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## Title

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THE $\Sigma / \Lambda$ BRANCHING RATIO OF $\mathrm{Y}_{1}^{*}$
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April 25, 1961

# THE $\Sigma / \Lambda$ BRANCHING RATIO OF $Y_{1}^{*}+$ 

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April 25. 1961

Recently a $T=1$ resonance in the $\Lambda \pi$ aystem called $X_{1}^{*}$ has been observed with a mass of $1385 \mathrm{Mev} .^{1-6}$ Two types of resonances have been predicted that might relate thi observation to other elementary-particle interactions:(1) $P 3 / 2$ resonances in the $\Lambda \pi$ and $\mathcal{E n}$ systems predicted by global symmotry ${ }^{7,8}$ corresponding to the ( $3 / 2,3 / 2$ ) resonance of the $\pi N$ ystem, (2) a spin-1/2 Y- $\pi$ resonance resulting from a bound state in the RNsystem. ${ }^{9,10}$ The position and width of the observed $Y_{1}^{*}$ resonance agree with both theories but since the spin and parity have not yet been determined, it is impossible at present to distinguish between the two theoretical interpretations.

Global symmetry ${ }^{11}$ predicts a theoretical branching ratio $\left(\mathrm{Y}_{1}^{*+} \rightarrow \Sigma^{0}+\pi^{+}\right) /\left(\mathrm{Y}_{1}^{*+} \rightarrow \Lambda+\nabla^{+}\right)=1 / 4$ for the $\mathrm{T}=1$ resonance. The phase-space factor $\left(P_{\Sigma} / P_{\Lambda}\right)^{3}=(126 / 207)^{3}=0.225$ reduces the expected branching ratiofor this process to $R=(1 / 4) \times 0.225 \sim 5 \%$. Furthermore, as a consequence of charge independence the rates $\mathrm{Y}_{1}^{*+} \rightarrow \Sigma^{ \pm}+\mathrm{m}^{0}, \mathrm{Y}_{1}^{* \pm} \rightarrow \Sigma^{0}+\pi^{\boldsymbol{*}}$. and $Y_{1}^{* 0} \rightarrow \Sigma^{ \pm}+\pi^{F}$ are equal. In addition to the $T=1$ resonance, a $T=2 \Sigma-\pi$ resonance with a total energy of 1540 Mev and a half width, $\mathrm{r} / 2$, of 60 Mev is predicted by global symmetry. ${ }^{8}$

[^0]The $\bar{K}-N$ bound-state model suggests values of $R$ considerably larger than $5 \%$. However, when non-zero offective fangoe atcotakdn into account ${ }^{12}$ R can become quite small, especially if the ( 2 A ) parity should be odd.

To investigate these possibilities, we have continued our study of $\mathbb{K}^{-}-\mathrm{p}$ interactions at $1.15 \mathrm{Bev} / \mathrm{c}$ in the Lawrence Radiation Laboratory 15 -in. hydrogen bubble chamber by atudying events in which $a \Sigma$ is observed. The total cross sections for these interactions are shown in Table $I$; only statistical errors are indicated. The separation of $\Sigma^{ \pm}+\pi^{\overline{+}}+\pi^{0}$ and $\Sigma^{ \pm}+\pi^{\bar{q}}+2 \pi^{0}$ events was difficult because many of the latter events can also be fitted to the first hypothesig. The numbers given in Table 1 and in the Dalitz and mass plots below were corrected to account for this ambiguity. The correction factor was estimated by using our $\Sigma^{ \pm}+\pi^{\mp}+\pi^{+}+\pi^{-}$events.

Dalitz plots for the thresbody reactions are shown in Fig. 1. The $\mathrm{Y}_{1}^{*}$ resonance of mass 1385 Mev should appear as a bunching of events about both horizontal and vertical lines correaponding to $T \pi=282 \mathrm{Mev}$. To obtain an upper limit for the branching ratio $R$, we combined the events into different charge atates of the $\Sigma \pi$ system. All charged $\Sigma$ were observed; however, in the $\Sigma^{0}$ cases only two-thirds of the events were observable because of the neutral decays of the $\Lambda^{0}$. Furthermore, we had estimated that about onethird of the $\Sigma^{0} \pi^{+}+\pi^{-}$evente alsofitted a $\Lambda \pi^{+} \pi^{-}$interpretation and had been included in already published data. ${ }^{1}$ Consequently each $\Sigma^{0} \pi^{+} \pi^{-}$ovent was: given a weight of 2.25. The resultant mass spectra are shown in Fig. 2. In the cases of $(\Sigma \pi)^{+}$and $(\Sigma \pi)^{-}$there appears to be no excess of events in the region of $M_{1}=1385 \mathrm{Mev}$. Using the number of $\left(\Lambda^{+}{ }^{+}\right.$) and $\left(\Lambda \pi^{-}\right)$events with $1355 \mathrm{Mev}<\mathrm{M}_{\mathrm{A} \pi}<1415 \mathrm{Mev}$ from reference l, and assuming that all $\Sigma \pi$ events in the game regions of Fig. 2 are $Y_{i}^{*}$, we obtain $R_{m a x} \leqslant 8 \%$.

This treatment yields an unrealistic upper limit, since there ie no evidence of any peaking above background. The results are consistent with $R=0$. The $\Sigma^{ \pm}+\pi^{+}+2 \pi^{0}$ evente possibly misidentified as $\Sigma^{ \pm}+\nabla^{\mp}+\pi^{0}$ (or vice versa) do not fall into the mass band used in this analysis, since they yield apparently high masses of the $\Sigma^{*}{ }_{\square}{ }^{0}$ syatem.

We conclude that the $\Sigma / \Lambda$ branching ratio $R$ for the strong decay of the $T=1 Y_{1}^{*}$ is at most a few percent and is consistent with zero. This result agrees with the value of $R$ obtained by Berge. ${ }^{3}$ As indicated above this value of $R$ does not rule out either the global symmetry or the KN bound-state model of the $Y_{1}^{*}$ resonance. No ovidence for the resonance with $T=2$ predicted by global symmetry at $M=1540 \mathrm{Mev}$ is observed; however, this wide resonance would be hard to separate from background.

The authors wish to thank the many members of the Bevatron and 15 -in. bubble chamber crew and the scanners who made this experiment possible. One of us, Philippe Eberhard, wishes to thank the Philippe Foundation, Inc. and the Commisariat $L$ ' Energie Atomique for a fellowahip.

Table I. Cross sections for the $\Sigma$ producing interactions at $1.15 \mathrm{Bev} / \mathrm{c}$

Cross sections
(mb)

$$
\begin{array}{llr}
\text { Reaction } & \begin{array}{c}
\text { No. of events } \\
\text { (uncorrected) }
\end{array} & \begin{array}{c}
\text { Cross sectiol } \\
\text { (mb) }
\end{array} \\
\hline \mathrm{K}^{-}+\mathrm{p} \rightarrow \Sigma^{-}+\pi^{+} & 87 & 1.40 \pm 0.16 \\
\rightarrow \Sigma^{+}+\pi^{-} & 84 & 1.34 \pm 0.18 \\
& \rightarrow \Sigma^{+}+\pi^{-}+\pi^{0} & 57 \\
& \rightarrow \Sigma^{-}+\pi^{+}+\pi^{0} & 54 \\
& \rightarrow \Sigma^{0}+\pi^{+}+\pi^{-} & 0.97 \pm 0.16 \\
& \rightarrow \Sigma^{+}+\pi^{-}+\pi^{0}+\pi^{0} & 0.83 \pm 0.20 \\
& \rightarrow \Sigma^{-}+\pi^{+}+\pi^{0}+\pi^{0} & 13
\end{array}
$$

## FOOTNOTES

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## FIGURE LEGENDS

Fig. 1. Dalitz plots for the reactions:
(a) $K^{-}+p=\mathbf{\Sigma}^{+}+{\mathbf{a}^{-}}^{+}+\boldsymbol{n}^{0}$ (57 vents)
(b) $\mathrm{K}^{*}+\mathrm{p} \rightarrow \mathrm{\Sigma}^{-}+\boldsymbol{\pi}^{+}+\mathbf{D}^{0}$ (54 evente)
(c) $\mathrm{K}^{-}+\mathrm{p} \rightarrow \Sigma^{0}+\pi^{+}+\pi^{-}$ (27 events).

Fig. 2. Mase plots of the charged and neutral $\mathbb{\Sigma}-\boldsymbol{\pi}$ eytems, including curves vepreaenting phase-space distributions.
(a) Mass of $\left(\Sigma_{\bar{w}}\right)^{-}$, from the reactions: $\mathrm{K}^{-}+\mathrm{p} \rightarrow \mathrm{x}^{0}+\mathbf{E}^{-}+{ }^{+}$

$$
\rightarrow \Sigma^{-}+\pi^{0}+\pi^{+}
$$



$$
\rightarrow \Sigma^{+}+\pi^{0}+\pi^{-}
$$

(c) Masa of $(\Sigma \pi)^{0}$, from the reactions: $\mathrm{K}^{-}+\mathrm{p} \rightarrow \mathrm{\Sigma}^{+}+\pi^{-}+\mathrm{p}^{0}$

$$
\rightarrow \Sigma^{-}+\pi^{+}+\pi^{0}
$$




$T_{\pi^{+}}(\mathrm{Mev})$


Adjusted number of events


Adjusted number of events




[^0]:    ${ }^{F}$ Work done under the auspices of the U.S. Aromic Energy Commiseion.
    ${ }^{\$}$ Presently at Laboratoire de Physique Atomique, College de France. Paris, France. **Presently at University of Wisconsin, Madison, Wisconsin.
    ${ }^{+t} p_{r e s e n t l y}$ at the University of Callfornia at Los Angeles. Los Angeles, California

