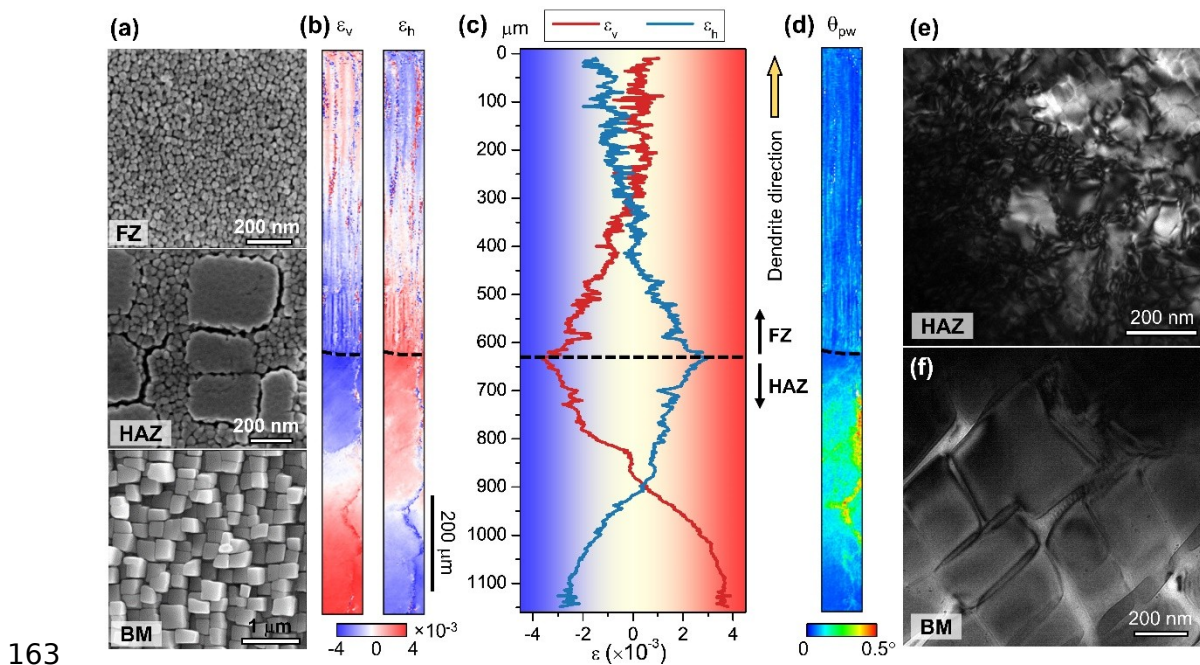
118 The EBM sample (Figure 2b) without HT does not fit the bill in this regard. There  
 119 the fast cooling rate experienced produces not only much finer dendrites in the FZ  
 120 compared to the cast counterpart, but also much smaller  $\gamma'$  precipitates - the average  
 121 size is < 50 nm (Figure 3a), only a quarter of that in the BM. A microstructure  
 122 gradient is observed in the HAZ (in between the FZ and the BM). The upper HAZ  
 123 near the fusion line has experienced a temperature sufficiently high to dissolve all the  
 124 primary  $\gamma'$ . Thus the small particles in the upper HAZ are re-precipitated  $\gamma'$ , quite  
 125 similar to those in the FZ. Moving down away from the fusion line, the peak  
 126 temperature is lower and the period of time that the local temperature is above the  $\gamma'$   
 127 solvus is shorter, such that the primary  $\gamma'$  only dissolves partially. The subsequent re-  
 128 precipitation at lowered temperatures forms fine secondary  $\gamma'$ . Together they constitute

129 a bimodal  $\gamma'$  size distribution. The population of the secondary  $\gamma'$  decreases gradually  
130 with increasing distance, becoming completely undetectable in the region about 300  
131  $\mu\text{m}$  away from the fusion line, leaving only cuboidal primary  $\gamma'$  with side length in the  
132 200 – 300 nm range.

133 After EBM, the crystal is riddled with dislocations and the elastic strain  
134 associated with them is distributed in an inhomogeneous manner. To see the latter, we  
135 have taken synchrotron X-ray microdiffraction ( $\mu\text{XRD}$ ) across the FZ and the HAZ,  
136 covering an area of 80  $\mu\text{m}$  (horizontal)  $\times$  1150  $\mu\text{m}$  (vertical), as indicated using the  
137 boxed rectangle in Figure 2a, at 2  $\mu\text{m}$  spatial resolution. From the orientation inverse  
138 pole figure (IPF) maps (Figure S2 in Supporting Information), it appears that the  
139 entire scanned area remains a single crystal and the dendrite growth is along the [001]  
140 direction. The vertical and horizontal components of the elastic strain tensor, denoted  
141 as  $\epsilon_v$  and  $\epsilon_h$ , respectively, are displayed in Figure 3b. In all regions from the BM  
142 through the HAZ to the FZ, these strains are found to vary considerably in magnitude  
143 from location to location, and even change sign. The BM substrate is under tension in  
144 the vertical direction but is compressed in the horizontal direction; this is believed to  
145 be related to the pre-printing thermal history of the superalloy. The opposite is  
146 observed for strains in the HAZ:  $\epsilon_h$  is tensile while  $\epsilon_v$  is compressive. The transition  
147 from the BM to the HAZ is smooth and gradual. Both  $\epsilon_h$  and  $\epsilon_v$  change sign near the  
148 BM/HAZ interface.  $\epsilon_h$  and  $\epsilon_v$  reach their peak magnitudes at the HAZ/FZ interface,  
149 and then decrease together upon entering the FZ. The magnitude of the strain is highly  
150 variable in the FZ, and changes sign from the interdendritic regions to dendrite cores,



151 probably due to non-uniform chemical distribution. A high density of dislocations in  
 152 the HAZ is reflected by the obviously-broadened Laue peak width, see the colored  
 153 peak width map in Figure 3d (although the dislocation density is difficult to quantify  
 154 accurately from the peak width). In the FZ, the peak width map also exhibits fine  
 155 stripe features, presumably attributable to the chemical and microstructural  
 156 heterogeneities between dendrite cores and interdendritic regions. This evidence of a  
 157 large population of dislocations in many local regions is consistent with the  
 158 observation in the transmission electron microscope (TEM) image of Figure 3e, in  
 159 which the dislocation density is measured to be about  $8 \times 10^{14} \text{ m}^{-2}$ . This value is for a  
 160 small local region, but from the colored peak width map the dislocation density is  
 161 inhomogeneous and should be higher in many other regions of the HAZ. Similar to  
 162 previous reports for cast materials <sup>24</sup>, the BM is almost defect free (Figure 3f).



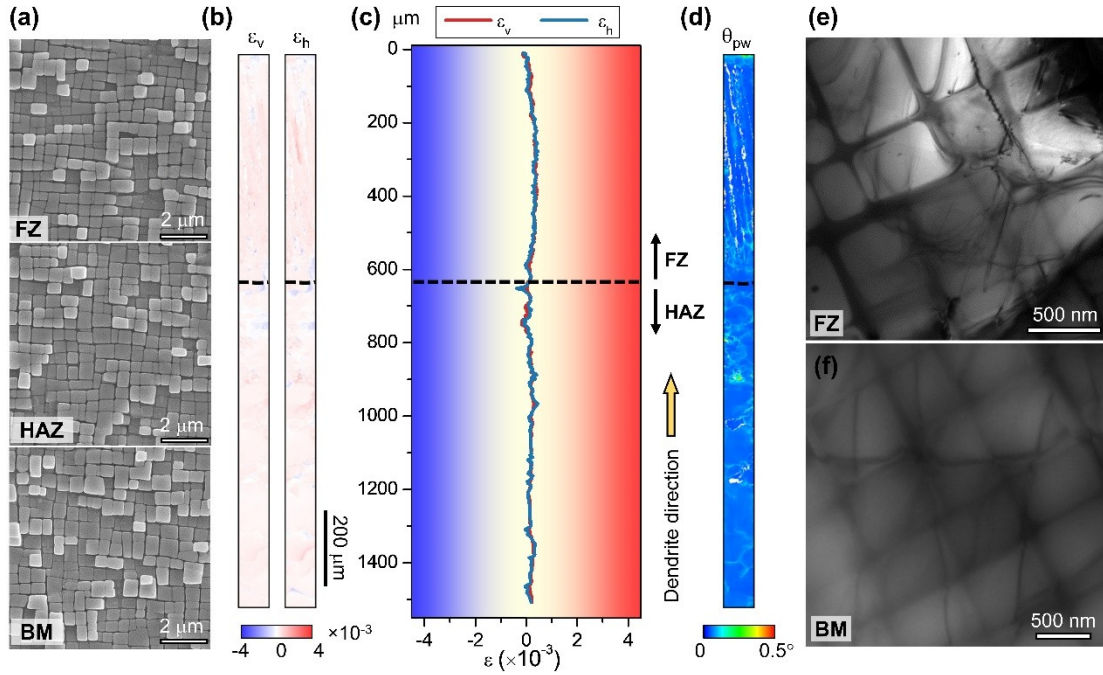
164 Figure 3 Inhomogeneous distributions of  $\gamma'$  morphology, elastic strain and dislocations in the EBM  
 165 sample. (a) Distinctly different  $\gamma'$  morphologies in the FZ, HAZ, and BM, as displayed in their  
 166 corresponding SEM images. (b) 2D and (c) 1D lattice elastic strain as well as (d) diffraction peak width

167 maps indicate non-uniform deformation. The yellow-red region in (d) is believed to represent a  
168 dislocation density of  $> 8 \times 10^{14} \text{ m}^{-2}$ . The fusion line is marked using a dashed line in (b-d). TEM  
169 images indicate that dislocations of high density are present in the HAZ (e), while the BM is almost  
170 defect free (f).  
171

172 We stress again that the dislocations and elastic energy stored are the root cause  
173 that drives RX upon the ensuing solutionizing treatment. After the EBM specimen  
174 was solution treated at 1300 °C for 30 min, EBSD found RX grains in not only the  
175 HAZ but also the FZ (Figure 2b). Shortening of the solution treatment time reduced  
176 the area fraction of the RX grains, but was not able to avoid RX, unless the solution  
177 treatment is conducted at a temperature well below that needed to reach complete solid  
178 solution. Such a temperature, however, would not be able to solutionize the  
179 interdendritic region, where the elements promoting  $\gamma'$  formation are known to be  
180 enriched and thus the  $\gamma'$  solvus temperature is higher than that in the dendrite cores. In  
181 other words, without the proposed pre-solution recovery annealing, either RX occurs  
182 (at super-solvus temperatures) or the solution treatment is ineffective (at sub-solvus  
183 temperatures) to remove the chemical and microstructural heterogeneities from  
184 dendrites.

185 After incorporating our recovery annealing before solution HT, the microstructure  
186 becomes completely different from not only the reference sample (full HT after EBM  
187 without recovery annealing; as seen in Figure 2c, there RX is obvious, although the  
188 precipitates inside the grains may become uniform in size), but also the EBM sample  
189 (non-uniform microstructure, as discussed above with Figure 3). The  $\gamma'$  precipitates in  
190 the FZ, HAZ, and BM are uniform, exhibiting identical morphology and size. As seen

191 in Figure 4a, they all have cubical shape with sharp vertices, straight edges, and  
192 uniform side length of about 500 nm. The  $\gamma'$  precipitates grow during the aging HT,  
193 becoming larger than those in the BM of the EBM sample. From the micro-  
194 indentation test results shown in Figure S3, after HT with recovery the hardness is also  
195 uniform throughout, all the way from the FZ to the HAZ and to the BM, in stark  
196 contrast to the non-uniform hardness distribution in the EBM sample due to the  
197 pronounced spatial variation of the  $\gamma'$  precipitates size and morphology (Figure 3a) as  
198 well as of the dislocation density.  $\mu$ XRD results prove that the residual strain is fully  
199 released, evidenced by the uniform light color in the 2D maps, Figure 4b. There is  
200 only slight variation in the 1D strain profile of Figure 4c. The dislocation density in  
201 the HAZ and the FZ is also brought down significantly. The TEM image in Figure 4e  
202 shows a dislocation density of  $2 \times 10^{13} \text{ m}^{-2}$ . Multiple TEM images are taken from the  
203 HAZ, two of which are displayed in Figure S4. From these different regions the  
204 average dislocation density in the HAZ is found to be about  $3 \times 10^{13} \text{ m}^{-2}$ , and in all  
205 local regions it never exceeds  $5 \times 10^{13} \text{ m}^{-2}$ . In other words, compared to the as-  
206 prepared EBM sample, the dislocation density decreases by more than 20 times in the  
207 HAZ. The BM, similar to the as-prepared sample, stays dislocation free (Figure 4f).



208

209 Figure 4 Uniform distributions of  $\gamma'$  morphology, elastic strain and dislocations after full HT following  
 210 our new protocol incorporating recovery annealing. (a) Indistinguishable SEM images of  $\gamma'$  precipitates  
 211 in the FZ, HAZ, and BM.  $\mu$ XRD proves (b, c) completely relieved strain, and (d) uniform and sharp  
 212 diffraction peaks indicating low defect density (except at dendrite boundaries). The fusion line is  
 213 marked with a dashed line in (b-d). TEM images indicate significant reduction in dislocation density in  
 214 the HAZ (e), while BM remains almost defect free (f).  
 215

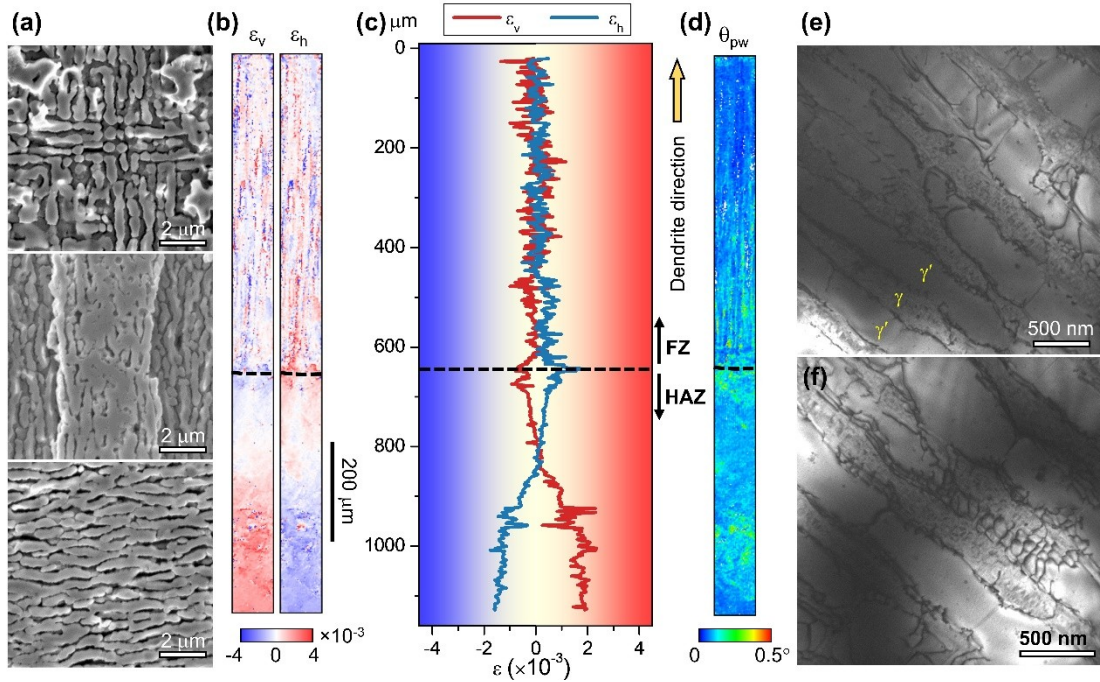
216 All the evidences above, including the  $\gamma'$  shape, size, and micro-hardness  
 217 distributions, as well as the reduced strain and dislocation density, demonstrate that  
 218 our new protocol is effective in producing a uniform microstructure. The key step,  
 219 absent for the standard HT, is the sub-solvus annealing at 1100 °C for 6 h before  
 220 solution treatment. It lowers the defect density that would have otherwise triggered  
 221 RX during solution HT. Our discussion in the following will therefore focus on the  
 222 mechanism involved to accomplish dislocation removal during this important recovery  
 223 annealing step.

224 As shown in Figure 5, recovery annealing at 1100 °C for 6 h leads to rafting

225 almost everywhere in the FZ and the HAZ. Rafting is a common occurrence well  
226 known in superalloys subjected to external stresses and elevated temperatures. It can  
227 also be induced by residual stresses at certain temperatures <sup>25</sup>. In our case the rafted  
228 microstructure is caused by the residual stresses built-in during the EBM. The  
229 recovery annealing is carried out at a temperature close to the typical creep testing  
230 temperature for this superalloy <sup>26</sup>. Three typical rafted  $\gamma'$  microstructures  
231 (demonstrated in Figure 5a), which are recorded in the upper FZ, near the FZ/HAZ  
232 interface, and in the lower HAZ, respectively, are examined in detail. For AM3, a  
233 superalloy with negative lattice mismatch, the rafting direction is expected to be  
234 perpendicular to the tensile stress direction. In the upper FZ region, both horizontally  
235 and vertically rafted structures are observed, agreeing well with the observed spatial  
236 variation of the strain tensor and the sign (stress direction) displayed in Figure 3. Near  
237 the interface between the FZ and HAZ, vertically rafted structure forms. This is also  
238 consistent with the measured horizontal tensile strain in this region. As the strain  
239 direction reverses sign in the lower HAZ, the rafting direction changes as well. In  
240 confirmation, the strains in the recovery-annealed specimen are significantly reduced,  
241 suggesting that the elastic strain energy stored in the EBM sample is effectively  
242 released along with rafting when the  $\gamma/\gamma'$  phase boundaries migrate (Figure 5b, c). The  
243 peak width map indicates that there are residual defects in the HAZ and the lower FZ  
244 (Figure 5d), thus TEM is employed for direct observation. As shown in Figure 5e and  
245 5f, dislocations now line up at  $\gamma/\gamma'$  phase boundaries, and dislocation networks are  
246 also observed occasionally. In the literature, these two types of dislocation

247 configurations were reported in superalloys after creep testing <sup>27,28</sup>. The dislocation  
248 density near these interfaces is measured to be  $9 \times 10^{13} \text{ m}^{-2}$  and  $2 \times 10^{14} \text{ m}^{-2}$ ,  
249 respectively, but still up to one order of magnitude lower than that in the HAZ of the  
250 EBM sample. Taken together, the  $\mu$ XRD and TEM observations suggest that the  
251 combined rafting-recovery is akin to that due to creep <sup>25,29,30</sup>. Note here again that  
252 previous annealing efforts were only able to achieve limited recovery, especially for  
253 regions that experienced rather high pre-strains in plastic deformation <sup>31,32</sup>. Most of the  
254 previous attempts apparently have missed the appropriate temperature window, as they  
255 were probably unaware of the potential role that could be played by the rafting of  $\gamma'$   
256 particles.

257       Regarding rafting, it is in fact the migration of the  $\gamma/\gamma'$  phase boundaries that has  
258 ushered in a new mechanism to facilitate dislocation recovery. First, many dislocations  
259 sink into the interphase boundaries as they sweep by. Second, more spaces are opened  
260 for dislocation motion, as rafting widens some  $\gamma$  channels to accommodate  
261 dislocation movement and interactions that lead to annihilation. Third, the residual  
262 dislocations rearrange into lower-energy configurations at the  $\gamma/\gamma'$  interfaces (Figure  
263 5) and these aggregated dislocation bundles re-configure more readily upon  
264 subsequent solution treatment at higher temperatures, as they no longer need to run  
265 across  $\gamma'$  precipitates to annihilate. This reduction of concentrated defects storing high  
266 energy leaves few spots as RX nucleation sites, such that the ensuing solution  
267 treatment no longer sets off RX.



268

269 Figure 5  $\gamma'$  morphology, elastic strain and dislocations after 1100 °C recovery annealing. (a) Rafted  $\gamma'$   
 270 observed in various directions.  $\mu$ XRD shows lowered strain in (b) 2D and (c) 1D profiles, as well as (d)  
 271 reduced diffraction peak width. The fusion line is marked by dashed line in (b-d). TEM observation  
 272 illustrates dislocations rearranged (e) at  $\gamma/\gamma'$  interfaces and (f) as networks. These are presumably low-  
 273 energy dislocation configurations that can readily annihilate upon solution treatment, and the  
 274 dislocation densities in these two images are almost one order of magnitude lower than that in the HAZ  
 275 of the EBM sample.  
 276

## 277 Discussion

278 The recovery HT protocol reported in this study is applicable to other 3D printing  
 279 scenarios. To demonstrate this, epitaxial AM3 layers were also manufactured using the  
 280 direct laser forming method on top of the cast BM with the same dimensions as those  
 281 for the EBM case. The subsequent heat treatments are the same, incorporating the  
 282 same recovery annealing. Homogeneous  $\gamma'$  precipitates are again obtained in all the  
 283 FZ, HAZ, and BM regions, without any sign of RX, as shown with EBSD  
 284 examination results in Figure S5 in Supporting Information.

285        Furthermore, although the recovery annealing costs 6 h, it is only a minor  
286 extension to the standard HT, which involves solution treatment at 1300 °C for 3 h  
287 followed by two steps of ageing at 1100 °C and 870 °C for 6 h and 20 h, respectively.  
288 The inserted recovery step is thus only a simple addition, with little additional cost  
289 while effectively curtailing the undesirable RX.

290        Finally, we mention in passing a tweak to the recovery approach to lower the  
291 defect density and hence the RX driving force: the substrate can be heated *in situ*  
292 along with the 3D printing (using electron beam, laser or other heating resource)  
293 process. This “preheating” may lower the temperature for recovery annealing and  
294 shorten its duration, depending on the preheating conditions. However, this substrate  
295 preheating is expected to demand a stringent temperature control: a substrate  
296 temperature too high would diminish the temperature gradient, jeopardizing the  
297 epitaxial deposition; A temperature too low, on the other hand, would still incur too  
298 much deformation in local regions and hence too many stored dislocations in the  
299 HAZ. As such one would still need to add rafting-assisted recovery to suppress RX  
300 altogether. Our results in this paper demonstrate that the delicate control of preheating  
301 is not necessary. As a limiting case, simply using post-3D-printing recovery annealing  
302 alone, with no preheating at all, is already adequate to get rid of the normally expected  
303 RX. The post-electron-beam-melting, pre-solution recovery parameters (annealing  
304 temperature and time) can be adjusted depending on the plastic strains that need to be  
305 recovered inside the printed/repared alloy; a systematic documentation would  
306 however exceed the space limit of this Letter.



307

## 308 **Conclusions**

309 In summary, we have designed a new heat treatment protocol to satisfy the  
310 requirement of “no RX together with uniform  $\gamma$ ”, mandated for 3D-printing repair of  
311 Ni-based superalloy single crystals. Most essential in our strategy is the realization of  
312 sub-solvus recovery prior to solution treatment, eliminating most of the stored defect  
313 energy that drives nucleation of new crystals during solution treatment. This is made  
314 possible by dislocation rearrangement and annihilation that are otherwise inactive in  
315 the absence of rafting  $\gamma'$ . Previous 3D-printing work to produce single crystalline  
316 superalloys was not aware of, and did not take advantage of, this vehicle that enables  
317 substantial recovery. Our finding thus opens an avenue to make additive manufacturing  
318 a widely applicable tool when dealing with the manufacture and repair of single-  
319 crystal superalloy part. We envision this could be used for the welding of single-  
320 crystal parts as well.

321

## 322 **Methods**

323 EBM with no powder feeding was carried out using a DMAMS Zcomplex3™  
324 electron-beam 3D-printing system operated in  $10^{-3}$  mbar vacuum. The substrate (BM)  
325 was cut into a cylinder 13 mm in diameter and 4 mm in height, from [001] cast SX  
326 boules after solid solution heat treatment. Electron beam of 15 mA was accelerated to  
327 60 keV and focused onto the BM surface to form a melt pool. Line scanning was  
328 programmed at the velocity of 10 to 15 mm/s to ensure epitaxial dendrite growth in

329 the melt pool. A fusion zone (FZ) of about 1500  $\mu\text{m}$  in width and 800  $\mu\text{m}$  in depth  
330 was generated. The EBM sample was then recovery-annealed at 1100  $^{\circ}\text{C}$  for 6 h,  
331 solution-treated at 1300  $^{\circ}\text{C}$  for 0.5 h, and then aged at 1100  $^{\circ}\text{C}$  and 870  $^{\circ}\text{C}$  for 6 h and  
332 20 h, respectively. Note that although the solution treatment temperature was the same  
333 as the standard HT protocol, the duration needed was significantly shorter, because  
334 solute segregation in the EBM sample is much less than in its cast counterpart.  
335 Comparisons are made with an identical EBM sample heat treated without recovery,  
336 skipping the 1100  $^{\circ}\text{C}$  annealing step. This is the standard HT sample serving as the  
337 reference.

338       Micro-hardness test was carried out using a Vickers hardness indenter under the  
339 force control mode, after EBM as well as after recovery annealing. On each specimen,  
340 a matrix of indentations covered the area from the BM to the FZ. Each indent was at  
341 least 4  $\mu\text{m}$  deep to exclude the surface effect, and neighboring indents were 105  $\mu\text{m}$   
342 apart to make sure the hardness value was not influenced by the plastic deformation  
343 around adjacent indents.

344       Before and after recovery-HT, the microstructure was examined under secondary  
345 electron mode in a SEM after etching in 25% phosphoric acid water solution at the  
346 voltage of 5 V for 10 s. RX was monitored by mapping the crystal orientation of the  
347 sample surface using EBSD, after electrochemical polishing in 10% perchloric acid  
348 alcohol solution at the voltage of 20~30 V for about 60 s.  $\mu\text{XRD}$  sample was electro-  
349 polished the same way, and then scanned using micro-focused synchrotron  
350 polychromatic X-ray beam at the Advanced Light Source of Lawrence Berkeley

351 National Laboratory <sup>33</sup>. The collected Laue diffraction data were processed using a  
352 custom-developed software based on the peak position comparison method to measure  
353 the strain tensor distribution accurately <sup>34</sup>. The defect density maps were obtained  
354 semi-quantitatively by plotting the Laue peak width distribution <sup>35</sup>. TEM specimens  
355 were prepared using the conventional twin-jet electropolishing.

356

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### 367 **Author contributions**

368 K.C., E.M., Z.W.S. and J.L. designed the project. R.H. conducted the EBM and direct  
369 laser forming experiments. R.H. and Y.L. developed the HT protocol for the EBM  
370 specimens. S.L. performed the HT on the direct laser forming sample. W.Z. carried  
371 out TEM and  $\mu$ XRD characterizations, and then analyzed and interpreted the  $\mu$ XRD

372 data with R.H. under the supervision of K.C. and N.T. The paper was written by K.C.,  
373 R.H., E.M., Z.W.S. and J.L. All authors contributed to discussions of the results.

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