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### Title

A MEASUREMENT OF THE POSITIVE  $\pi^-$   $\mu$  DECAY LIFETIME

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O. Chamberlain, R.F. Mozley, J. Steinberger, and C. Wiegand

May 10, 1950

Berkeley, California

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-3-

A MEASUREMENT OF THE POSITIVE  $\pi^+$  DECAY LIFETIME

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May 10, 1950

The lifetime for the decay of a  $\pi$  meson into a  $K$  meson and neutral particle was first measured by Richardson<sup>1</sup> and later by Martinelli and Panofsky<sup>2</sup>. The method was the same in both cases: The fraction of  $\pi$  mesons surviving various times of flight is measured by placing photographic detectors at various path lengths from the target.

In the experiment reported here we observe the time lag between the two bursts of fluorescence due to mesons decaying in a scintillation crystal. The first burst is due to the stopping of the entering  $\pi$  meson, the second to the  $K$ -meson. As is shown in Fig. 1, a particle penetrating the first and into the second crystal starts the sweep ( $10^{-8}$  sec/mm) of an oscilloscope. The pulses in the second crystal are delayed  $0.5 \times 10^{-6}$  sec. to allow the sweep to start and brighten and are then photographed. If the responsible particle is a  $\pi^+$  meson which stops in the crystal, it undergoes  $\pi^+ \rightarrow K^+$  decay and two pulses appear on the trace. The  $K^+$  meson has a range of only 2 mm in the crystal. If its decay electron is detected some time ( $.5-2.5 \times 10^{-6}$  sec) later; a neon light flashes and is photographed together with the scope trace. Only such marked traces are measured. Of these marked traces, 650 or roughly one-half, show the two pulses of the  $\pi^+ \rightarrow K^+$  event. Five percent are calculated to be due to random delayed coincidences, and another 3 percent

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<sup>1</sup>J. Richardson, Phys. Rev. 74, 1720 (1948).

<sup>2</sup>E.A. Martinelli and W.K.H. Panofsky, Phys. Rev. 77, 465 (1950).

-4-

due to  $\pi$  mesons which have decayed in flight and come to rest in the second crystal as  $\pi$  mesons. The remaining traces are due to  $\pi$  decays which are too fast to be resolved. The sweep speed of the oscilloscope is calibrated periodically with an oscillator of known frequency.

In Fig. 2 the data are presented both in differential and integral form. The integral data have the advantage of greater statistical accuracy, since they make use of all starred traces after the background subtraction. They are, however, more open to systematic error.

The mean life is (standard deviation):

$$\pi = 2.65 \pm .12 \times 10^{-8} \text{ sec. (differential data)}$$

$$\pi = 2.65 \pm .08 \times 10^{-8} \text{ sec. (integral data)}$$

The previously reported values are:

$$\pi = 1.11 \pm_{-.35}^{+.45} \times 10^{-8} \text{ sec. (Richardson}^1)$$

$$\pi = 1.97 \pm_{-.25}^{+.21} \times 10^{-8} \text{ sec. (Panofsky and Martinelli)}$$

The reasons for the discrepancies are not understood.

We wish to thank Professors E. McMillan and E. Segrè for their encouragement. This work was supported by the Atomic Energy Commission.

Figure Captions

Fig. 1 Block Diagram of Apparatus.

Fig. 2 Decay rate of the positive  $\pi$  meson. The zero-time point on the integral curve represents all marked traces minus accidental traces. The zero-time point on the differential curve represents those traces with pulse separation between  $2.17$  and  $3.97 \times 10^{-8}$  sec, the next point those between  $3.98$  and  $5.78 \times 10^{-8}$  sec, etc.



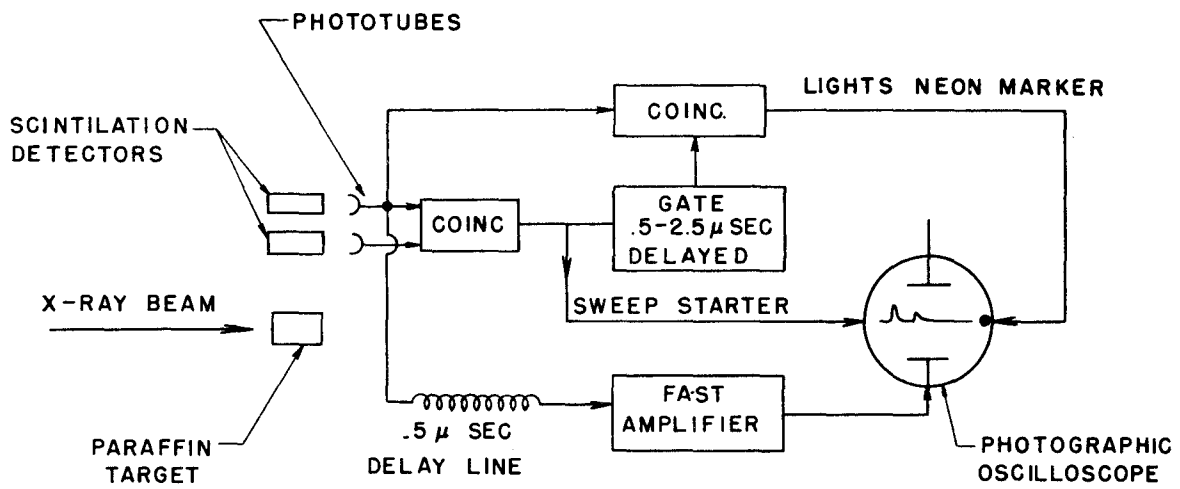


FIG. 1

MU 232

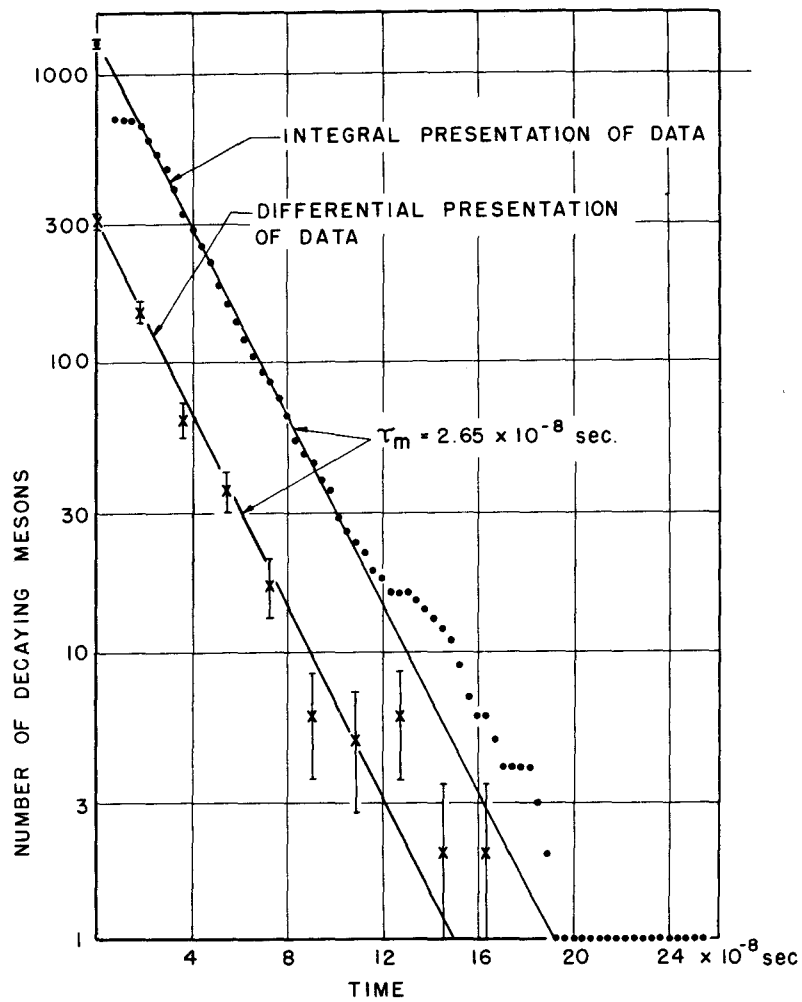


FIG. 2

MU 233